

# Recent results on strange systems from QCD sum rules and lattice QCD

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Chiral 09  
06/07/09

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with P. Gubler, D. Jido, T. Kojo, T. Nishikawa

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QCD Sum Rules: T. Doi, H. Kim, Y. Kondo, G. Erkol

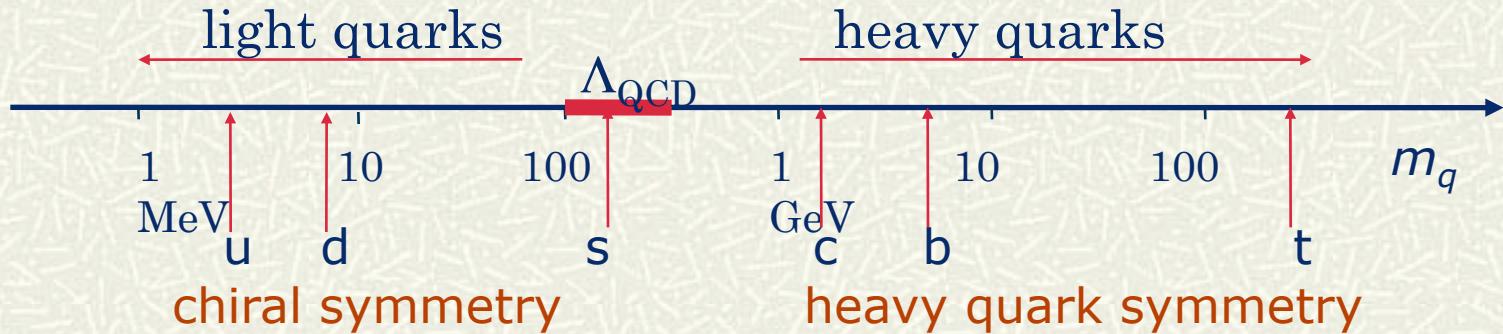
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# Strangeness in QCD

$$\mathcal{L}_{QCD} = -\frac{1}{4}(G_{\mu\nu}^a)^2 + \sum_q \bar{q}(i D_\mu \gamma^\mu - m_q)q ,$$

If the quark masses of  $N_f$  flavors are equal, QCD has the exact flavor  $SU(N_f)$  symmetry.



The dynamics of both the light and heavy quarks can be described by effective theories, but the **strange** quark is sensitive to the QCD details.

# Strangeness in QCD

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*The quark model* with  $SU(3)/SU(6)$  flavor symmetry has played a major role in the history of hadron physics, but its limitation has become clear in the spectroscopy of excited and exotic hadrons.

Technical developments in the direct QCD approaches start to allow reliable predictions to the hadronic observables from QCD.

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# Evidence of the $\Theta^+$ in the $\gamma d \rightarrow K^+ K^- pn$ reaction

T. Nakano,<sup>1</sup> N. Muramatsu,<sup>1</sup> D.S. Ahn,<sup>1</sup> J.K. Ahn,<sup>2</sup> H. Akimune,<sup>3</sup> Y. Asano,<sup>4</sup>

PRC79 (2009) 025210 arXiv:0812.1035 (4 Dec 2008)

