Status of the Hadronic Light by Light Contribution to Muon $g-2$

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Sixth Workshop on Chiral Dynamics, 9 July 2009, Bern

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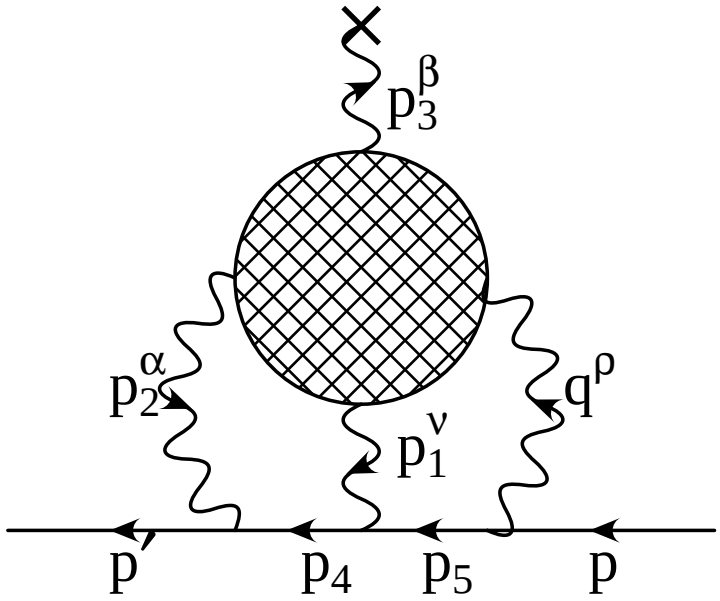
Introduction

Hadronic light-by-light contribution to $g-2$



$$\mathcal{M} = |e|^7 A_\beta \int \frac{\mathrm{d}^4 p_1}{(2\pi)^4} \int \frac{\mathrm{d}^4 p_2}{(2\pi)^4} \frac{1}{q^2\, p_1^2\, p_2^2 (p_4^2 - m^2)\,(p_5^2 - m^2)}$$
$$\times \underline{\Pi^{\rho\nu\alpha\beta}(p_1, p_2, p_3)}\; \bar{u}(p')\gamma_\alpha(\not{p}_4 + m)\gamma_\nu(\not{p}_5 + m)\gamma_\rho u(p)$$

Introduction

Need

$$\Pi^{\rho\nu\alpha\beta}(p_1,p_2,p_3) = i^3 \int \mathrm{d}^4x \int \mathrm{d}^4y \int \mathrm{d}^4z \, \exp^{i(p_1\cdot x + p_2\cdot y + p_3\cdot z)} \times$$
$$\times \langle 0|T[V^\rho(0)V^\nu(x)V^\alpha(y)V^\beta(z)]|0\rangle$$

with $V^\mu(x) = [\overline{q}\,\widehat{Q}\gamma^\mu q](x)$ and $\widehat{Q} = \frac{1}{3}\mathrm{diag}(2,-1,-1)$

full four-point function with $p_3 \longrightarrow 0$ •

Using current conservation at $p_3 = 0$

$$\Pi^{\rho\nu\alpha\lambda}(p_1,p_2,p_3) = -p_{3\beta}\frac{\delta\Pi^{\rho\nu\alpha\beta}(p_1,p_2,p_3)}{\delta p_{3\lambda}}$$

one just needs derivatives at $p_3 = 0$ •

Introduction

★ Many scales involved: Impose low energy and several OPE limits ⇨ Not full first principle calculation at present •

(Two lattice groups just starting: Not clear final uncertainty) •

Large N_c and CHPT counting:

Organizes different degrees of freedom contributions •

E. de Rafael

- Goldstone boson exchange: $\mathcal{O}(N_c)$ and $\mathcal{O}(p^6)$ •
- Quark Loop and non-Goldstone boson exchange: $\mathcal{O}(N_c)$ and $\mathcal{O}(p^8)$ •
- Goldstone bosons Loop: $\mathcal{O}(1)$ in $1/N_c$ and $\mathcal{O}(p^4)$ •

Introduction

Based on this organization:

- Two full calculations

  J. Bijnens, E. Pallante, J.P. (BPP)

  M. Hayakawa, T. Kinoshita, A. Sanda (HKS)

- Dominant pseudo-scalar exchange: Extensive analytic analysis •

  M. Knecht, A. Nyffeler (KN)

  Found sign mistake ✔

  M. Knecht, A. Nyffeler, M. Perrottet, E. de Rafael

Introduction

★ New four-point form factor short-distance constraint:

