

... for a brighter future

## **Baryon chiral extrapolation** in SU(3)

Ross D. Young

Chiral Dynamics 2009, University of Bern, Switzerland 6-10 July 2009

#### see Young & Thomas, arXiv:0901.3310



Argonne

**U.S.** Department



A U.S. Department of Energy laboratory managed by UChicago Argonne, LLC

#### Outline

- Motivating the physics
- Chiral extrapolation problem
  - 3-flavour treatment of effective field theory (EFT)
  - poor convergence of SU(3) expansion
- Solution: Finite-Range Regularisation (FRR)
  - constrained by lattice
- New determination of the strangeness sigma term





HPQCD / UKQCD / MILC / Fermilab, PRL92,022001(2004)



## 2003

#### "High-Precision Lattice QCD Confronts Experiment"



HPQCD / UKQCD / MILC / Fermilab, PRL92,022001(2004)



## 1999

#### "Quenched Light Hadron Spectrum"



CP-PACS, PRL84,238(2000)



## **Ratio plot - Quenched QCD**





## **Ratio plot - Quenched QCD**

#### New (2004) FLIC fermion results





## **Ratio plot - Quenched QCD**

#### New (2004) FLIC fermion results





## Latest results in lattice QCD

## Ab Initio Determination of Light Hadron Masses

S. Dürr,<sup>1</sup> Z. Fodor,<sup>1,2,3</sup> J. Frison,<sup>4</sup> C. Hoelbling,<sup>2,3,4</sup> R. Hoffmann,<sup>2</sup> S. D. Katz,<sup>2,3</sup> S. Krieg,<sup>2</sup> T. Kurth,<sup>2</sup> L. Lellouch,<sup>4</sup> T. Lippert,<sup>2,5</sup> K. K. Szabo,<sup>2</sup> G. Vulvert<sup>4</sup>

More than 99% of the mass of the visible universe is made up of protons and neutrons. Both particles are much heavier than their quark and gluon constituents, and the Standard Model of particle physics should explain this difference. We present a full ab initio calculation of the masses of protons, neutrons, and other light hadrons, using lattice quantum chromodynamics. Pion masses down to 190 mega—electron volts are used to extrapolate to the physical point, with lattice sizes of approximately four times the inverse pion mass. Three lattice spacings are used for a continuum extrapolation. Our results completely agree with experimental observations and represent a quantitative confirmation of this aspect of the Standard Model with fully controlled uncertainties.

21 NOVEMBER 2008 VOL 322 SCIENCE www.sciencemag.org



## Latest results in lattice QCD



21 NOVEMBER 2008 VOL 322 SCIENCE www.sciencemag.org



## Latest results in lattice QCD



21 NOVEMBER 2008 VOL 322 SCIENCE www.sciencemag.org



## New (public) lattice results





#### **Octet-baryon masses**

 $N,\Lambda,\Sigma,\Xi$  SU(2) symmetry

- SU(3) chiral limit:  $m_u = m_d = m_s = 0$ 
  - Octet are degenerate (one mass):  $M_0$
- Chiral EFT perturbs about this limit
  - Leading term is a single insertion of quark mass operator
    - 3 possible ways:



– Lagrangian (3 parameters)

 $\mathcal{L}_{BBq} = 2\alpha_M(\bar{B}B\mathcal{M}) + 2\beta_M(\bar{B}\mathcal{M}B) + 2\sigma_M\bar{B}B\operatorname{Tr}(\mathcal{M})$ 



#### **Octet-baryon masses**

• Leading-order expansion O(1)

$$M_{N} = M_{0} + 2(\alpha_{M} + \beta_{M})m_{q} + 2\sigma_{M}(2m_{q} + m_{s})$$

$$M_{\Lambda} = M_{0} + (\alpha_{M} + 2\beta_{M})m_{q} + \alpha_{M}m_{s} + 2\sigma_{M}(2m_{q} + m_{s})$$

$$M_{\Sigma} = M_{0} + \frac{1}{3}(5\alpha_{M} + 2\beta_{M})m_{q} + \frac{1}{3}(\alpha_{M} + 4\beta_{M})m_{s} + 2\sigma_{M}(2m_{q} + m_{s})$$