$m_\Theta = 1.524(2)(3) \text{ GeV}/c^2$  with statistics  $5.1 \sigma$

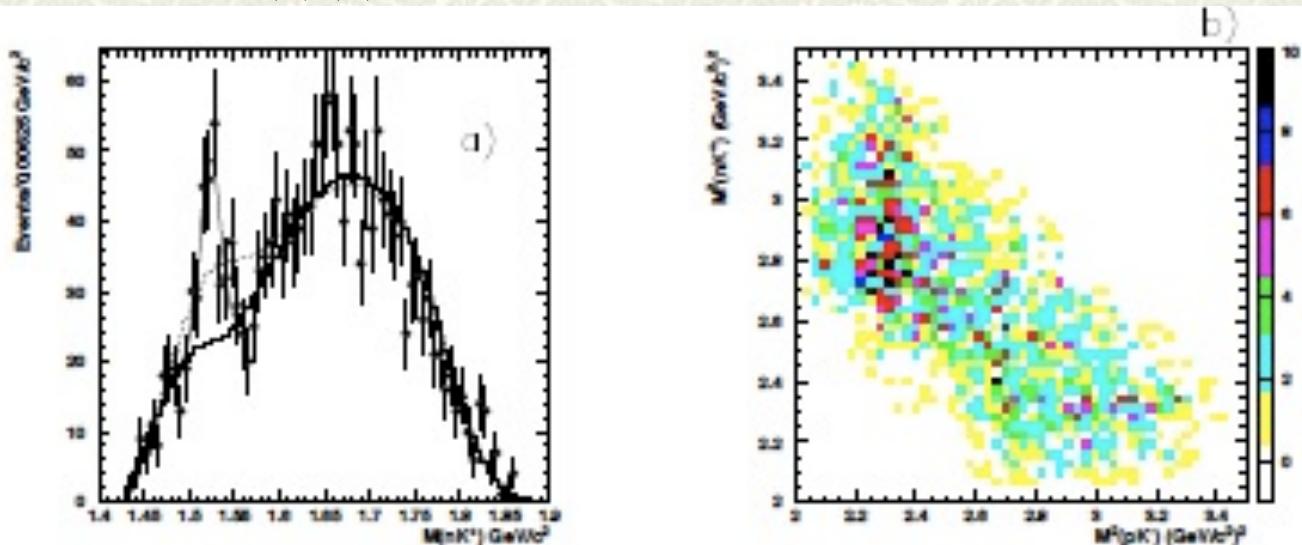


FIG. 12: (a)  $M(nK^+)$  distribution with a fit to the RMM background spectrum only (dashed line) and with a Gaussian function (solid line). The dotted line is the background. (b) Dalitz plot of  $M^2(nK^+)$  vs.  $M^2(pK^-)$ .

# Lattice QCD for $\Theta^+$

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## # Lattice QCD for $\Theta^+$ (so far *quenched*)

F. Csikor, S. Sasaki, T.-W. Chiu, N. Mathur

N. Ishii, T. Takahashi, B. Lasscock

- No low-lying  $1/2^+ (I=0)$  resonance state.
- May exists a  $1/2^- (I=0)$  resonance near the NK threshold, but it requires careful study to distinguish from NK scattering states.
- Some contradictory results for the  $3/2^+$  or  $3/2^-$  states.

## # Anticipate full-QCD calculations with physical quark masses in which scattering states and resonances are clearly distinguished.

# QCD SR for $\Theta^+$

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P. Gubler, D. Jido, T. Kojo, T. Nishikawa, M.O., PRD 79  
(2009) 114011, arXiv:0902.2049

- # QCD SR predicts  $(J, I) = (3/2^+, 0)$  and  $(3/2^+, 1)$  penta-quarks in the mass region of  $1.4\text{-}1.6 \text{ GeV}/c^2$ .
- # They decay into NK P-wave states. The decay widths can be relatively small. No QCD prediction of the decay width is available.