K. Melnikov, A. Vainshtein

(see also M. Knecht, S. Peris, M. Perrottet, E. de Rafael)

Model:

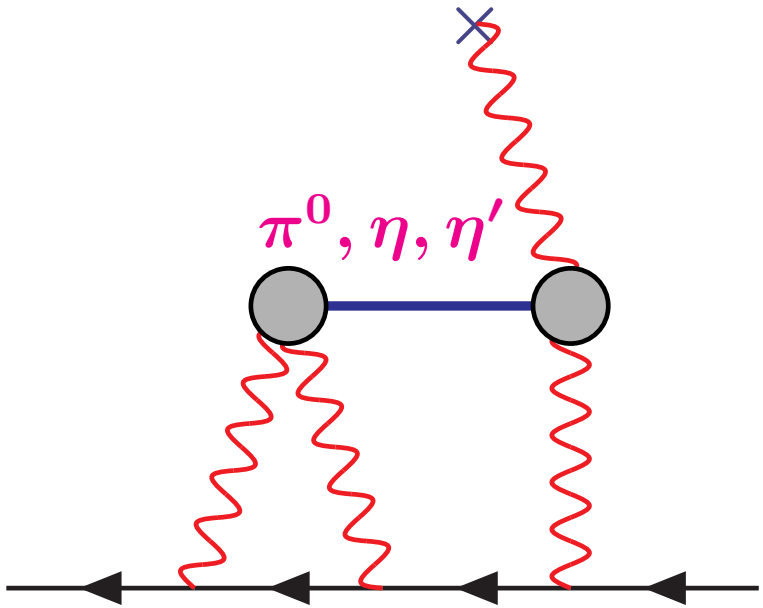
Full light-by-light saturated by pseudo-scalar and pseudo-vector pole exchanges •

Very recently, A. Nyffeler used an $\pi^0\gamma^*\gamma^*$ off-shell form factor model • (see next talk)

First step, one needs more work to have the full light-by-light •

"Old" Calculations: Pseudo-Scalar Exchange

Dominant contribution ⇨ pseudo-scalar exchange •



Nambu-Goldstone π^0 makes special enhacement at low energy

$$a(\pi^0) = \left(\frac{\alpha}{\pi}\right)^3 N_c \frac{m^2 N_c}{48\pi^2 f_\pi^2} \left[\ln^2 \frac{M_\rho}{m_\pi} + \mathcal{O}(\ln \frac{M_\rho}{m_\pi}) + \mathcal{O}(1)\right]$$

M. Knecht, A. Nyffeler, M. Perrottet, E. de Rafael

"Old" Calculations: Pseudo-Scalar Exchange

Here, I discuss work in J. Bijnens, E. Pallante, J.P. •

We used a variety of $\pi^0\gamma^*\gamma^*$ form factors

$$\mathcal{F}^{\mu\nu}(p_1,p_2) = \frac{N_c}{6\pi}\frac{\alpha}{f_\pi}\, i\varepsilon^{\mu\nu\alpha\beta}\, p_{1\alpha}p_{2\beta}\, \underline{\mathcal{F}(p_1^2,p_2^2)}$$

fulfilling as many as possible QCD constraints • (Short-distance, data, $U_A(1)$ normalization and slope at the origin). In particular,

$$\begin{aligned}\mathcal{F}(Q^2,Q^2) &\rightarrow \frac{A}{Q^2}\\ \mathcal{F}(Q^2,0) &\rightarrow \frac{B}{Q^2}\end{aligned}$$

for Q^2 Euclidean and very large

"Old" Calculations: Pseudo-Scalar Exchange

All form factors we used converge for $\mu \sim (2-4)$ GeV and the numerical difference between them is small ✔

Somewhat different $\pi^0\gamma^*\gamma^*$ form factors used in M. Hayakawa, T. Kinoshita, A. Sanda and M. Knecht, A. Nyffeler •

Results agree very well (after correcting a mistake in the sign of the phase space) •

Adding π^0, η and η' contributions

	$10^{10} \times a_\mu$
BPP	(8.5 ± 1.3)
HKS	(8.3 ± 0.6)
KN	(8.3 ± 1.2)