$$M_{\Xi} = M_{0} + \frac{1}{3}(\alpha_{M} + 4\beta_{M})m_{q} + \frac{1}{3}(5\alpha_{M} + 2\beta_{M})m_{s} + 2\sigma_{M}(2m_{q} + m_{s})$$

$$m_{\pi}^2 = 2Bm_q \quad m_K^2 = B(m_q + m_s)$$

$$m_q \to \frac{m_\pi^2}{2B}, \quad m_s \to \frac{2m_K^2 - m_\pi^2}{2B} \qquad \{\alpha, \beta, \sigma\} \to B\{\alpha', \beta', \sigma'\}$$



#### **Beyond first derivative: Loop corrections**

• At O(3/2) contributions from meson dressing



• Up to overall factors, (HB) loop integral reduces to

$$\int_{0}^{\infty} dk \frac{k^{4}}{k^{2} + m_{\pi}^{2}} = \int_{0}^{\infty} dk \frac{k^{4} - m_{\pi}^{4} + m_{\pi}^{4}}{k^{2} + m_{\pi}^{2}}$$
$$= \int dk \left(k^{2} - m_{\pi}^{2}\right) + \left(\int dk \frac{m_{\pi}^{4}}{k^{2} + m_{\pi}^{2}}\right) = \left(\frac{\pi}{-m_{\pi}^{3}}\right)$$

- RENORMALIZATION:
  - Absorb divergences into LECs



#### What about Decuplet baryons?

- EFT integrates out all non-dynamical degrees of freedom

   EFT can only be valid at energy scales below any physical threshold that has been integrated out
- Decuplet-less EFT cannot describe meson masses greater than Octet-Decuplet splitting
- Physical quark masses:  $M_\Delta M_N \sim 0.3\,{
  m GeV}$   $m_K,\,m_\eta \sim 0.5\,{
  m GeV}$ 
  - For physical strange-quark mass, an EFT that includes all dynamical degrees of freedom must include Decuplet
- I don't know if there is a unique way to include Decuplet
  - Including Decuplet: perhaps model dependent
  - Not including Decuplet: no longer an EFT



#### My power-counting for including Decuplet

- Assume Octet and Decuplet-baryons degenerate
- Include all (Octet and Decuplet) loop contributions to a given order in the quark mass
- Evaluate Decuplet loop integrals with explicit mass-splitting in the relevant propagators
  - Ensures physical threshold and IR branch structure maintained
- Renormalize such that the Decuplet-less EFT is recovered in the limit  $m_{PS} \rightarrow 0 \ll \Delta$



#### My power-counting for including Decuplet

- Assume Octet and Decuplet-baryons degenerate
- Include all (Octet and Decuplet) loop contributions to a given order in the quark mass
- Evaluate Decuplet loop integrals with explicit mass-splitting in the relevant propagators
  - Ensures physical threshold and IR branch structure maintained
- Renormalize such that the Decuplet-less EFT is recovered in the limit  $m_{PS} \rightarrow 0 \ll \Delta$

I am not claiming this is better than anyone else's method for including decuplet



### **Everything to O(3/2)**





### **Everything to O(3/2)**





#### **Correct for lattice volume**

- Same loop integrals describe leading finitevolume correction
- FV corrections purely infrared – should not be sensitive to UV regularisation
- Error estimate reflects FV corrections evaluated with and without a cutoff





#### **Power counting estimate for O(2)**



#### **Power counting estimate for O(2)**

• If we adopt conventional wisdom "4 pi fpi"

- Physical point

$$\mathcal{O}(2) \sim \frac{m_{\eta}^4}{(4\pi f_{\pi})^4} \sim 5\%$$



#### **Power counting estimate for O(2)**

- If we adopt conventional wisdom "4 pi fpi"
  - Physical point

$$\mathcal{O}(2) \sim \frac{m_{\eta}^4}{(4\pi f_{\pi})^4} \sim 5\%$$

Lattice masses

$$\mathcal{O}(2) \sim \frac{m_{\eta}^4}{(4\pi f_{\pi})^4} \sim 11\%$$



#### Lattice Simulation Results: LHPC





#### Best fit to lightest 2 quark masses

- Poor fit  $m_\pi \lesssim 0.35 \,{
  m GeV}$  $\chi^2/{
  m dof} \sim 40$   $m_K \lesssim 0.6 \,{
  m GeV}$
- "Best" fit  $M_0 \sim 0.27 \,\mathrm{GeV}$
- Empirical suggestion