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# Nuclear force

## # Three Regions of Nuclear Force

**N.F. = OPE + Med. Attr. + SR Repulsion**

**Energy scales  $\sim 100 \sim 500$  MeV**

$\Leftrightarrow$  **B.E. of the deuteron  $\sim 2$  MeV**

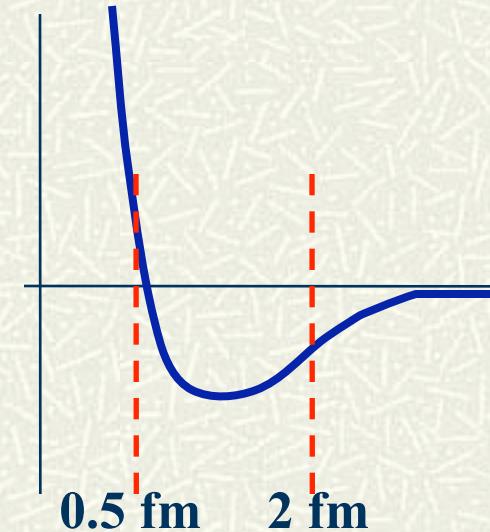
**Range of the OBE  $\sim 1$  fm**

$\Leftrightarrow$  **Size of the deuteron  $\sim 4$  fm**

## # Are the other baryonic potential similar?

**Use of SU(3) symmetry to generalize meson exchanges**

**How can we extrapolate the SR part? Its origin?**



# Generalized Nuclear Force

$$\text{SU}(6) \supset \text{SU}(3) \text{ flavor} \times \text{SU}(2) \text{ spin}$$

$$[3] \underline{56} \leftarrow \begin{array}{l} 8 (\text{S}=1/2): \text{ N } \Lambda \Sigma \Xi \\ 10 (\text{S}=3/2): \Delta \Sigma^* \Xi^* \Omega \end{array}$$

## SU(3) flavor

$$8 \times 8 = 1 + 8_S + 27 + 8_A + 10 + 10^*$$

↑ Symmetric      ↑ Antisymm      ↑  
ΛΛ - NΞ - ΣΣ (I=0)      NN (I=1)      NN (I=0)

M. O., K. Shimizu, K. Yazaki, PLB130 (1983) 365, NPA464 (1987) 700

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# Generalized Nuclear Force

27

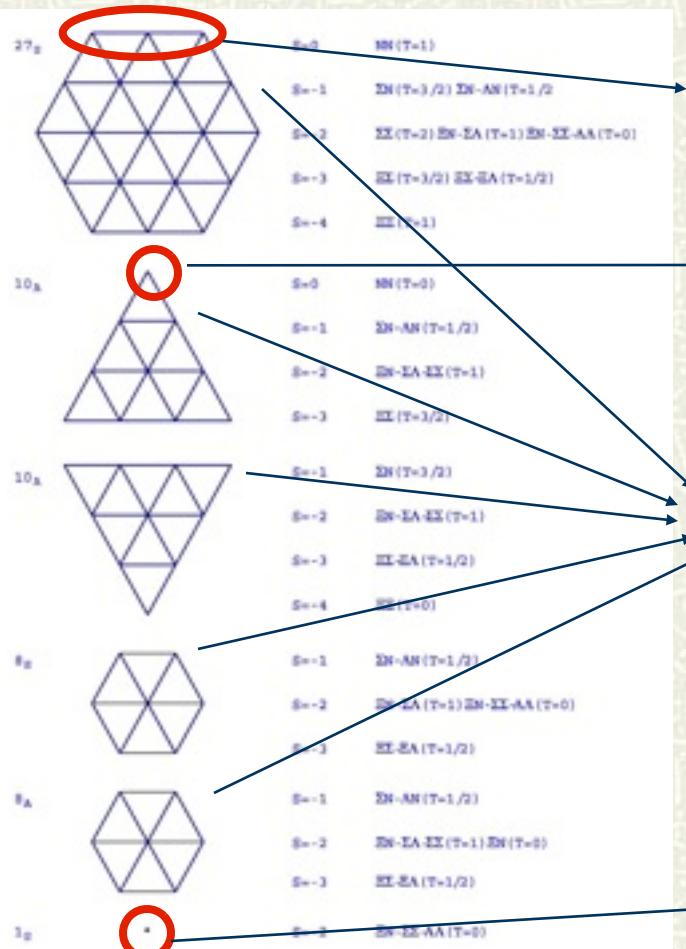
$10^*$

10

$8_S$

$8_A$

1



$NN(I=1) \ ^3S_1$

$NN(I=0) \ ^1S_0$

$N\Lambda, N\Sigma$

*Hypernucleus*

H dibaryon  $\Lambda\Lambda$  -  $N\Xi$  -  $\Sigma\Sigma$  ( $I=0$ )  
6-quark states?

