

Single Nucleon experiments

A. Deur

Thomas Jefferson National Accelerator Facility

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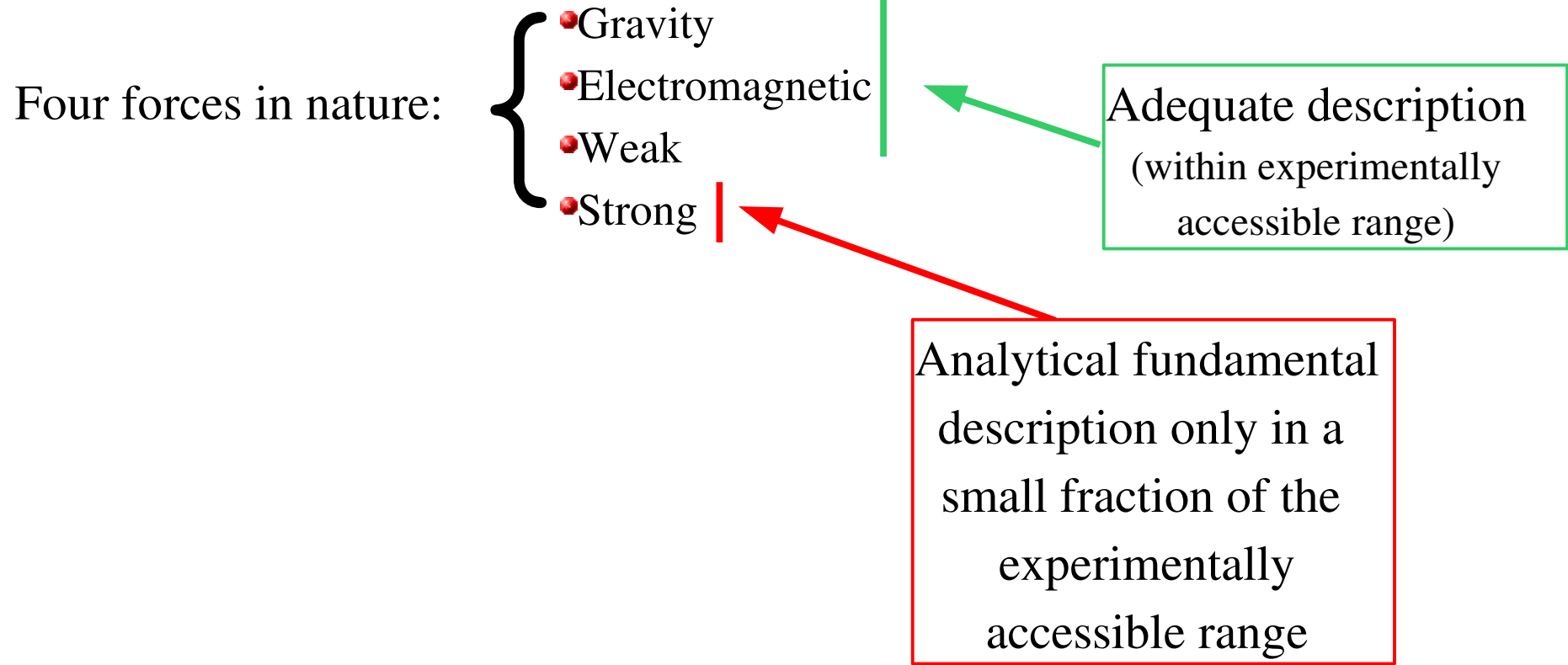
- Single nucleon (neutron or proton)
- Single particle detection (lepton detection, inclusive experiments)
- Single spin direction for each type of particle (doubly polarized experiments)
- Single Lab (focus on the role of Jefferson Lab)

Context (why should we care about the nucleon spin structure at low Q^2 ?)

Four forces in nature: {

- Gravity
- Electromagnetic
- Weak
- Strong

Context (why should we care about the nucleon spin structure at low Q^2 ?)



We need to know the strong force better for a complete understanding of nature.

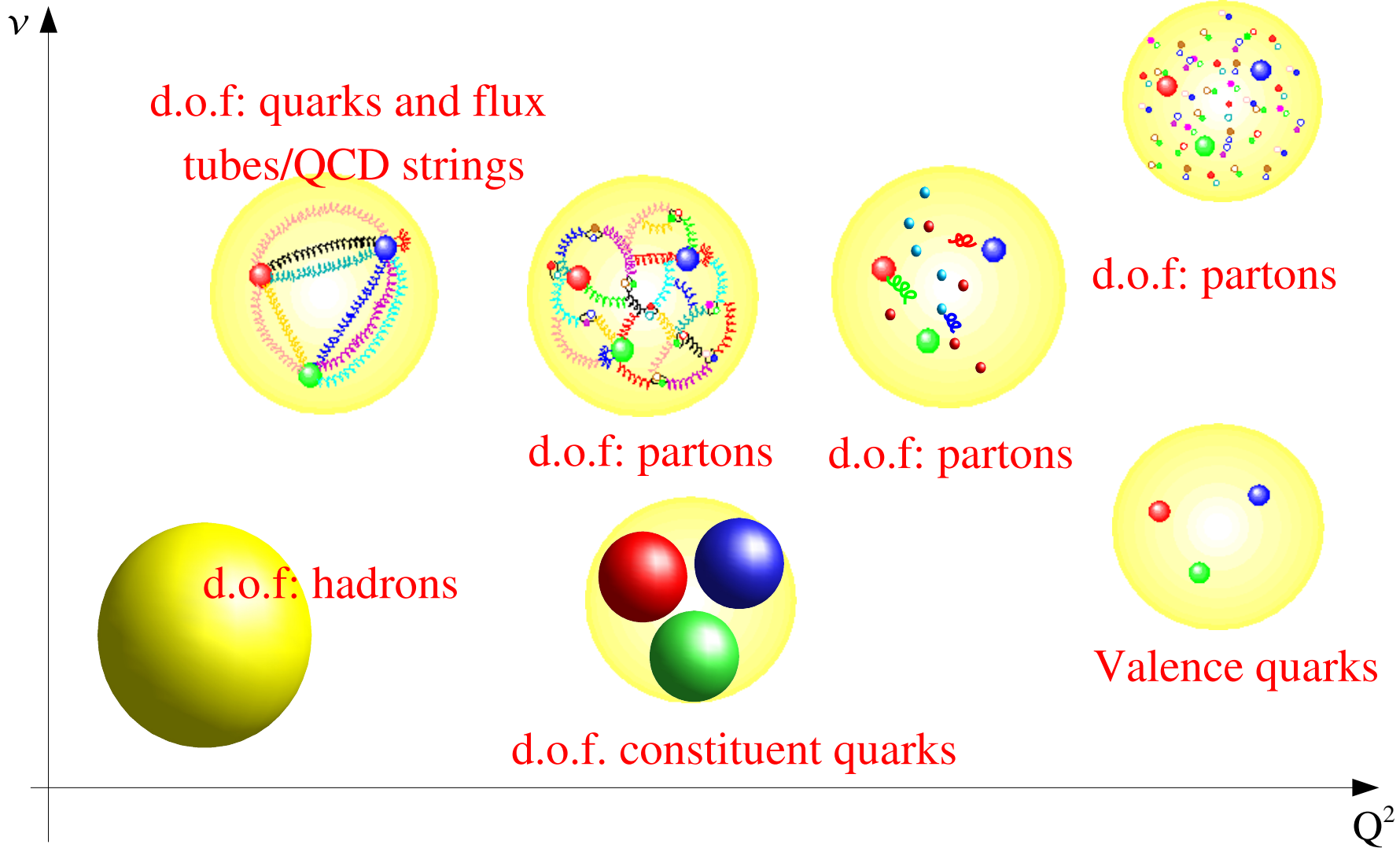
The nucleus is a natural laboratory for such study because its structure is governed by the strong force.