 $\mathcal{O}(2) \sim \left(\frac{m_{\eta}}{\Lambda_B}\right)^4 \sim 300\%$ 

 $\Lambda_B \sim 0.6 \,\mathrm{GeV}$ 



#### Best fit to lightest 2 quark masses

• Poor fit  $\gamma \chi^2/{
m dof} \sim 40$ 

 $m_\pi \lesssim 0.35 \,\mathrm{GeV}$  $m_K \lesssim 0.6 \,\mathrm{GeV}$ 

- "Best" fit  $M_0 \sim 0.27 \,\mathrm{GeV}$
- Empirical suggestion

$$\mathcal{O}(2) \sim \left(\frac{m_{\eta}}{\Lambda_B}\right)^4 \sim 300\%$$

 $\Lambda_B \sim 0.6 \,\mathrm{GeV}$ 





#### What about Finite-Range Regularisation (FRR)?

- Introduce a resummation of higher-order terms with a single parameter
- Chiral loop integrals modified to cut off divergences

$$\int_0^\infty dk \, \frac{k^4}{k^2 + m^2} \left(\frac{\Lambda^2}{\Lambda^2 + k^2}\right)^4$$

Upon renormalisation gives identical expansion to O(3/2)

Text book:  $M_B^{(3/2)} = M_0 + \delta M^{(1)} + \delta M^{(3/2)} + 0$ 

FRR: 
$$M_B^{(3/2)} = M_0 + \delta M^{(1)} + \delta M^{(3/2)} + \mathcal{O}(\frac{m_{PS}^4}{\Lambda})$$



#### **Regularisation parameter?**

- Model-indepence of EFT only exists if results independent of this cutoff
- Can the lattice results select a preferred scale to regularise the EFT?



#### **Regularisation parameter?**

- Model-indepence of EFT only exists if results independent of this cutoff
- Can the lattice results select a preferred scale to regularise the EFT?





#### Fits to 2 lightest LHPC points





#### Meson masses - LHPC





#### Fits to 2 lightest LHPC points





#### Fits to 2 lightest LHPC points





#### More new lattice results: PACS-CS





#### Fit to 2 (blue) PACS-CS points - fixed strange mass





#### Fit to 2 (blue) PACS-CS points - fixed strange mass





#### **Consistency in LECs?**



Consistency of lattice at finite "a"

	$M_0$	$lpha_M$	$\beta_M$	$\sigma_M$
LHPC	0.82(6)	1.69(31)	1.29(26)	0.56(13)
PACS-CS	0.83(8)	1.64(33)	1.17(26)	0.52(15)



#### Test of scale determination





 $0.9 \ 1 \ 1.1$ 

#### Test of scale determination



# Excellent confirmation of consistent scale determination



• Gell-Mann–Okubo Relation and Pion-Nucleon sigma term  $M_N, M_\Lambda, M_\Sigma, M_\Xi$ 

## $m_s \langle N | \bar{s}s | N \rangle \simeq 335 \pm 132 \,\mathrm{MeV}$ Nelson & Kaplan PLB(1987)

 $\sim M_N^{phys} - M_N^{SU(3)chiral limit}$ 



• Gell-Mann–Okubo Relation and Pion-Nucleon sigma term  $M_N, M_\Lambda, M_\Sigma, M_\Xi$ 

$$m_s \langle N | \bar{s}s | N \rangle \simeq 335 \pm 132 \, {
m MeV}$$
  
Nelson & Kaplan PLB(1987)

 $\sim M_N^{phys} - M_N^{SU(3)chiral limit}$ 

QCD Lagrangian  $\sim \dots \bar{s}(D + m_s)s$ 

$$m_s \langle N | \bar{s}s | N \rangle = m_s \frac{\partial M_N}{\partial m_s}$$

evaluated at physical point!



