# Chiral low-energy couplings from lattice computations in the epsilon-regime

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Working Group 1 - Goldstone bosons: Light quark parameters

## Outline



• Chiral effective theory in the  $\epsilon$ -regime

#### 2 Lattice QCD in the epsilon regime

#### Results

- LECs from Dirac eigenvalues distribution
- 2-point correlators

### 4 Conclusions

### Introduction

Match lattice QCD with the chiral effective theory:

- Test if chiral symmetry is broken according to the expected pattern
- Extract the Low Energy Couplings (LECs) of the effective theory from first principles
- Ideally:  $a \rightarrow 0, m \rightarrow 0, V \rightarrow \infty$
- Practically: simulations with  $N_f = 2, 2 + 1$  light quarks

 $a\simeq 0.05-0.12$  fm,  $M_\pi\simeq 200-500$  MeV ,  $L\sim 1.7-3$  fm,  $M_\pi L\simeq 3-4.$ 

Need for good control over systematic effects:

- Iattice artefacts
- finite-volume effects
- higher order contributions in the effective theory

# Introduction

Chiral regimes at finite volume V:

#### p-regime

- $(\lim_{m\to 0})_{M_{\pi}L \geq \mathcal{O}(1)}$
- $L \gg 1/(4\pi F)$
- $m \sim p^2$ ,  $\frac{1}{L}$ ,  $\frac{1}{T} \sim p$
- finite-volume effects exponentially suppressed



L

#### *ϵ*-regime

- $(\lim_{m\to 0})_{m\Sigma V=\mathcal{O}(1)}$
- $L \gg 1/(4\pi F); \quad (M_{\pi}L)^2 < 1$
- zero modes contribution non-perturbative

• 
$$m \sim \epsilon^4$$
,  $\frac{1}{L}$ ,  $\frac{1}{T} \sim \epsilon$ 

 finite-volume effects sizeable (polynomial), mass effects suppressed



### *ϵ*-regime

Reorganization of the power counting:

 $\rightarrow$  at a given order, less LECs appear with respect to p-regime

 $\rightarrow$  In principle: ideal regime to compute LO LECs ( $\Sigma$ , F)

#### Chiral symmetry breaking in a finite volume V

For  $\mu = m\Sigma V \ll 1$ 

$$\begin{split} \mathcal{Z} &= \int_{\mathrm{SU}(\mathrm{N}_{\mathrm{f}})} dU_0 \exp\left[\mu \operatorname{ReTr} U_0\right]; \quad U_0 \to \text{ zero mode} \\ \Sigma(\mu) &= \frac{\Sigma}{N_f} \frac{\partial}{\partial \mu} \ln \mathcal{Z} \propto \Sigma \mu \end{split}$$

 $\Sigma(\mu)_{\mu 
ightarrow 0} 
ightarrow 0$ : no symmetry breaking at finite volume

but finite-size scaling gives informations about spontaneous symmetry breaking in infinite-volume

Hasenfratz & Leutwyler ('90)

## Two-point functions in $\epsilon$ -regime

Example: pseudoscalar and axial correlator at NLO

$$C_{P}^{ab}(t) = \frac{1}{L^{3}} \int d^{3}x \langle P^{a}(x)P^{b}(0) \rangle = \delta^{ab}\Sigma^{2} \left[ a_{P} + \frac{T}{F^{2}L^{3}}b_{P}h_{1}(t/T) \right]$$

$$C_{A}^{ab}(t) = \frac{1}{L^{3}} \int d^{3}x \langle A_{0}^{a}(x)A_{0}^{b}(0) \rangle = \delta^{ab}\frac{F^{2}}{V} \left[ a_{A} + \frac{T}{F^{2}L^{3}} \left( b_{A}h_{1}(t/T) + c_{A} \right) \right]$$

• 
$$h_1\left(\frac{t}{T}\right) = \frac{1}{2}\left[\left(\frac{t}{T} - \frac{1}{2}\right)^2 - \frac{1}{12}\right]$$
 parabolic time dependence  
•  $a_P, b_P, a_A, b_A, c_A$ : functions of  $\mu = m\Sigma V, L, T$ 