1st step:

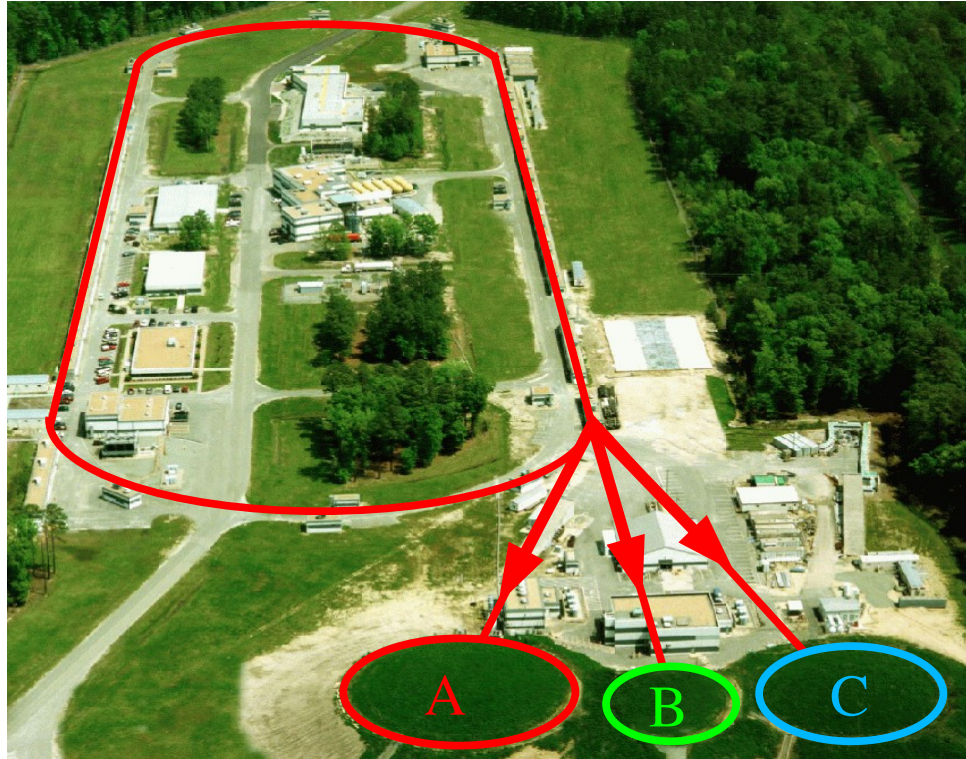
- Test:
- Gauge theory of strong force (QCD) in calculable domain (perturbative domain, pQCD).
 - Effective descriptions in non perturbative domain (e.g. χ pT).

2nd step:

Connection between fundamental description (QCD)
and effective descriptions.

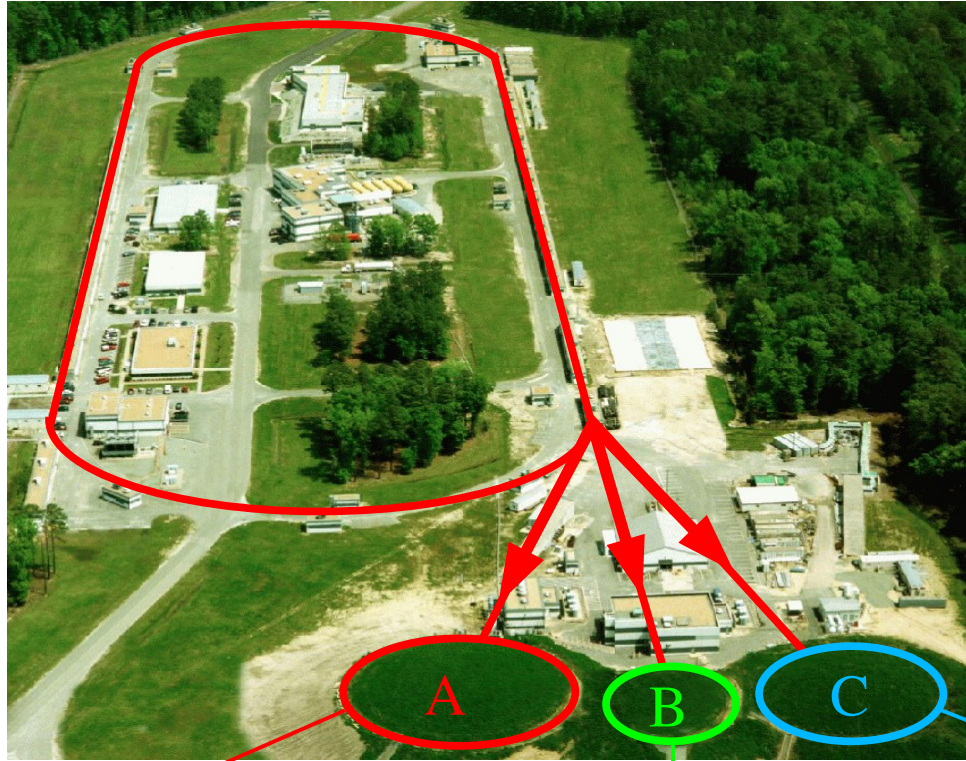


Jefferson Lab

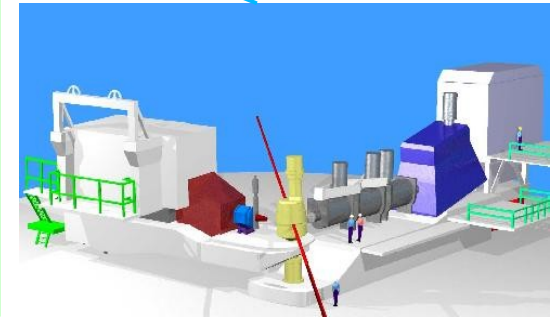
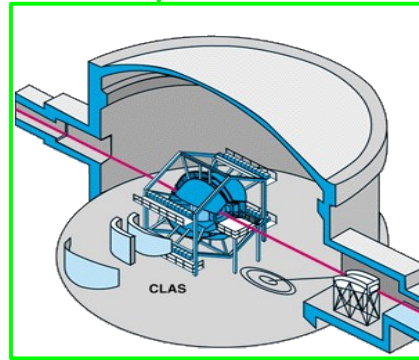
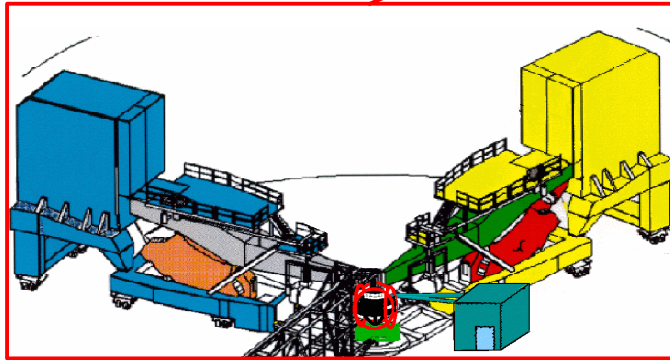


Continuous e^- beam.
1 to 6 GeV.
Polarization: $\sim 85\%$
Up to $200 \mu\text{A}$.

Jefferson Lab



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Up to $200 \mu\text{A}$.