"Old" Calculations: Pseudo-Vector Exchange

Need $a_1^0\gamma\gamma^*$ and $a_1^0\gamma^*\gamma^*$ form factors •

➪ related to $\pi^0\gamma\gamma^*$ and $\pi^0\gamma^*\gamma^*$ by anomalous Ward identities ✔

Pseudo-vector exchange

	$10^{10} \times a_\mu$
BPP	(0.25 ± 0.10)
HKS	(0.17 ± 0.10)

"Old" Calculations: Scalar Exchange

Need $S^0\gamma\gamma^*$ and $S^0\gamma^*\gamma^*$ form factors •

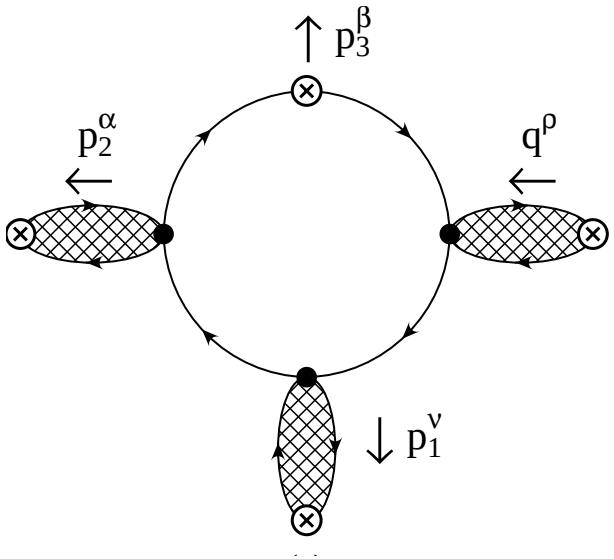
They are constrained by CHPT at $\mathcal{O}(p^4)$: L_i's reproduced ✔

Within ENJL: Ward identities impose relations between Quark loop and Scalar exchange •

$$a_\mu(\text{Scalar}) = -(0.7 \pm 0.2) \cdot 10^{-10}$$

Not included by M. Hayakawa, T. Kinoshita and A. Sanda nor by K. Melnikov and A. Vainshtein •

"Old" Calculations: Non-Meson Exchange "Quark-Loop"

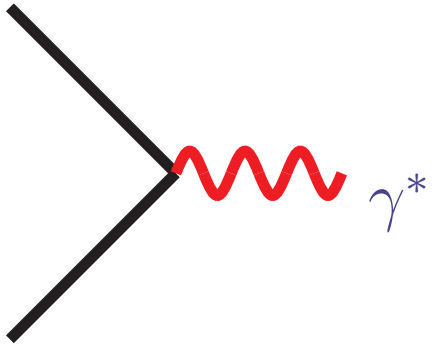


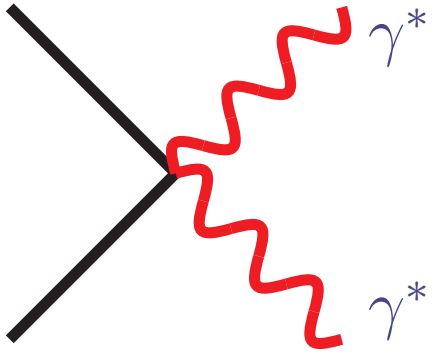
Λ [GeV]	$10^{10} \times a_\mu$
0.7	2.2
1.0	2.0
2.0	1.9
4.0	2.0

- Low Energy (0 to Λ): ENJL model •
- High Energy (Λ to ∞): Bare heavy quark loop with $m_Q = \Lambda$ •
- Numerical matching ✔

"Old" Calculations: Pion- and Kaon-Loop

Leading contribution in chiral counting, suppressed by $1/N_c$

OK: 

Problem: 

Low-intermediate energy (0.5-1.5) GeV $\gamma^*\gamma^* \to \pi\pi$ data not available: Models needed !