Hansen ('90)

 $\rightarrow$  only LO LECs F and  $\Sigma$  appear

 in the ε-regime, topology plays a relevant role: predictions at fixed topology

$$\mathcal{Z}_{
u} = \int_{\mathrm{U}(\mathbf{N}_{\mathrm{f}})} dU_{0} (\det U_{0})^{
u} \exp\left[\mu \mathrm{Re}\mathrm{Tr}U_{0}
ight]$$

 $ightarrow \Sigma_{
u}(\mu), \, C^{ab}_{P,
u}(t), C^{ab}_{A,
u}(t)...$ 

Leutwyler, Smilga ('92), Damgaard et al ('01,'02)

## Lattice QCD in the $\epsilon$ -regime

Simulate lattice QCD with  $L \gg 1/(4\pi F)$  and  $m\Sigma V \lesssim 1$ 

 $\rightarrow$  need small quark masses

Which discretization?

**Ginsparg-Wilson fermions** 

Ginsparg, Wilson ('82)

Neuberger, Fixed-point actions, Domain Wall with  $\textit{L}_5 \rightarrow \infty$ 

- © Chiral Symmetry preserved at finite lattice spacing
- © Topology naturally defined through index theorem
- © Dynamical simulations numerically challenging

 $\rightarrow$  many quenched simulations in the  $\epsilon$ -regime *P.* Hernández et al ('99), *T.* DeGrand ('01), *W.* Bietenholz et al ('04), *L.* Giusti et all ('04,'07)

→ still few dynamical simulations JLQCD, JLQCD + TWQCD ('07,'08), P. Hasenfratz et al ('07), DeGrand and Schaefer ('07)

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Which discretization?

#### Wilson fermions

- © Chiral symmetry explicitly broken at finite lattice spacing
- ③ Small quark masses problematic: Dirac spectrum not bounded from below. Empirical stability bound:  $m > m_{min}$ ,  $m_{min} \propto a/\sqrt{V}$

Del Debbio et al ('06)

- No natural definition of topological charge
- © Dynamical simulations feasible

# Lattice QCD in the $\epsilon$ -regime

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Possible solution: reweighting  $\rightarrow$  first simulations in the  $\epsilon$ -regime

A. Hasenfratz et al ('08)

For reweighting see also Lüscher and Palombi ('08)

#### Wilson Twisted-Mass fermions

 $\bullet\,$  Dirac spectrum bounded from below  $\rightarrow \epsilon\text{-regime}$  can be reached

K. Jansen et al (ETM) ('07,'08)

# Matching lattice QCD with chiral effective theory

- Ideally: first perform continuum extrapolation
   → then match with chiral effective theory
- Alternative: incorporate discretization effects in the chiral effective theory  $\rightarrow$  Symanzik approach

Symanzik ('75)

(Disadvantage: new unknown LECs)

- GW fermions: described by continuum ChPT; LECs( $a^2$ )
- Wilson-like fermions: WChPT, TMChpt

Sharpe & Singleton ('98), Rupak & Shoresh ('02), Sharpe & Wu ('04), Aoki & Bär ('04)...

Can be extented to the  $\epsilon$ -regime

O. Bär, S.N., S. Schaefer ('09), A. Shindler ('09)

for  $m \sim a\Lambda_{QCD}^2$  (GSM regime) : explicit breaking of chiral symmetry still dominated by quark mass  $\rightarrow$  lattice spacing effects suppressed (NNLO)

## Results: LECs from Dirac eigenvalues distribution

 $\chi$ PT at LO in the  $\epsilon$ -expansion  $\leftrightarrow$  Random Matrix Theory  $\downarrow$ probability distribution of single eigenvalues  $p_{k,\nu}(\zeta_k,\mu)$ ;  $\zeta_k = \lambda_k \Sigma V, \mu = m \Sigma V$   $\downarrow$ match low-lying spectrum of Dirac operator:  $\langle \zeta_k \rangle^{RMT} = \langle \lambda_k \rangle^{QCD} \Sigma V \rightarrow \text{extract } \Sigma$  (F)