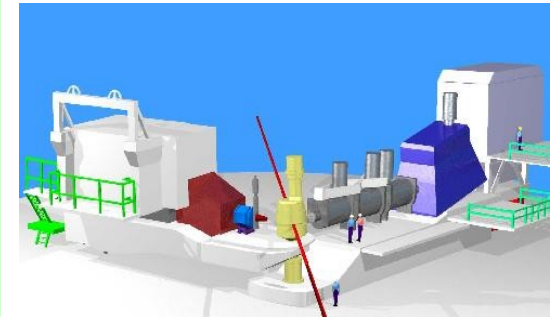
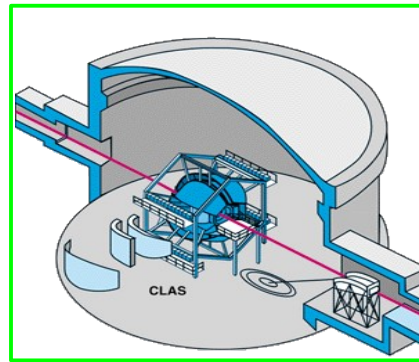
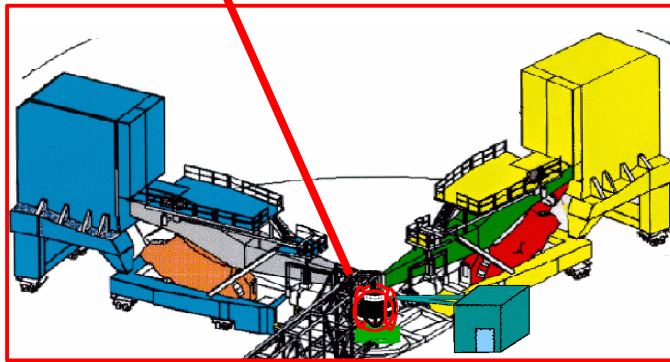
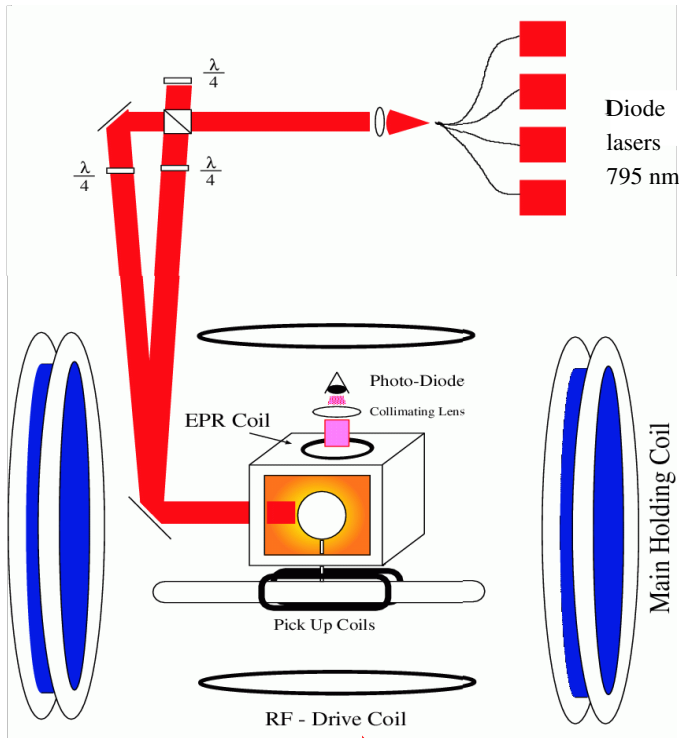


Polarized targets

^3He target:

- Effective polarized neutron target
- High luminosity: $10^{36} \text{ s}^{-1} \text{ cm}^{-2}$
- Low dilution: $\sim 30\%$
- Excellent polarization: $\sim 60-70\%$
- Any polarization directions

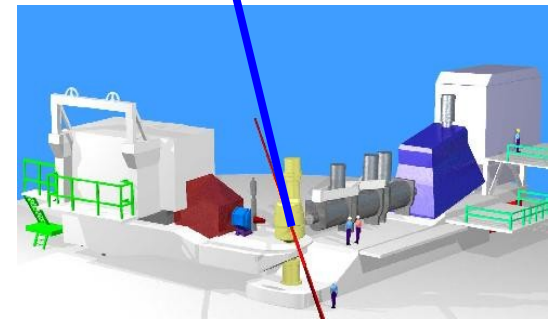
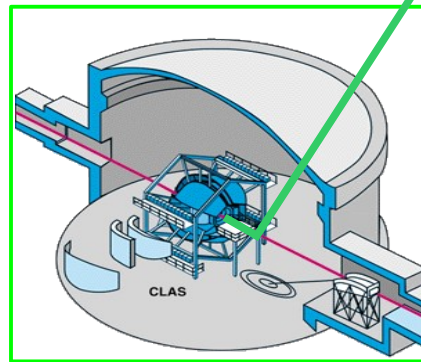
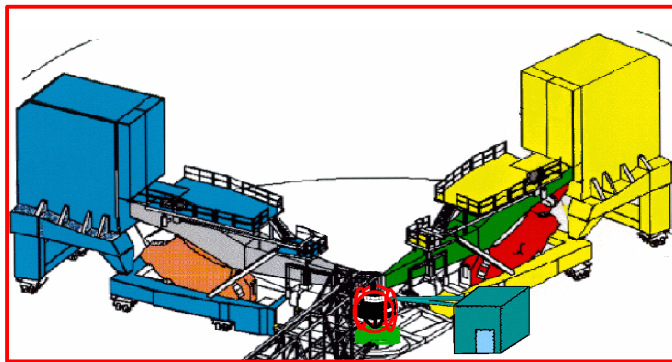
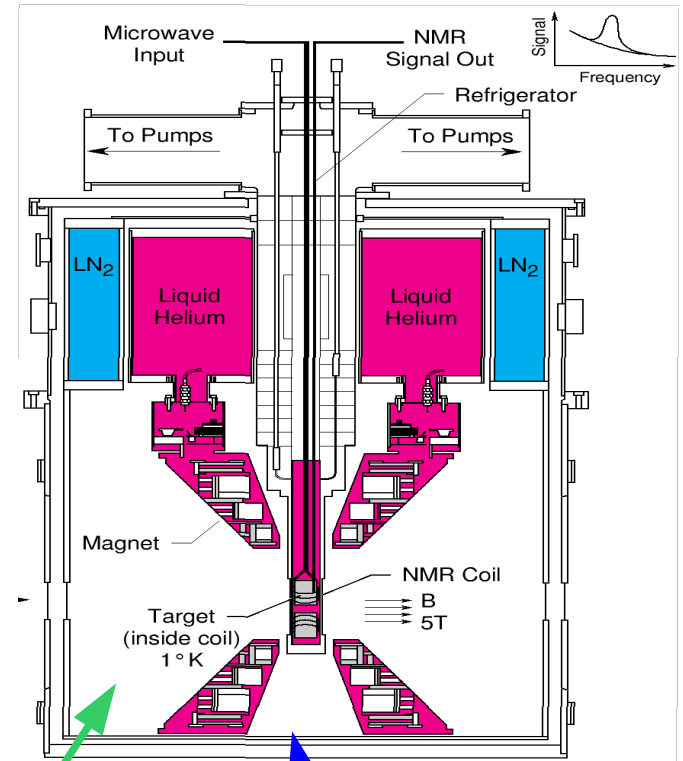
Results on ^3He structure available as well.
See K. Slifer's talk yesterday.