Model for $\pi\pi\gamma(\gamma)$ form-factor	$10^{10} \times a_\mu$
Point-like	-4.6
BPP (Full VMD)	-1.8
HKS (HGS)	-0.4

Kaon loop is much smaller: -0.05×10^{-10} •

New Short Distance Constraints

K. Melnikov and A. Vainshtein

New short-distance constraint on four-point function form factor

$$\langle 0|T[V^\nu(p_1)V^\alpha(p_2)V^\rho(-(p_1+p_2+p_3))]|\gamma(p_3\to 0)\rangle$$

using OPE with $-p_1^2 \simeq -p_2^2 >> -(p_1+p_2)^2$ Euclidean and large,

$$T[V^\nu(p_1)V^\alpha(p_2)] \sim \frac{1}{\hat{p}^2}\,\varepsilon^{\nu\alpha\mu\beta}\,\hat{p}_\mu\,[\overline{q}\,\widehat{Q}^2\gamma_\beta\gamma_5 q](p_1+p_2)$$

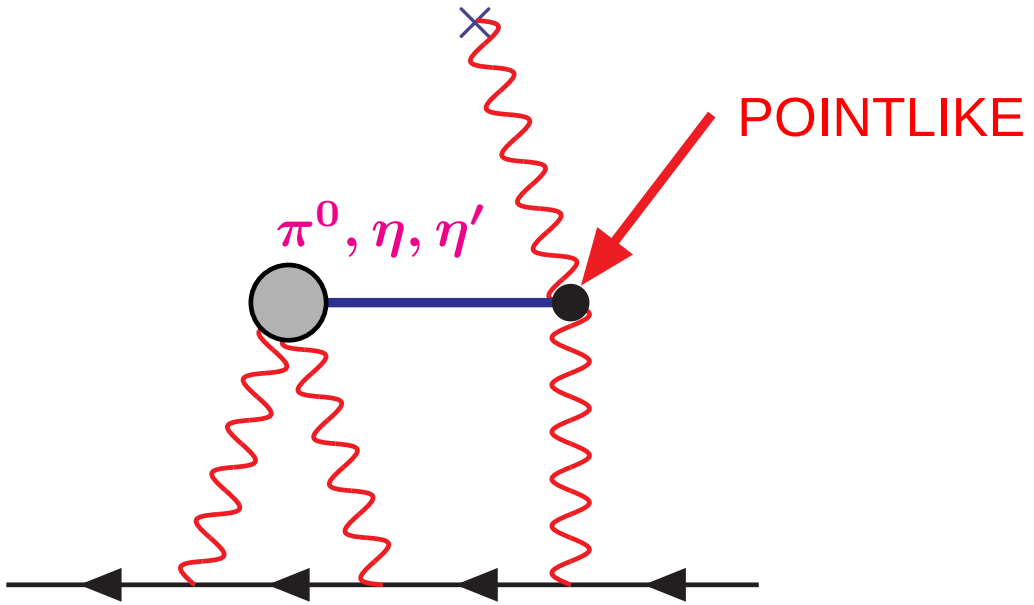
with $\hat{p} = (p_1-p_2)/2 \simeq p_1 \simeq -p_2$

New OPE Constraint: Pseudo-scalar exchange

New OPE constraint saturated by pseudo-scalar exchange

➪ Model uses a point-like vertex when $p_3 \to 0$ •

Not all OPE constraints satisfied: Negligible numerically •



New OPE Constraint: Axial-Vector exchange

Axial-Vector exchange depends very much on the resonance mass mixing •

K. Melnikov and A. Vainsthein:
Ideal mixing for $f_1(1285)$ and $f_1(1420)$ •

Mass mixing	$10^{10} \times a_\mu$
No New OPE (Nonet symmetry)	0.3 ± 0.1
M=1.3 GeV (Nonet symmetry)	0.7
M=M_ρ (Nonet symmetry)	2.8
Ideal mixing	2.2 ± 0.5

Comparison: Leading order in N_c

Leading order in N_c:

Quark Loop + Pseudo-Scalar + Pseudo-Vector + Scalar Exchanges •

Total at $\mathcal{O}(N_c)$	$10^{10} \times a_\mu$
BPP (Nonet symmetry)	(10.9 $\pm$ 1.9) + -(0.7 $\pm$ 0.1) = (10.2 $\pm$ 1.9)
HKS (Nonet symmetry)	(9.4 $\pm$ 1.6) + ??Scalar??

MV: Hadronic model saturated by pole exchanges:
Cannot compare individual contributions •

Total at $\mathcal{O}(N_c)$	$10^{10} \times a_\mu$
MV (Nonet symmetry)	(12.1 $\pm$ 1.0) + ??Scalar??
MV (Ideal mass mixing)	(13.6 $\pm$ 1.5) + ??Scalar??