# Generalized Nuclear Force

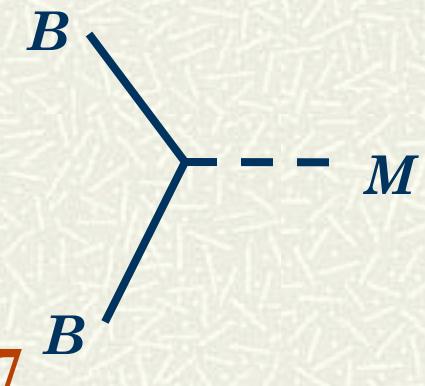
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- # Understanding generalized nuclear force from the first principle is a key to disentangle origins and mechanisms of hadronic interactions.
- # As the first step, we explore the SU(3) structure of the meson exchange part of the potential.

# SU(3) Symmetry of M-B-B vertex

Baryon octet 8  
Meson octet 8  
singlet 1

N Λ Σ Ε  
 $\pi$   $\eta_8$  K  
 $\eta_1$



$$8 \times 8 = 1 + 8_D + 8_F + 10 + 10 + 27$$

Two independent couplings for  $8 \times 8 = 8$

$$\begin{aligned} F \text{ Tr}[ [ B, B ] M ] \\ D \text{ Tr}[ \{ B, B \} M ] \end{aligned}$$

F coupling  
D coupling

$\alpha_F = F / (F+D)$  is a free parameter in SU(3)

# Spin-Flavor SU(6) Symmetry

Based on the NR quark model ( $u\uparrow u\downarrow d\uparrow d\downarrow s\uparrow s\downarrow$ )

Baryons  $56 = (8, 1/2) + (10, 3/2)$   
 $N \Lambda \Sigma \Xi \Delta \Sigma^* \Xi^* \Omega$

Mesons  $1 = (1,0) \quad \eta_1$   
 $35 = (8,0) + (8,1) + (1,1)$   
 $\pi \eta_8 K \rho \omega K^* \Phi$

$B \quad B \quad M$   
 $56 \times 56 = 1 + 35 + 405 + 2695$   
(8,0) (unique)      (8,0)

F/D ratio is fixed by the SU(6) symmetry  
 $\rightarrow F/D = 2/3$  or  $a_F = 2/5$

# Spin-Flavor SU(6) Symmetry

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But,

QCD has no reason to adopt the SU(6) symmetry.

The SU(6) symmetry is **broken!!**

by mass difference of quarks

by spin-dependent interactions of quarks

ex N (J=1/2) – Δ (J=3/2) mass difference

(8,1/2)      (10, 3/2)    56-dim. SU(6) rep.

# QCD calculation of coupling constants

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- # The SU(3) invariance for the coupling constants is not established, although the phenomenological models often assume the invariance. The F/D ratios of the coupling constants are the fitting parameters in the models.
- # How good is the SU(3) symmetry in the coupling constants?
- # What does QCD predict for F/D ratio, if SU(3) is valid?
- # Direct computation of the coupling constants from QCD QCD sum rule T. Doi, H. Kim, Y. Kondo, G. Erkol, M. O. Lattice QCD T.T. Takahashi, G. Erkol, M.O.