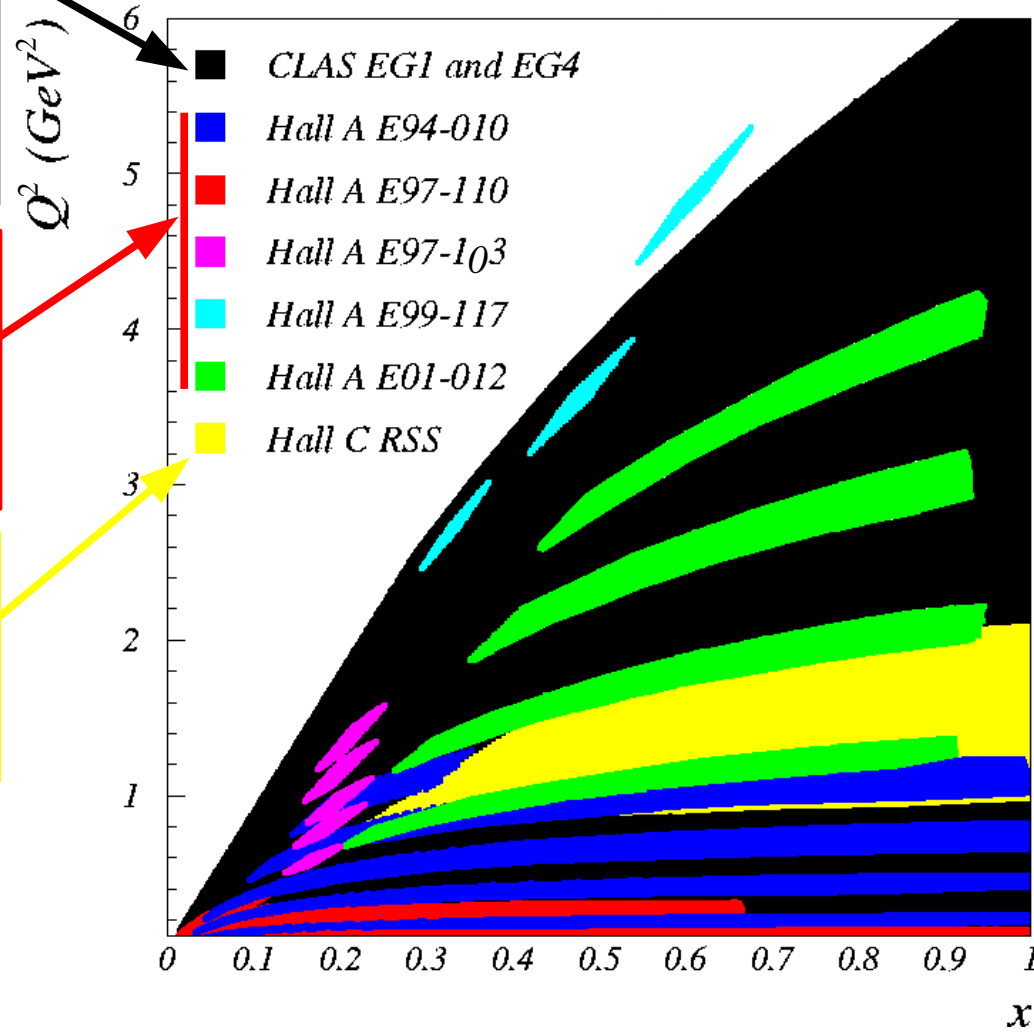
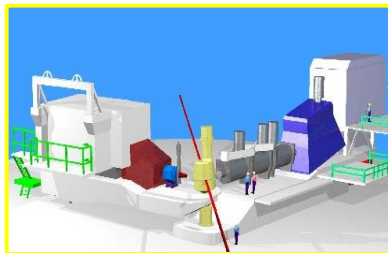
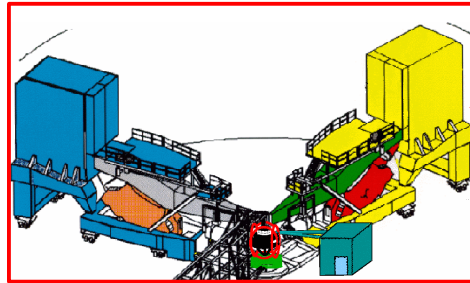
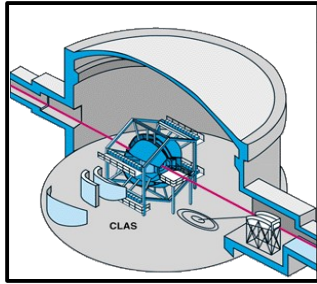
Polarized targets

Ammonia targets:

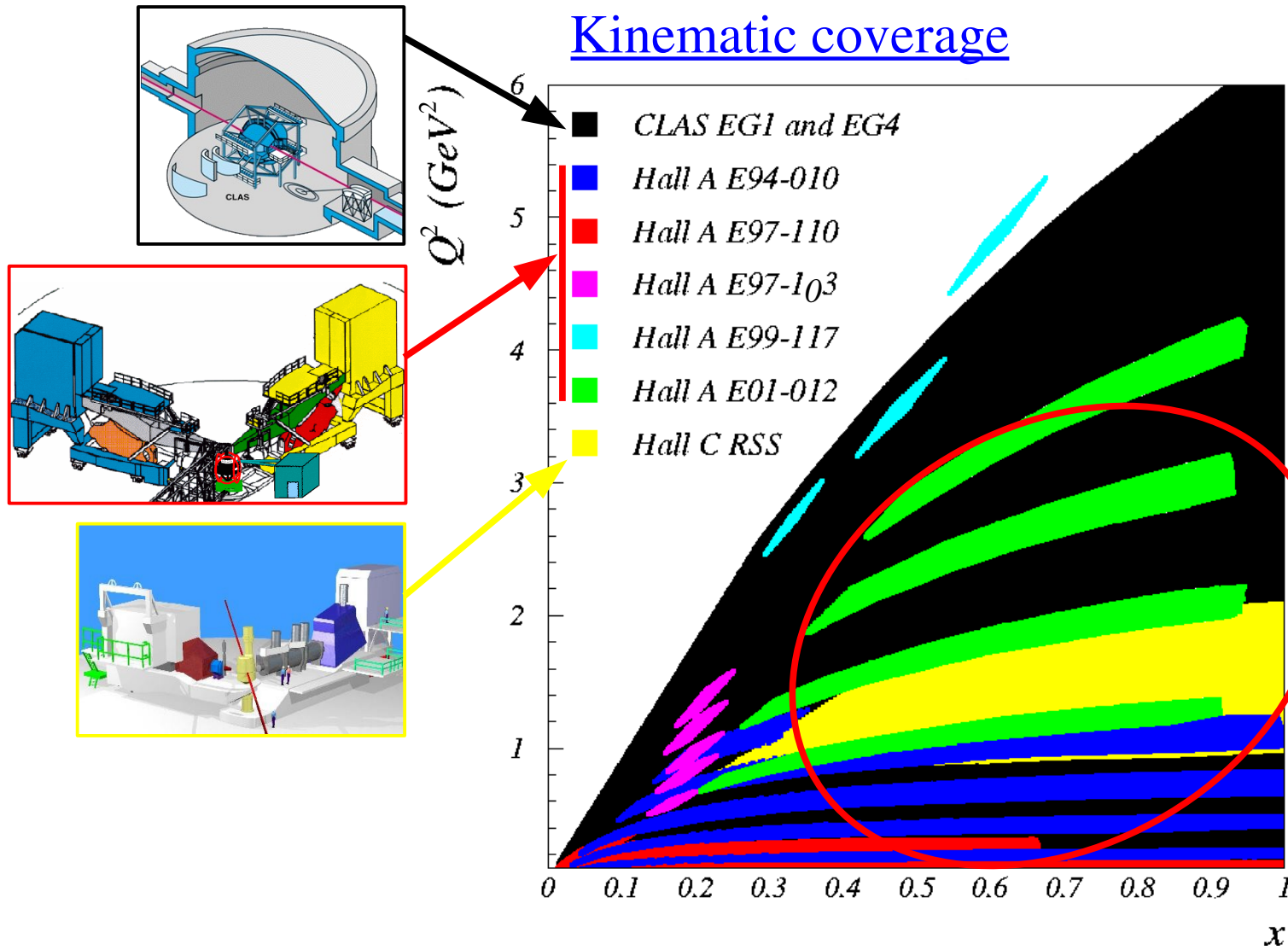
- Polarized proton & deuteron
- Good luminosity: 10^{34} (B) & 10^{35} (C) $s^{-1}cm^{-2}$
- High dilution: ~15%
- High polarization: ~80% (p)
~40% (d)
- Longitudinal polarization (B)
Longitudinal and transverse polarization (C)