Shuryak and Verbaarschot ('93), Damgaard ('98), Akemann and Damgaard ('98), Wilke et al ('98), Damgaard and Nishigaki ('98), Osborn et al ('99), Akemann and Damgaard ('08)

Possible issue on systematic errors:

for the distribution of single eigenvalues, higher order corrections are not known

Recent unquenched computations with GW fermions:

JLQCD+TWQCD ('07), DeGrand and Schaefer ('07), Lang et al ('06), P. Hasenfratz et al ('07)

## Results: LECs from Dirac eigenvalues distribution



P. Hasenfratz et al ('07-'09): Fixed-point action

 $N_{\rm f} = 2 + 1$ , a=0.129(5) fm,  $L \simeq 1.6$  fm,  $\mu_{u,d} = m_{u,d} \Sigma V \simeq 1.4$ ,  $\mu_{s} = m_{s} \Sigma V \simeq 12.3$ 

$$\Sigma^{\overline{MS}}(\mu = 2 \text{ GeV}) = (239(11) \text{ MeV})^3$$

## Results: LECs from 2-point correlators



NLO Fit:

 $\Sigma^{\overline{MS}}(\mu = 2 \text{ GeV}) = (239.8(4.0) \text{ MeV})^3; \quad F = 87.3(5.6) \text{ MeV}$ 

### Results: LECs from 2-point correlators



A. Hasenfratz, Hoffmann, Schaefer ('08)

NHYP Wilson fermions

 $N_{\rm f}=$  2,  $a\simeq$  0.115 fm,  $L\simeq$  1.84, 2.8 fm,

 $\mu = m\Sigma V \simeq 0.7 - 5$ ,  $\nu$  not fixed

• Continuum NLO fit: (L = 2.8 fm)

$$F = 90(4) \text{ MeV},$$
  
 $\Sigma^{\overline{MS}}(2 \text{ GeV}) = (248(6) \text{ MeV})^3$ 

• GSM\* NLO fit: (L = 2.8 fm)Include leading  $(O(a^2))$ correction at NLO

O. Bär, S.N., S. Schaefer('09)

 $\begin{array}{l} F = 88(3) \; \text{MeV} \\ \Sigma^{\overline{\text{MS}}}(2 \; \text{GeV}) = (249(4) \; \; \text{MeV})^3 \\ c_2 = 0.02(8) \; \text{GeV}^4 \end{array}$ 

## Summary of lattice results for $\boldsymbol{\Sigma}$



 $[\Sigma^{\overline{\text{MS}}}(2\text{GeV})]^{1/3}$ 

see FLAG report, in preparation

### Summary of lattice results for F



see FLAG report, in preparation

### Conclusions

Lattice simulations in the  $\epsilon$ -regime give important informations on low-energy QCD  $\rightarrow$  LECs

systematic errors are different with respect to infinite-volume determinations

#### Still numerically challenging

 $\rightarrow$  but simulations with Wilson-like fermions seem feasible

Lattice effects due to explicit breaking of chiral symmetry are under control, at least for mesonic 2-point correlators

More studies needed → volumes, lattice spacings

- check if the predicted NLO scaling is verified
- systematic errors under control?

### Conclusions

#### • Other applications:

- Mixed regimes:
  - "heavy" quarks in the *p*-regime, "light" quarks in the  $\epsilon$ -regime

Bernardoni & Hernández ('07), Fukaya & Damgaard ('07)

Spectral density: connect p- and ε-regime

Fukaya & Damgaard ('09), see also Giusti & Lüscher ('09)

 $\rightarrow$  See plenary talk by Hashimoto

Heavy Meson chiral perturbation theory for heavy-light systems

Burdman & Donoghue ('92), Wise ('92)

• LECs of effective weak Hamiltonian  $\rightarrow K \rightarrow \pi\pi$  amplitutes quenched investigations:

Giusti et al ('04)