# QCD Sum rule for coupling constants

T. Doi, Y. Kondo, M.O., Phys. Rept. 398 (2004) 253

## # $\pi NN$ coupling constant

$$\Pi^{\alpha\beta}(q, p) = i \int d^4x e^{iq \cdot x} \langle 0 | T \left[ J_N^\alpha(x) \bar{J}_N^\beta(0) \right] | \pi(p) \rangle$$

$$\Pi(q, p) = i\gamma_5 \not{p} \Pi^{\text{PV}} + i\gamma_5 \Pi^{\text{PS}} + \boxed{\gamma_5 \sigma^{\mu\nu} q_\mu p_\nu \Pi^{\text{T}}} + i\gamma_5 \not{q} \tilde{\Pi}^{\text{PV}}$$

tensor structure

## # nucleon interpolation field

$$J_N(x; t) = 2\epsilon_{abc} [ (u_a^T(x) C d_b(x)) \gamma_5 u_c(x) + t (u_a^T(x) C \gamma_5 d_b(x)) u_c(x) ]$$

two independent terms mixed by  $\tan \theta = t$

# F/D ratio *v.s.* $\cos\theta$ for T sum rule

## Tensor sum rule

T. Doi, H. Kim, M.O., Phys.Rev. C62 (2000) 055202

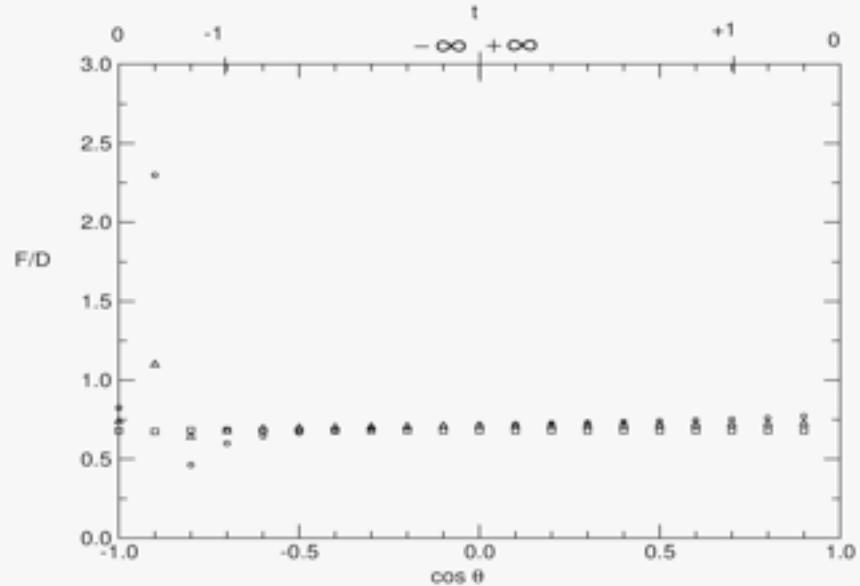
$$F/D = 0.65 \pm 0.10$$

$$= 2/3 \text{ for SU(6)}$$

$$= 0.57 \text{ from } g_A \text{ (exp)}$$

$$\alpha_F = F/(F+D) = 0.394$$

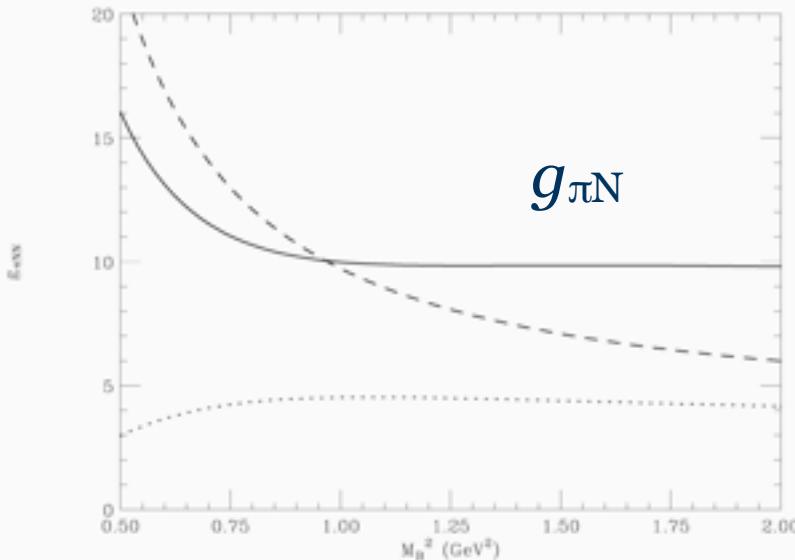
$$= 2/5 \text{ for SU(6)}$$



The SU(3) is fairly good for  $\pi BB'$  coupling consistent with the phenomenological models, ex. Nijmegen potential.