Kinematic coverage

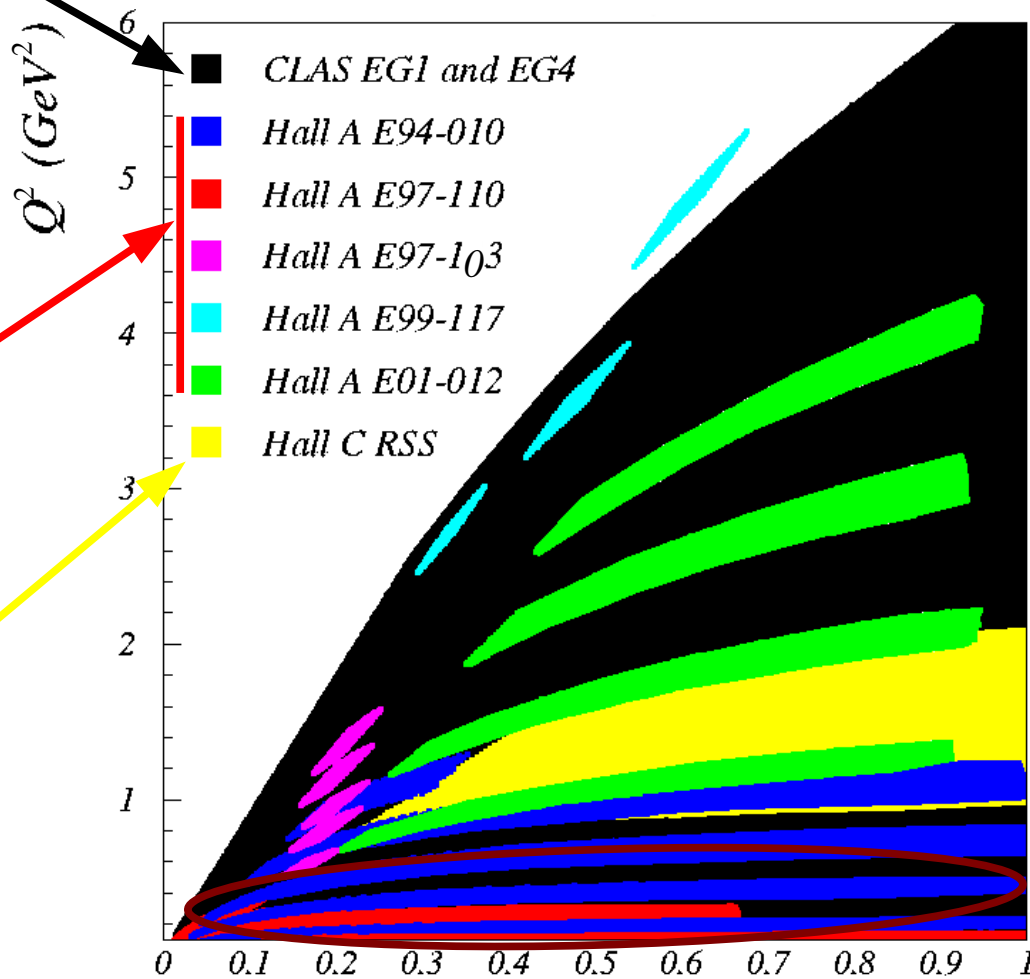
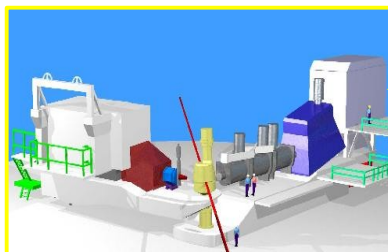
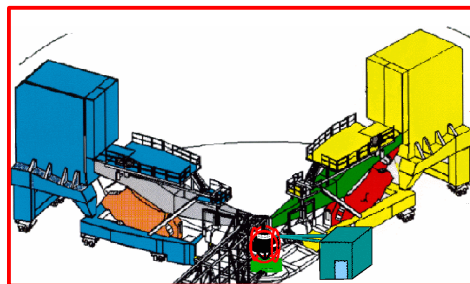
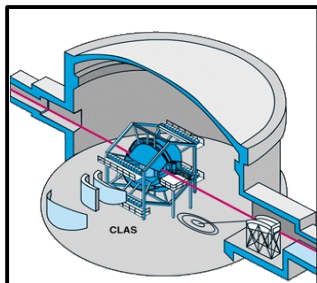


Kinematic coverage



Transition from
short scales
(pQCD)
to large scales

Kinematic coverage



Effective theories
of strong
interaction
at large distances

Moments of spin structure functions

$$N^{\text{th}}\text{-moments: } \left\{ \begin{array}{l} \int g_1 x^{n-1} dx \\ \int g_2 x^{n-1} dx \end{array} \right. \quad \text{First moments: } \Gamma_1, \Gamma_2$$

$$\star \Gamma_1^N: \left\{ \begin{array}{l} \text{Ellis-Jaffe sum rule (large } Q^2) \\ \text{Gerasimov-Drell-Hearn (GDH) sum rule (} Q^2=0) \end{array} \right.$$

$$\star \Gamma_1^{p-n}: \text{Bjorken sum rule (large } Q^2)$$

$$\star \Gamma_2^N: \text{Burkhardt-Cottingham (BC) sum rule (any } Q^2)$$

$$\left. \begin{array}{l} \star d_2 \\ \star \text{Spin polarizability} \end{array} \right\} \text{No low-}x \text{ extrapolation issue}$$

In this talk, I will focus on **moments**. Structure Functions are (obviously) available too.

Transition from short to large scales: results on $\int g_1^p - g_1^n dx$

At large Q^2 , proportional to axial charge of the nucleon g_a (**Bjorken sum rule**):

$$\int g_1^p - g_1^n dx = \frac{1}{6} g_a (1 + f(Q^2))$$

$f(Q^2)$: series in α_s fully calculable within pQCD.

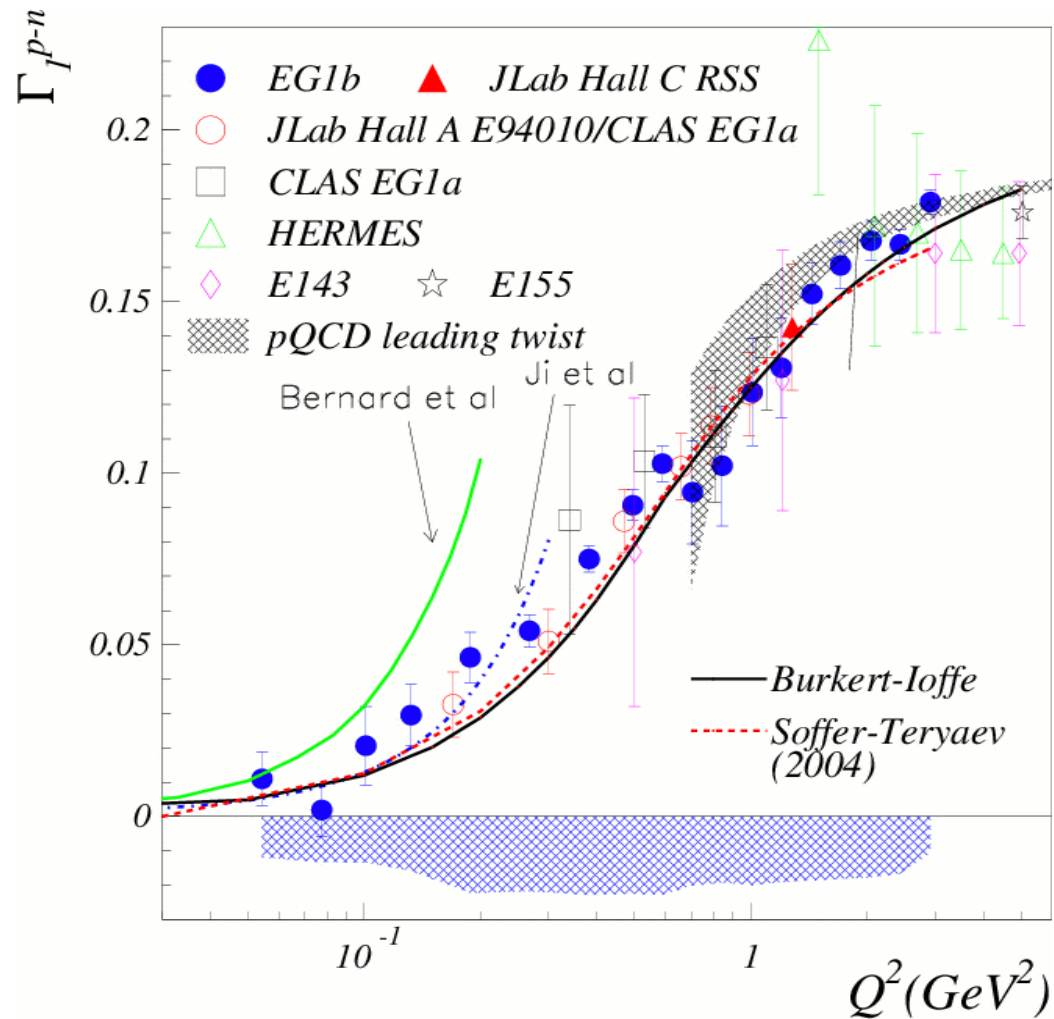
At intermediate Q^2 , (lattice QCD) proportional to spin-dependent Compton amplitude.

