# Determination of Low Energy Constants and Testing Chiral Perturbation Theory at Next to Next to Leading Order

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## Why are we looking for relations between observables?

Chiral Perturbation Theory  $\rightarrow$  every observable can be written as a sum of terms of increasing importance in the Chiral expansion.

$$O = O^{(2)} + O^{(4)} + O^{(6)}$$

The  $p^6$  part can be split as

$$O^{(6)} = O_{C_i(\text{tree level})} + O_{L_i(\text{one loop})} + O_{F_0(\text{two loops})}$$

We look for relations between observables such that the first contribution cancels out. Using these

- we can check how large is the loop contribution and test ChPT convergence in a *C<sub>i</sub>* independent way
- we hoped to perform a fit of the *L<sub>i</sub>* at NNLO not depending on the *C<sub>i</sub>*. Unfortunately in most of the relations the NLO *L<sub>i</sub>* contributions cancel too (the dependence on the *L<sub>i</sub>* is only through the NNLO pieces)
- in this way we isolated combinations of the  $C_i$

# Overview of the processes considered and relations found

process	# observables	# relations
$\pi\pi$ scattering	11	5
$\pi K$ scattering	14	5
$\pi K$ and $\pi \pi$ scattering	no extra observables	2
$K_{\ell 4}$ (with $\pi K$ scattering)	10	1
$\eta \rightarrow 3\pi \text{ (with } \pi K)$	6	2
scalar form factors $F_S^{\pi/K}(t)$	18	6
$F_S^{\pi/K}(t)$ , $\pi\pi$ and $\pi K$ scattering	no extra observables	2
$F_S^{\pi/K}(t), K_{\ell 4}, \pi \pi$ and $\pi K$ scattering	no extra observables	1
$F_{S}^{\pi/K}(t)$ , masses and decay constants	6	4
Vector form factors $F_V^{\pi/K}$	11	7
Total	76	35



3 Summary and Future Steps

### Numerical analysis explanation

- Evaluation of each side of the relation using experimental data and/or dispersive analysis:
   [CGL] G. Colangelo, J. Gasser and H. Leutwyler, *Nucl. Phys.* B 603 (2001) 125 (ππ scattering)
   [BDM] Büttiker, Descotes-Genon, Moussallam *Eur. Phys. J.* C 33 (2004) 409 (πK scattering)
   [NA48/2] NA48/2 coll., *Eur. Phys. J.* C 54 (2008) 411-423 (*K*<sub>ℓ4</sub>)
   [E865] S. Pislak *et al.*, *Phys. Rev.* D 67 (2003) 072004 (*K*<sub>ℓ4</sub>)
- Solution Using ChPT up to  $p^6$  results;  $L_i = \text{fit10}$  and  $C_i = 0$ . For references see J. Bijnens, Prog. Part. Nucl. Phys. **58** (2007) 521
- Solution We quote the difference of the two  $\Rightarrow$  it contains only the  $p^6$  piece coming from the  $C_i$  and higher order terms.
- Errors obtained adding in quadrature the uncertainties from experiments/dispersive results. No theoretical uncertainty due to the values of L<sub>i</sub> or to higher orders has been added

• 
$$A(\pi^a \pi^b \to \pi^c \pi^d) = \delta^{a,b} \delta^{c,d} A(s,t,u) + \delta^{cd} \delta^{bd} A(t,u,s) + \delta^{ad} \delta^{bc} A(u,t,s)$$

• The isospin amplitudes  $T^{I}(s,t)$  (I = 0, 1, 2) are written in terms of the function A(s, t, u) and then expanded in partial waves:

$$T^{I}(s,t) = 32\pi \sum_{\ell=0}^{+\infty} (2\ell+1)P_{\ell}(\cos\theta)t^{I}_{\ell}(s)$$
  
Near threshold  $\rightarrow t^{I}_{\ell}(s) = q^{2\ell}(a^{I}_{\ell} + b^{I}_{\ell}q^{2} + \mathcal{O}(q^{4}))$ 
$$q^{2} = \frac{1}{4}(s - 4m_{\pi}^{2}) \qquad a^{I}_{\ell}, b^{I}_{\ell} \dots = \text{scattering lengths, slopes, }\dots$$

• We studied only those observables where a dependence on the  $C_i$  shows up  $\rightarrow$  11 threshold parameters