• Gell-Mann–Okubo Relation and Pion-Nucleon sigma term  $M_N, M_\Lambda, M_\Sigma, M_\Xi$ 

$$m_s \langle N | \bar{ss} | N \rangle \simeq 335 \pm 132 \,\text{MeV}$$
  
Nelson & Kaplan PLB(1987)  
 $\sim M_N^{phys} - M_N^{SU(3)chiral \ limit}$   
Lagrangian  $\sim \dots \bar{s}(D + m_s)s$   
 $m_s \langle N | \bar{ss} | N \rangle = m_s \frac{\partial M_N}{\partial m_s}$   
evaluated at *physical* point!

Improved Effective Field Theory estimate

$$m_s \frac{\partial M_N}{\partial m_s} = 113 \pm 108 \,\mathrm{MeV}$$

Borasoy & Meissner (1997)



QCD

• Gell-Mann–Okubo Relation and Pion-Nucleon sigma term  $M_N, M_\Lambda, M_\Sigma, M_\Xi$ 

$$m_s \langle N | \bar{s}s | N \rangle \simeq 335 \pm 132 \,\text{MeV}$$
  
Nelson & Kaplan PLB(1987)  
 $\sim M_N^{phys} - M_N^{SU(3)chiral\ limit}$   
QCD Lagrangian  $\sim \dots \bar{s}(D + m_s)s$   
 $m_s \langle N | \bar{s}s | N \rangle = m_s \frac{\partial M_N}{\partial m_s}$   
evaluated at *physical* point!

Improved Effective Field Theory estimate

Lattice?  $\left(m_s \frac{\partial M_N}{\partial m_s}\right) = 113 \pm 108 \,\mathrm{MeV}$ Borasoy & Meissner (1997)



#### **Beyond the masses**

- Absolute masses competitive precision with recent Science article
- Can determine the sensitivity of observables to strange-quark mass
  - Important for lattice QCD: fine-tuning the strange-quark mass is computationally expensive
- Can extract strangeness nucleon sigma term





#### **Direct dark matter detection**

- "Hadronic uncertainties in the elastic scattering of supersymmetric dark matter", Ellis *et al.* PRD77(2008)
- Spin-independent neutralino scattering cross section

$$\sigma_{\rm SI} = \frac{4m_r^2}{\pi} [Zf_p + (A - Z)f_n]^2,$$

$$\frac{f_N}{m_N} = \sum_{q=u,d,s} f_{T_q}^{(N)} \frac{\alpha_{3q}}{m_q} + \frac{2}{27} f_{TG}^{(N)} \sum_{q=c,b,t} \frac{\alpha_{3q}}{m_q}$$

• Light quarks:  $\alpha_{3q} \sim cm_q$ 

$$\bar{\sigma}_{\{u,d,s\}} = f_{T\{u,d,s\}}^{(p)} \sim \{0.027, 0.039, 0.363\}$$



#### **Direct dark matter detection**

- "Hadronic uncertainties in the elastic scattering of supersymmetric dark matter", Ellis *et al.* PRD77(2008)
- Spin-independent neutralino scattering cross section

$$\sigma_{\rm SI} = \frac{4m_r^2}{\pi} [Zf_p + (A - Z)f_n]^2,$$

$$\frac{f_N}{m_N} = \sum_{q=u,d,s} f_{T_q}^{(N)} \frac{\alpha_{3q}}{m_q} + \frac{2}{27} f_{TG}^{(N)} \sum_{q=c,b,t} \frac{\alpha_{3q}}{m_q}$$

• Light quarks:  $\alpha_{3q} \sim cm_q$ 

$$\bar{\sigma}_{\{u,d,s\}} = f_{T\{u,d,s\}}^{(p)} \sim \{0.027, 0.039, 0.363\}$$
  
New result: ~0.033

please wait: Giedt et al.



#### **Summary**

- Robust chiral extrapolation
  - Excellent precision in absolute mass determination
  - Confirmation of scale determination efforts by HPQCD/MILC etc.
  - Moderates the fine-tuning problem of the strange quark mass
- Strangeness sigma term is small
- New opportunities for precision baryon studies in lattice QCD



#### Finite-volume correction (estimated in EFT)

Lightest LHPC simulation