At small Q^2 , (χ pT) proportional to spin-dependent Compton amplitude.

At $Q^2 \rightarrow 0$, proportional to anomalous magnetic moments squared of the nucleons (**Gerasimov-Drell-Hearn sum rule**, applies also to individual nucleons):

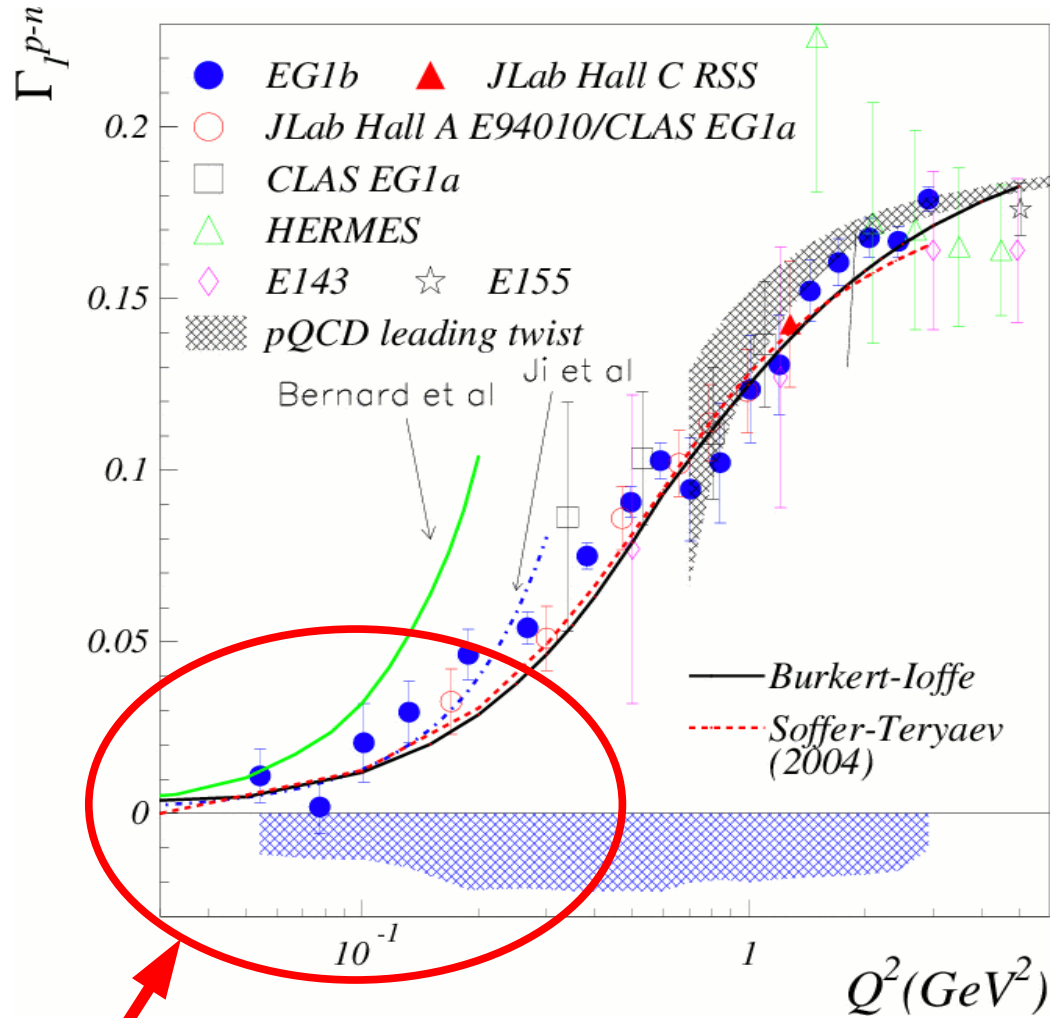
$$\int g_1^p - g_1^n dx = \frac{-Q^2}{8} \left(\frac{\kappa_p^2}{M_p^2} - \frac{\kappa_n^2}{M_n^2} \right)$$

Transition from short to large scales: results on $\int g_1^p - g_1^n dx$



Δ contribution suppressed
 \Rightarrow Easier check of χpT .

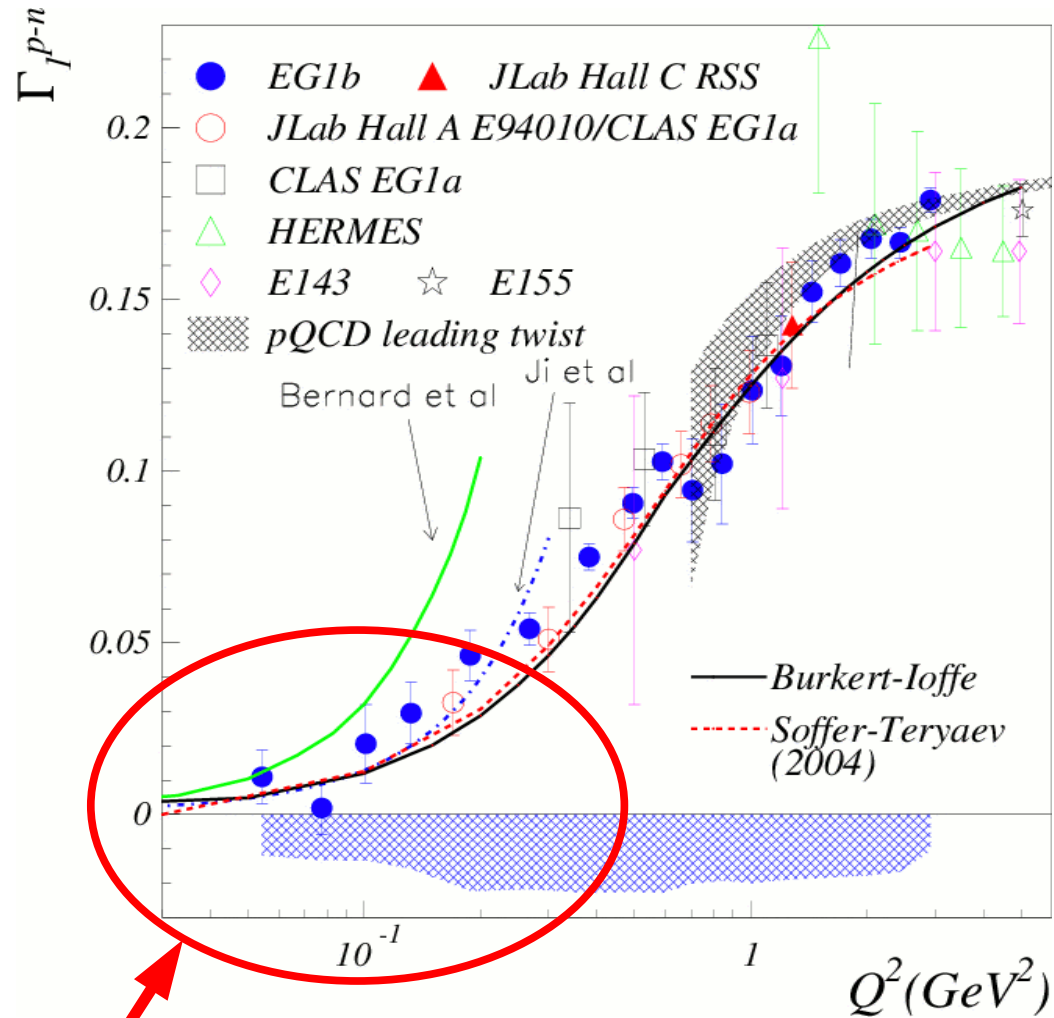
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Nice agreement with χpT (Δ suppressed?)