# Projected correlation function

- # The most reliable estimate of the absolute value of the pi-N-N coupling is by the projected correlated function method: Kondo-Morimatsu Nucl. Phys. A717 (2003)



$$g_{\pi N} = \mathbf{9.6 \pm 1.6}$$

v.s.

$$g_{\pi N} (\text{exp.}) \sim \mathbf{12.8}$$

**slightly  
underestimated**

# QCDSR summary

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## # Other mesons

- K and  $\eta$ :  $p^2$  expansion is marginally justifiable.
- scalar mesons:  $\sigma, a_0, f_0, \kappa$   
G. Erkol, M. O., T. Rijken, R. Timmermans,  
PRC73 (2006) 044009  
 $\alpha_F \sim 0.55$ , but observe significant SU(3) breaking
  - The  $\Sigma$ - $\sigma$  and  $\Xi$ - $\sigma$  couplings are enhanced by 30-50%.
- The baryon-scalar-meson coupling strengths can be a good indicator of the 4-quark structure of the scalar mesons.

## # A challenging subject for LQCD.

# Lattice QCD calculation of the ps-meson-octet-baryon couplings

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- # **Unquenched lattice QCD is applied to the pseudoscalar-meson-octet-baryon coupling form factors.**

T. T. Takahashi, G. Erkol, MO PRD 79 (2009) 074509

- CP-PACS gauge configuration: 2-flavor dynamical quarks on the  $16^3 \times 32$  lattice
- RG improved gauge action + the mean-field improved clover quark action
- $\beta = 1.95 \rightarrow a = 0.16 \text{ fm} \quad a^{-1} = 1.267 \text{ GeV}$
- The ratio and absolute values of the coupling constants are obtained for several quark masses:  
 $m_q \sim 150, 100, 65, 35 \text{ MeV}$

# Lattice QCD calculation

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## # 3-point function

$$\begin{aligned} \langle G^{\mathcal{B}\mathcal{P}\mathcal{B}'}(t_2, t_1; \mathbf{p}', \mathbf{p}; \Gamma) \rangle &= -i \sum_{\mathbf{x}_2, \mathbf{x}_1} e^{-i\mathbf{p} \cdot \mathbf{x}_2} e^{i\mathbf{q} \cdot \mathbf{x}_1} \\ &\times \Gamma^{\alpha\alpha'} \langle \text{vac} | T[\eta_{\mathcal{B}}^\alpha(x_2) P(x_1) \bar{\eta}_{\mathcal{B}'}^{\alpha'}(0)] | \text{vac} \rangle, \end{aligned}$$

$$\eta_N(x) = \epsilon^{abc} [u^{Ta}(x) C \gamma_5 d^b(x)] u^c(x),$$

$$\eta_\Sigma(x) = \epsilon^{abc} [s^{Ta}(x) C \gamma_5 u^b(x)] u^c(x),$$

$$\begin{aligned} \eta_\Lambda(x) &= \frac{1}{\sqrt{6}} \epsilon^{abc} \{ [u^{Ta}(x) C \gamma_5 s^b(x)] d^c(x) - [d^{Ta}(x) C \\ &\quad \times \gamma_5 s^b(x)] u^c(x) + 2[u^{Ta}(x) C \gamma_5 d^b(x)] s^c(x) \}, \end{aligned}$$