• A(s, t, u) can be written in terms of 6 independent parameters

$$A(s,t,u) = b_1 + b_2 s + b_3 s^2 + b_4 (t-u)^2 + b_5 s^3 + b_6 s (t-u)^2$$
  
+non polynomial part

- $\Rightarrow$  5 relations among the scattering lengths.
- They hold for  $n_f = 2$ , 3, at NLO and NNLO: not only the  $p^6$  LECs cancel out, but also the tree level part involving the  $p^4$  LECs does. Still there is  $L_i$  or  $l_i$  dependence through the non polynomial part

$$\left[5b_0^2 - 2b_0^0 - 27a_1^1 - 15a_0^2 + 6a_0^0\right]_{C_i} = -18\left[b_1^1\right]_{C_i}$$
(1)

$$\left[3a_1^1 + b_0^2\right]_{C_i} = 20\left[b_2^2 - b_2^0 - a_2^2 + a_2^0\right]_{C_i}$$
(2)

$$\left[b_0^0 + 5b_0^2 + 9a_1^1\right]_{C_i} = 90 \left[a_2^0 - b_2^0\right]_{C_i}$$
(3)

 $a_\ell^I(b_\ell^I)$  expressed in unit of  $m_\pi^{2\ell}(m_\pi^{2\ell+2})$ 

	[CGL]	NLO	NLO	NNLO	NNLO	remainder
		1-loop	LECs	2-loop	1-loop	
LHS (1)	$0.009 \pm 0.039$	0.054	-0.044	-0.041	-0.002	$0.041 \pm 0.039$
RHS (1)	$-0.102 \pm 0.002$	-0.009	-0.044	-0.060	-0.008	$0.018 \pm 0.002$
10 LHS (2)	$0.334 \pm 0.019$	0.209	0.097	0.103	0.029	$-0.105 \pm 0.019$
10 RHS (2)	$0.322\pm0.008$	0.177	0.097	0.120	0.034	$-0.107 \pm 0.008$
LHS (3)	$0.216 \pm 0.010$	0.166	0.029	0.053	0.016	$-0.047 \pm 0.010$
RHS (3)	$0.189 \pm 0.003$	0.145	0.029	0.049	0.020	$-0.054 \pm 0.003$

	[CGL]	two-flavour	remainder
		[CGL]	
LHS (1)	$0.009 \pm 0.039$	-0.003	$0.007 \pm 0.039$
RHS (1)	$-0.102 \pm 0.002$	-0.097	$-0.005 \pm 0.002$
10 LHS (2)	$0.334 \pm 0.019$	0.332	$0.002 \pm 0.019$
10 RHS (2)	$0.322\pm0.008$	0.318	$0.004\pm0.075$
LHS (3)	$0.216\pm0.010$	0.206	$0.010\pm0.010$
RHS (3)	$0.189 \pm 0.003$	0.189	$0.000\pm0.003$

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$$\begin{bmatrix} 3b_1^1 + 25a_2^2 \end{bmatrix}_{C_i} = 10 \begin{bmatrix} a_2^0 \end{bmatrix}_{C_i}$$
(4)  
$$\begin{bmatrix} -5b_2^2 + 2b_2^0 \end{bmatrix}_{C_i} = 21 \begin{bmatrix} a_3^1 \end{bmatrix}_{C_i}$$
(5)

 $a_\ell^I(b_\ell^I)$  expressed in unit of  $m_\pi^{2\ell}(m_\pi^{2\ell+2})$ 

	[CGL]	NLO	NLO	NNLO	NNLO	remainder
		1-loop	LECs	2-loop	1-loop	
10 LHS (4)	$0.213 \pm 0.005$	0.137	0.032	0.053	0.035	$-0.043 \pm 0.005$
10 RHS (4)	$0.175 \pm 0.003$	0.121	0.032	0.050	0.029	$-0.057 \pm 0.003$
$10^3$ LHS (5)	$0.92\pm0.07$	0.36	0.00	0.56	-0.01	$0.00\pm0.07$
$10^3$ RHS (5)	$1.18\pm0.04$	0.42	0.00	0.57	0.03	$0.15\pm0.04$

	[CGL]	two-flavour	remainder
		[CGL]	
10 LHS (4)	$0.213\pm0.005$	0.204	$0.009\pm0.005$
10 RHS (4)	$0.175\pm0.003$	0.176	$-0.001 \pm 0.003$
$10^3$ LHS (5)	$0.92\pm0.07$	1.00	$-0.08\pm0.07$
$10^3$ RHS (5)	$1.18\pm0.04$	1.15	$0.04\pm0.04$