Transition from short to large scales: results on $\int g_1^{p-n} dx$



Δ contribution suppressed
 \Rightarrow Easier check of χ_{pT} .

Low Q^2 fit:

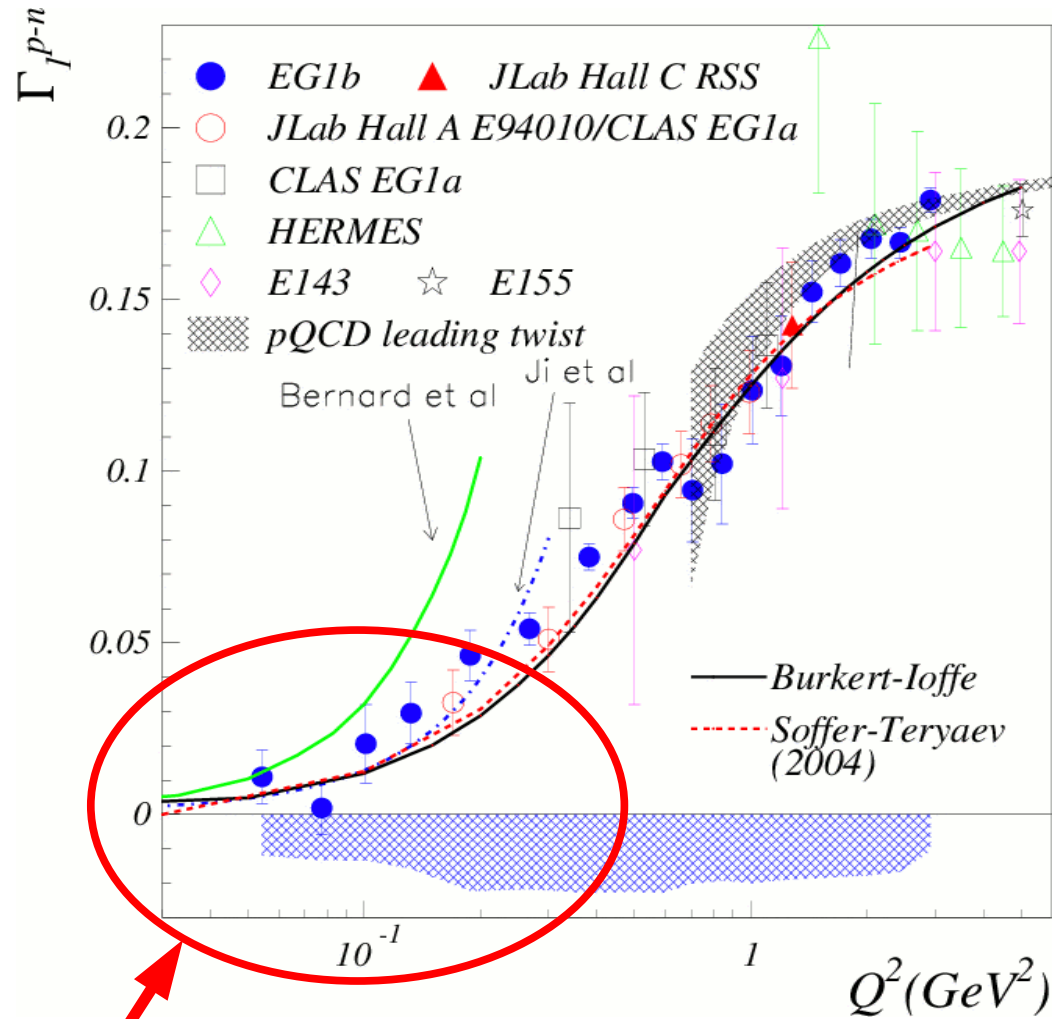
$$\Gamma_1^{p-n} = \frac{\kappa_n^2 - \kappa_p^2}{8M^2} Q^2 + aQ^4 + bQ^6$$

$$a = 0.80 \pm 0.07 \pm 0.23, \quad b = -1.13 \pm 0.16 \pm 0.39$$

$$a^{\chi_{pT, Ji}} = 0.74, \quad a^{\chi_{pT, B.}} = 2.4$$

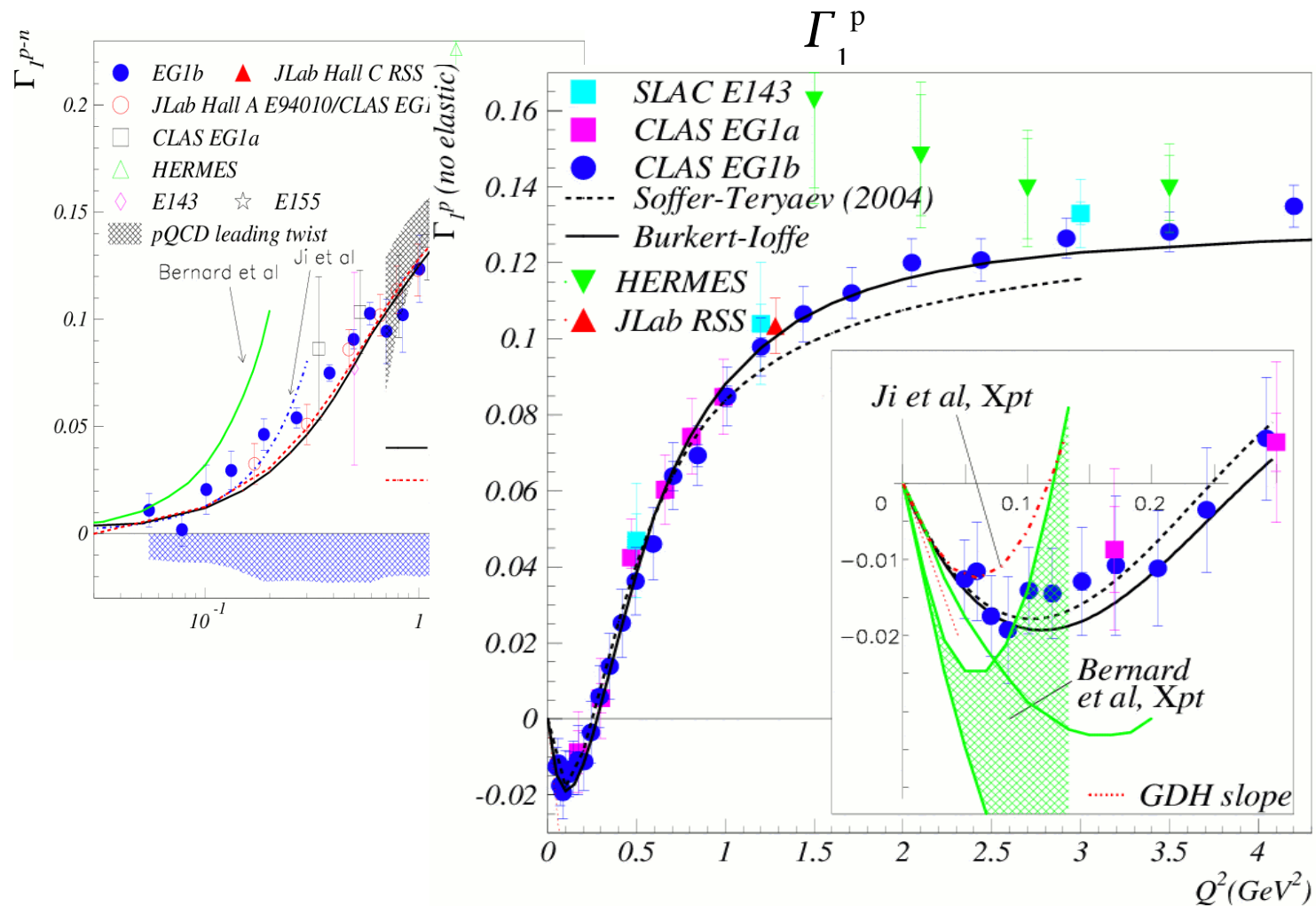
Nice agreement with χ_{pT} (Δ suppressed?)