# Lattice QCD calculation

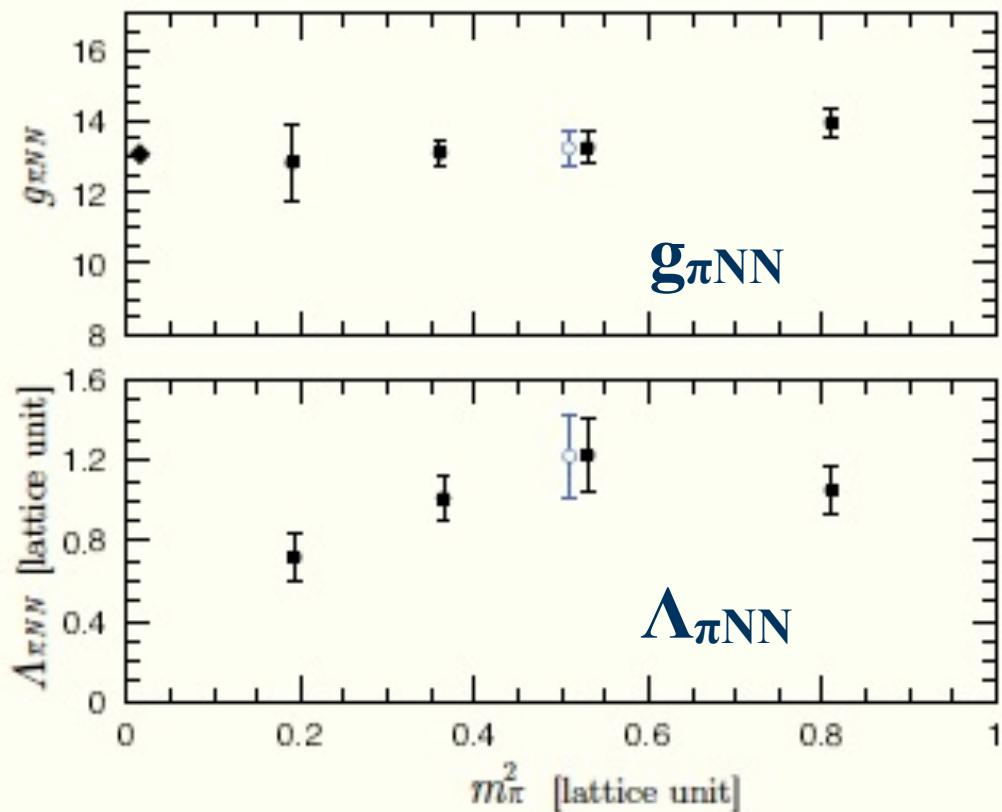
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## # PS-meson-baryon coupling form factors

$$R(t_2, t_1; \mathbf{p}', \mathbf{p}; \Gamma; \mu) = \frac{\langle G^{\mathcal{B}P\mathcal{B}'}(t_2, t_1; \mathbf{p}', \mathbf{p}; \Gamma) \rangle}{\langle G^{\mathcal{B}'}(t_2; \mathbf{p}'; \Gamma_4) \rangle} \left[ \frac{\langle G^{\mathcal{B}}(t_2 - t_1; \mathbf{p}; \Gamma_4) \rangle}{\langle G^{\mathcal{B}'}(t_2 - t_1; \mathbf{p}'; \Gamma_4) \rangle} \right. \\ \times \left. \frac{\langle G^{\mathcal{B}'}(t_1; \mathbf{p}'; \Gamma_4) \rangle \langle G^{\mathcal{B}'}(t_2; \mathbf{p}'; \Gamma_4) \rangle}{\langle G^{\mathcal{B}}(t_1; \mathbf{p}; \Gamma_4) \rangle \langle G^{\mathcal{B}}(t_2; \mathbf{p}; \Gamma_4) \rangle} \right]^{1/2},$$

$$R(t_2, t_1; \mathbf{0}, \mathbf{p}; \Gamma; \mu) \xrightarrow[t_2 - t_1 \gg a]{t_1 \gg a} \frac{g_P^L(q^2)}{[2E(E + m)]^{1/2}} q_3,$$