• Rel (4) and (5)  $\rightarrow$  ok at 2-sigma level

## $\pi K$ scattering: generalities

- $T^{I}(s, t, u) =$  scattering amplitude in isospin channel  $I = \frac{1}{2}, \frac{3}{2}$
- As for the  $\pi\pi$  scattering, it's possible to define scattering lengths  $a_{\ell}^{I}, b_{\ell}^{I}$ :

$$T^{I}(s,t,u) = 16\pi \sum_{\ell=0}^{+\infty} (2\ell+1) P_{\ell}(\cos\theta) t^{I}_{\ell}(s)$$
  
Near threshold  $\rightarrow t^{I}_{\ell} = \frac{1}{2} \sqrt{s} q^{2\ell}_{\pi K} (a^{I}_{\ell} + b^{I}_{\ell} q^{2}_{\pi K} + \mathcal{O}(q^{4}_{\pi K}))$ 
$$q^{2}_{\pi K} = \frac{s}{4} \left( 1 - \frac{(m_{K} + m_{\pi})^{2}}{s} \right) \left( 1 - \frac{(m_{K} - m_{\pi})^{2}}{s} \right)$$
$$t = -2q^{2}_{\pi K} (1 - \cos\theta), \quad u = -s - t + 2m^{2}_{K} + 2m^{2}_{\pi}$$

• Again we studied only those scattering lengths where a dependence on the  $C_i$  shows up  $\rightarrow$  14 threshold parameters

### $\pi K$ scattering: relations

On the other hand the isospin amplitudes T<sup>I</sup>(s, t, u) are written in terms of the crossing symmetric and antisymmetric amplitudes T<sup>±</sup>(s, t, u) which can be expanded around t=0, s=u (ν = s-u/4m<sub>κ</sub>) (subthreshold expansion):

$$T^{+}(s,t,u) = \sum_{i,j=0}^{\infty} c_{ij}^{+} t^{i} \nu^{2j} \qquad T^{-}(s,t,u) = \sum_{i,j=0}^{\infty} c_{ij}^{-} t^{i} \nu^{2j+1}$$

where  $16\rho^2 \left[c_{\overline{20}}\right]_{C_i} = 3 \left[c_{\overline{01}}\right]_{C_i}$ ,  $\rho = m_K/m_\pi$  and  $c_{ij}$  are expressed in unit of  $m_\pi^{2i+2j+1}$ 

- $\rightarrow$  9 independent subthreshold parameters.
- $\Rightarrow$  5 relations between the scattering lengths holding both at  $p^6$  and at  $p^4$ : no dependence of the  $L_i$  at NLO.

For simplicity we introduce the notation

$$\begin{aligned} a_{\ell}^{-} &= a_{\ell}^{1/2} - a_{\ell}^{3/2} \qquad b_{\ell}^{-} = b_{\ell}^{1/2} - b_{\ell}^{3/2} \\ a_{\ell}^{+} &= a_{\ell}^{1/2} + 2a_{\ell}^{3/2} \qquad b_{\ell}^{+} = b_{\ell}^{1/2} + 2b_{\ell}^{3/2} \end{aligned}$$

$$\left(\rho^{4} + 3\rho^{3} + 3\rho + 1\right) \left[a_{1}^{-}\right]_{C_{i}} = 2\rho^{2} \left(\rho + 1\right)^{2} \left[b_{1}^{-}\right]_{C_{i}} - \frac{2}{3}\rho \left(\rho^{2} + 1\right) \left[b_{0}^{-}\right]_{C_{i}} + \frac{1}{2\rho} \left(\rho^{2} + \frac{4}{3}\rho + 1\right) \left(\rho^{2} + 1\right) \left[a_{0}^{-}\right]_{C_{i}}$$

$$(6)$$

$$5\left(\rho^{2}+1\right)\left[a_{2}^{-}\right]_{C_{i}}=\left[a_{1}^{-}\right]_{C_{i}}+2\rho\left[b_{1}^{-}\right]_{C_{i}}$$

$$(7)$$

$$5(\rho+1)^{2} \left[b_{2}^{-}\right]_{C_{i}} = \frac{(\rho-1)^{2}}{\rho^{2}} \left[a_{1}^{-}\right]_{C_{i}} - \frac{\rho^{4} + \frac{2}{3}\rho^{2} + 1}{4\rho^{4}} \left[a_{0}^{-}\right]_{C_{i}} + \frac{\rho^{2} - \frac{2}{3}\rho + 1}{2\rho^{2}} \left[b_{0}^{-}\right]_{C_{i}}$$
(8)