Masses produce main difference in pseudo-vector exchange •

Comparison: NLO in $1/N_c$

Next to leading order in $1/N_c$ contributions:

Charged Pion and Kaon Loop •

Model for $\pi\pi\gamma(\gamma)$ form-factor	$10^{10} \times a_\mu$
BPP (Full VMD)	-1.9 ± 0.5
HKS (HGS)	-0.45 ± 0.8

K. Melnikov and A. Vainshtein:

Full NLO in $1/N_c$ estimate

$$a_\mu = (0 \pm 1) \cdot 10^{-10}$$

Comparison

BPP vs HKS:

Full	$10^{10} \times a_\mu$
BPP	8.3 ± 3.2
HKS	8.9 ± 1.7

No scalar exchange, different quark loop and different pion and kaon loops almost compensate •

Comparison

BPP vs MV:

Full	$10^{10} \times a_\mu$
BPP	8.3 ± 3.2
MV	13.6 ± 2.5

Several order $1.5 \cdot 10^{-10}$ differences,
in addition to new OPE effects •

$-1.5 \cdot 10^{-10}$ (Different pseudo-vector mass mixing)
$-0.7 \cdot 10^{-10}$ (No scalar exchange)
$-1.9 \cdot 10^{-10}$ (No pion+kaon loop)
= $-4.1 \cdot 10^{-10}$ •

Final [BPP-MV] difference: $-5.3 \cdot 10^{-10}$ •

Conclusions and Prospects

At present, large N_c results agree within 1 σ ✔

$$a_\mu^{N_c,\mathrm{lbl}} = (11.0 \pm 4.0) \times 10^{-10}$$

Based partly in this discussion: Recent new analysis of HLbL

J.P., E. de Rafael and A. Vainshtein

★ $1/N_c$ expansion works reasonably well ✔

★ Chiral enhancement factors demand more than the lightest Nambu-Golsdtone bosons •

Conclusions Prospects

Adding effects beyond leading order in $1/N_c$, in a conservative analysis, J.P., E. de Rafael and A. Vainshtein

π^0, η and η' exchanges: $(11.4 \pm 1.3) \times 10^{-10}$
(includes non-meson exchange in ENJL case)

Scalar exchange : $-(0.7 \pm 0.7) \times 10^{-10}$

Axial-vector exchange : $(1.5 \pm 1.0) \times 10^{-10}$

Pion and kaon loops: $-(1.9 \pm 1.9) \times 10^{-10}$

Charm quark loop: 0.23×10^{-10}

★ Our final result is $a_\mu^{\rm lbl} = (10.5 \pm 2.6) \times 10^{-10}$

Conclusions and Prospects

New muon $g-2$ experiment goal ➪ 1.6×10^{-10}

A new full calculation of $a_\mu^{\rm lbl}$ is desirable and possible: Goal: $\sim (1.5-2.0) \times 10^{-10}$ uncertainty in $a_\mu^{\rm lbl}$ •

★ At $\mathcal{O}(N_c)$: Study relevant reduced full four-point function with large N_c techniques •

☛ First step: Analytic expressions •

☛ Second step: Implement as many short-distance and low energy constraints as possible •

(possible problems J. Bijnens, E. Gámiz, E. Lipartia, J.P.)

Conclusions and Prospects

★ NLO in $1/N_c$: More theoretical work is clearly needed here !

⇨ Goldstone bosons at one loop: Need $\gamma^*\gamma^* \to \pi^+\pi^-$ vertex •
Experiment can help here •

⇨ Non-Goldstone bosons at one loop: Little is known •
(recent related work by A. Pich, I. Rosell, J.J. Sanz-Cillero) •

More work is certainly needed •

Conclusions and Prospects

More work needed to have the hadronic light-by-light contribution to muon $g-2$ with new experimental uncertainty goal: 1.6×10^{-10} •

Goal: To have under control model dependences •

Corolary: HLbL to electron g-2

★ New very accurate a_e leads to most precise fine structure constant α •

G. Gabrielse

New analysis, including subleading m_l^2 terms in π^0 exchange at leading logs approx togethe with m_e^2/m_μ^2 scaling leads to:

$$a_e^{\rm lbl} = (0.35 \pm 0.10) \times 10^{-13}$$

J.P., E. de Rafael and A. Vainshtein