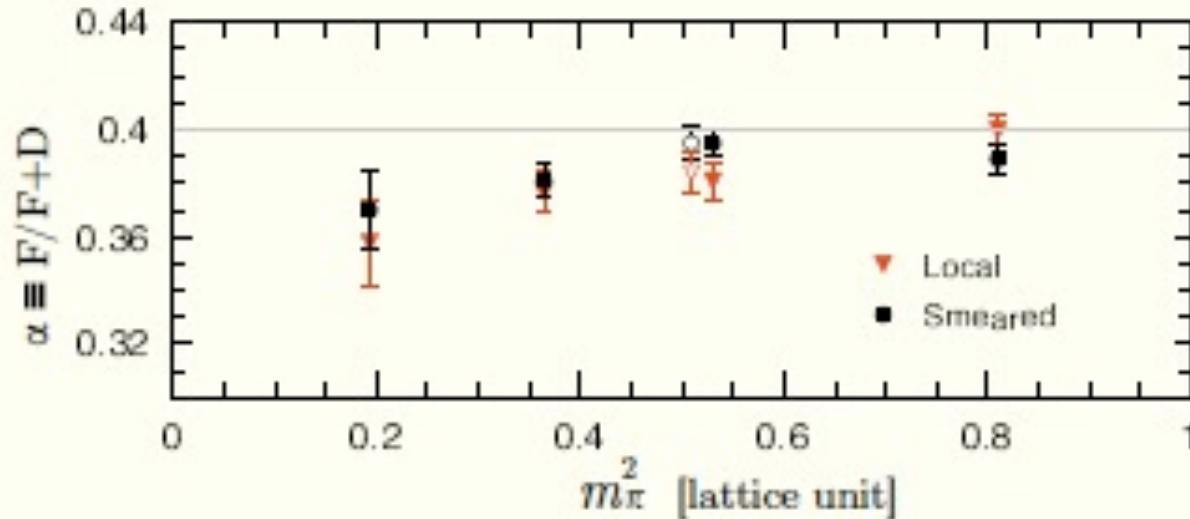
# pi-N-N coupling/form factor



- $g_{\pi NN}$  is consistent with  $g_{\pi NN}$  (pheno.)  $\sim 12.8$
- The monopole form factor is softer than the one used in the meson exchange models.

$$\begin{aligned}\Lambda_{\pi NN} &\sim 0.6 \text{ a}^{-1} \\ &\sim 0.75 \text{ GeV}\end{aligned}$$

# SU(3) couplings



$$\alpha = \frac{F}{F + D} = 0.395 \quad v.s. \text{ SU}(6) \text{ value } \alpha = 0.4$$

**SU(3) breaking effect is very small. The deviation  $\delta < 5\%$**

# Lattice QCD Summary

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- # The two-flavor full-QCD lattice calculation was performed for the ps meson-baryon coupling constants and form factors.
- #  $g_{\pi NN}$  is consistent with  $g_{\pi NN}$  (pheno.)  $\sim 12.8$ .
- # The monopole form factor is softer than the one used in the meson exchange models.  $\Lambda_{\pi NN} \sim 0.75$  GeV.
- # The SU(3) symmetry for the ps meson-baryon couplings happens to be “very” good.
- # F/D ratio  $a_F=0395$  is close to the SU(6) value.

# Conclusion and Outlook

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- # Serious QCD calculations are inevitable for revealing new (exotic) multi-quark world of hadrons.
- # Lattice QCD has reached the position in which the real quark-mass full QCD with chiral symmetry is accessible. (but not yet applied to pentaquarks)
- # QCD sum rules also give reliable results with careful evaluations of its validity.
- # Strange systems supply many ideal subjects for checking physical ideas of hadron dynamics.
- # JPARC is a dream machine for strangeness hadron physics.