All quantities are in the units of powers of  $m_{\pi^+}$ 

	[BDM]	NLO	NLO	NNLO	NNLO	remainder
		1-loop	LECs	2-loop	1-loop	
LHS (6)	$5.4 \pm 0.3$	0.16	0.97	0.77	-0.11	$0.6 \pm 0.3$
RHS (6)	$6.9\pm0.6$	0.42	0.97	0.77	-0.03	$1.8\pm0.6$
10 LHS (7)	$0.32\pm0.01$	0.03	0.12	0.11	0.00	$0.07\pm0.01$
10 RHS (7)	$0.37\pm0.01$	0.02	0.12	0.10	-0.01	$0.14\pm0.01$
100 LHS (8)	$-0.49\pm0.02$	0.08	-0.25	-0.17	0.05	$-0.21\pm0.02$
100 RHS (8)	$-0.85\pm0.60$	0.03	-0.25	0.11	-0.03	$-0.71 \pm 0.60$

- Rel (6)  $\rightarrow$  ok at 2 sigma
- Rel  $(7) \rightarrow ok$  with a theoretical error about half the NNLO contribution
- Rel (8)  $\rightarrow$  ok but large uncertainty

$$7\left(\rho^{2}+1\right)\left[a_{3}^{-}\right]_{C_{i}} = \left[a_{2}^{-}\right]_{C_{i}}+2\rho\left[b_{2}^{-}\right]_{C_{i}}$$
(9)  

$$7\left[a_{3}^{+}\right]_{C_{i}} = \frac{1}{2\rho}\left[a_{2}^{+}\right]_{C_{i}}-\left[b_{2}^{+}\right]_{C_{i}}+\frac{1}{5\rho}\left[b_{1}^{+}\right]_{C_{i}}-\frac{1}{60\rho^{3}}\left[a_{0}^{+}\right]_{C_{i}}$$
(10)

All quantities are in the units of powers of  $m_{\pi^+}$ 

	[BDM]	NLO	NLO	NNLO	NNLO	remainder
		1-loop	LECs	2-loop	1-loop	
100 LHS (9)	$0.13\pm0.01$	0.04	0.00	0.01	0.03	$0.05\pm0.01$
100 RHS (9)	$0.01\pm0.01$	0.01	0.00	0.00	0.00	$-0.01\pm0.01$
$10^3$ LHS (10)	$0.29\pm0.05$	0.09	0.00	0.06	0.01	$0.13\pm0.03$
$10^3$ RHS (10)	$0.31\pm0.07$	0.03	0.00	0.06	0.05	$0.17\pm0.07$

• Rel (9)  $\rightarrow$  ChPT seems to underestimate  $a_3^-$ 

#### $\pi\pi$ scattering and $\pi K$ scattering

Considering  $\pi\pi$  and  $\pi K$  scattering together two more relations appear. These are due to the following identities:

$$[b_5]_{C_i} = [c_{30}^+]_{C_i} + \frac{3}{4\rho} [c_{20}^-]_{C_i}, \qquad [b_6]_{C_i} = \frac{1}{4\rho} [c_{20}^-]_{C_i} + \frac{1}{16\rho^2} [c_{11}^+]_{C_i}$$

which in terms of the threshold parameters read

$$6 \left[ a_3^1 \right]_{C_i} = (1+\rho) \left[ a_3^+ + 3a_3^- \right]_{C_i}$$
(11)

$$3\left[(1+\rho)^{2}\left[b_{2}^{2}\right]_{C_{i}}+7\left(1-\rho\right)^{2}\left[a_{3}^{1}\right]_{C_{i}}\right] = (1+\rho)\left[7\left(1-4\rho+\rho^{2}\right)\left[a_{3}^{-}\right]_{C_{i}}+\left[a_{2}^{+}+2\rho b_{2}^{+}\right]_{C_{i}}\right]$$
(12)

All quantities in units of powers of  $m_{\pi^+}$ 

	[CGL]	NLO	NLO	NNLO	NNLO	remainder
	[BDM]	1-loop	LECs	2-loop	1-loop	
$10^3$ LHS (11)	$0.34\pm0.01$	0.12	0.00	0.16	0.00	$0.05\pm0.01$
$10^3$ RHS (11)	$0.38\pm0.03$	0.12	0.00	0.05	0.04	$0.16\pm0.03$
10 LHS (12)	$-0.13 \pm 0.01$	-0.12	0.00	-0.05	0.02	$0.01\pm0.01$
10 RHS (12)	$-0.09\pm0.02$	-0.05	0.00	-0.02	-0.01	$-0.01\pm0.02$

• Rel (11)  $\rightarrow a_3^-$  appears similar discrepancy seen in Kampf, Moussallam Eur. Phys. J. C. 47 (2006) 723 (see talk by Bijnens)

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### $K_{\ell 4}$ : generalities and relation

- In the transition amplitude 4 form factors appear: F, G, H, R (R in  $K_{e4}$  is suppressed  $\rightarrow$  only in  $K_{\mu4}$ )
- Using partial wave expansion and neglecting *d* wave terms (10 observables):

$$F = f_s + f'_s q^2 + f''_s q^4 + f'_e s_e / 4m_\pi^2 + f_t \sigma_\pi X \cos \theta + \dots ,$$
  

$$G_p = g_p + g'_p q^2 + g''_g q^4 + g'_e s_e / 4m_\pi^2 + g_t \sigma_\pi X \cos \theta + \dots$$

 $s_{\pi}(s_e) =$ invariant mass of dipion (dilepton)  $q^2 = (s_{\pi}/(4m_{\pi}^2) - 1)$ 

• 1 relation between  $\pi K$  scattering and  $K_{e4}$  observables:

$$\sqrt{2} \left[ f_s'' \right]_{C_i} = 64 \rho F_\pi \left[ c_{30}^+ \right]_{C_i}$$

in terms of the  $\pi K$  threshold parameters reads

$$\sqrt{2} \left[ f_s'' \right]_{C_i} = 32\pi \frac{\rho}{1+\rho} F_{\pi} \left[ \frac{35}{6} \left( 2+\rho+2\rho^2 \right) \left[ a_3^+ \right]_{C_i} - \frac{5}{4} \left[ a_2^+ + 2\rho b_2^+ \right]_{C_i} \right]$$

All quantities are in units of powers of  $m_{\pi^{+}}$ ,

#### $K_{\ell 4}$ : numerics

#### (see talk by Bijnens)

	[BDM], [E865],	NLO	NLO	NNLO	NNLO	remainder
	[NA48/2]	1-loop	LECs	2-loop	1-loop	
LHS (17)	$-0.73 \pm 0.10$	-0.23	0.00	-0.15	-0.05	$-0.29\pm0.10$
RHS (17)	$0.50\pm0.07$	0.19	0.00	0.10	0.03	$0.18\pm0.07$





3 Summary and Future Steps

- many observables at NNLO, depending on many correlated LECs
- we found relations among observables not depending on the NNLO constants (and most of them not depending on the NLO either) → a way to check ChPT
- although many relations work well, results for  $K_{\ell 4}$  and  $\pi K$  scattering show discrepancies: further investigation needed
- arXiv:0906.3118 [hep-ph]

Future steps:

- new fit of the  $L_i$  with a better treatment of the  $C_i$  and using new exp data available, dispersive analysis and lattice results (masses, scalar form factors)
- include corrections (e.g. isospin breaking)

# Status of the $L_i$ fit

# PRELIMINARY!!!!

- Program for fitting (Minuit) almost ready (pion scalar radius,  $\pi\pi$  and  $\pi K$  scattering observables now included)
- *C<sub>i</sub>* obtained through Resonance Estimates (Vector, Scalar, PseudoScalar) (see Amoros, Bijnens, Talavera, Nucl. Phys. B 602 (2001) 87 )

	fit 10 iso	NA48	$F_K/F_{\pi}$	All
$10^{3}L_{1}^{r}$	$0.40\pm0.12$	0.98	0.97	$0.99\pm0.13$
$10^{3}L_{2}^{r}$	$0.76\pm0.12$	0.78	0.79	$0.60\pm0.22$
$10^{3}L_{3}^{\overline{r}}$	$-2.40\pm0.37$	-3.14	-3.12	$-3.07\pm0.59$
$10^{3}L_{4}^{r}$	$\equiv 0$	$\equiv 0$	$\equiv 0$	$0.65\pm0.64$
$10^{3}L_{5}^{r}$	$0.97\pm0.11$	0.93	0.72	$0.53\pm0.10$
$10^{3}L_{6}^{r}$	$\equiv 0$	$\equiv 0$	$\equiv 0$	$0.07\pm0.65$
$10^{3}L_{7}^{r}$	$-0.30\pm0.15$	-0.30	-0.26	$-0.21\pm0.15$
$10^{3}L_{8}^{r}$	$0.61\pm0.20$	0.59	0.48	$0.37\pm0.17$
$\chi^2$ (dof)	0.25(1)	0.17(1)	0.19(1)	0.78 (4)

- NA48  $\rightarrow$  NA48 exp data. No change in the fit including curvature:  $f_s'' = -0.90$  (exp value:  $f_s'' = -1.58 \pm 0.064$ )
- $F_K/F_{\pi} = 1.19$  (value of  $L_5^r$  changes)
- All  $\rightarrow$  add  $a_0^0, a_0^2, a_0^{1/2}, a_0^{3/2}$ , scalar pion radius