Transition from short to large scales: results on $\int g_1^p - g_1^n dx$

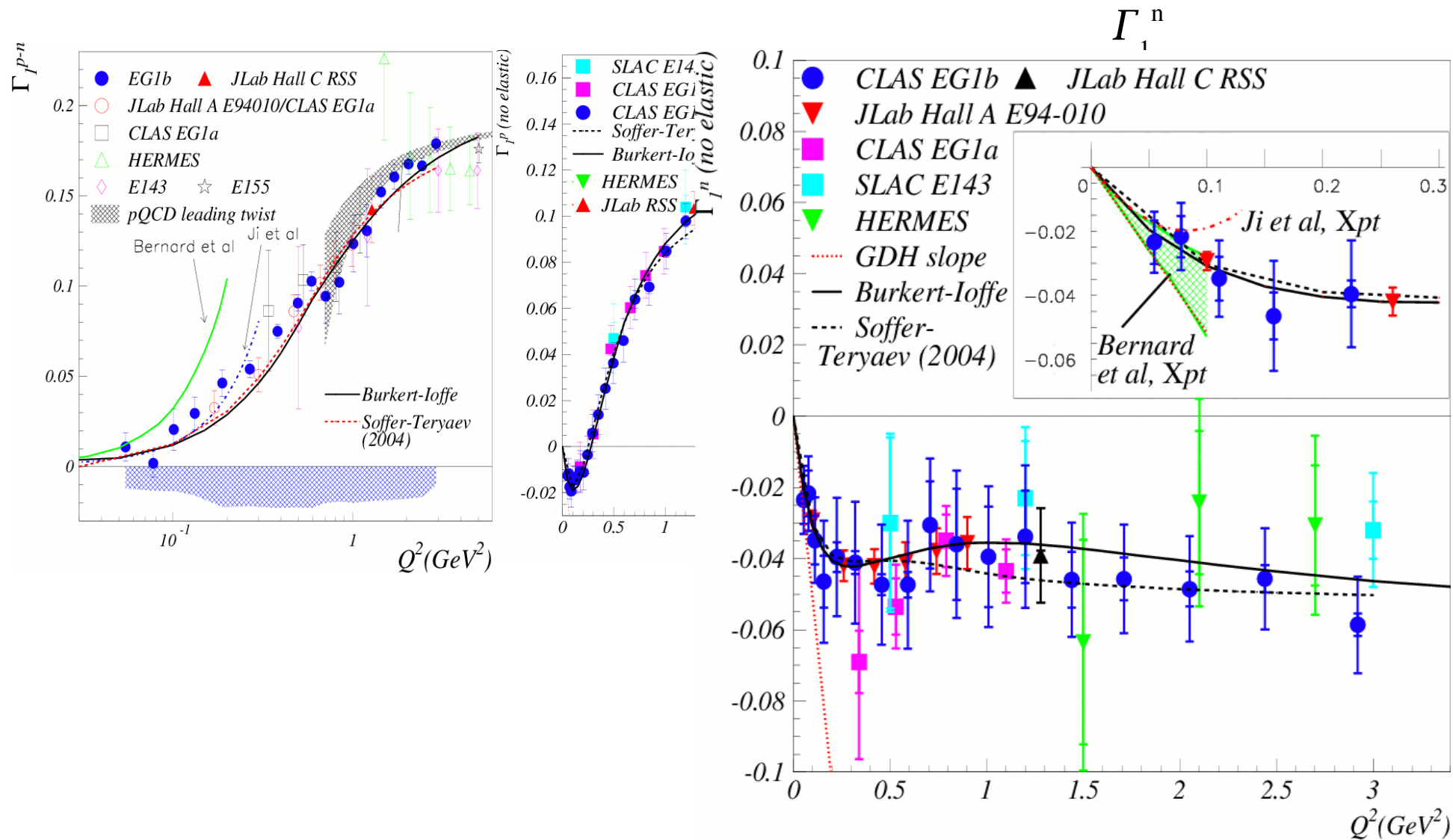


Nice agreement with χpT (Δ suppressed?)

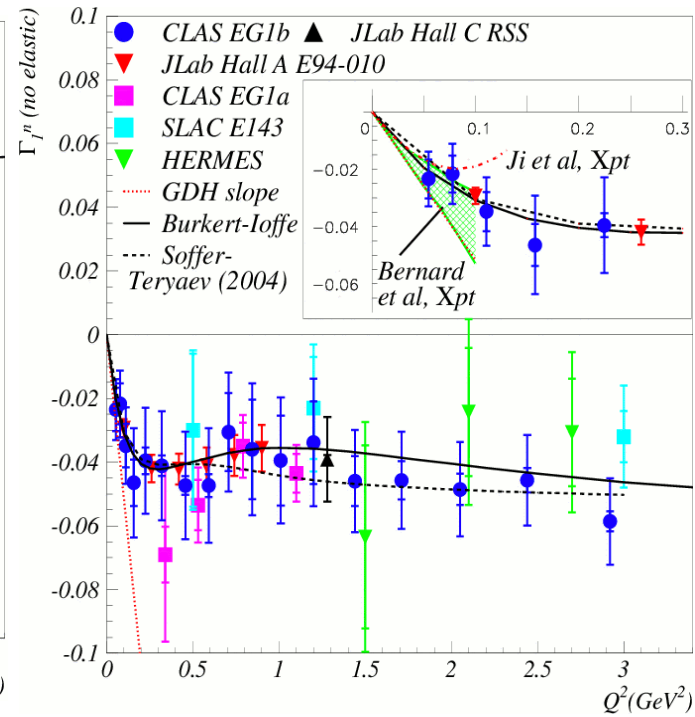
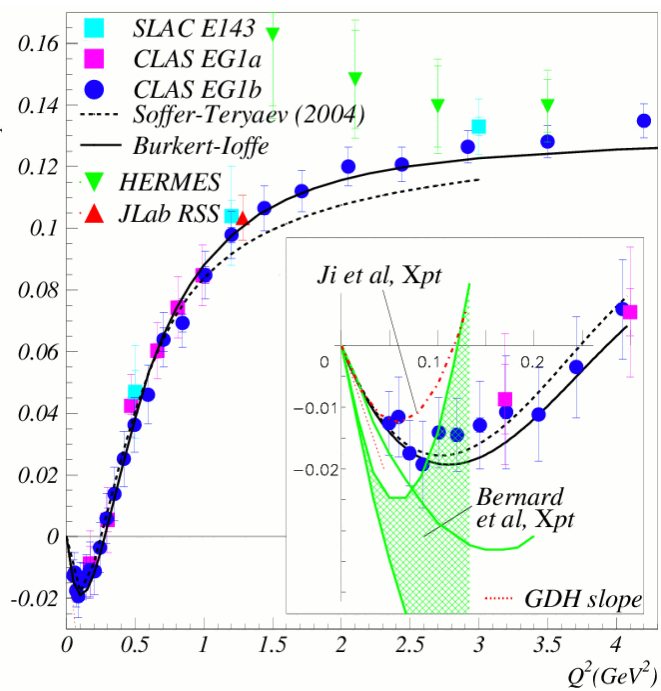
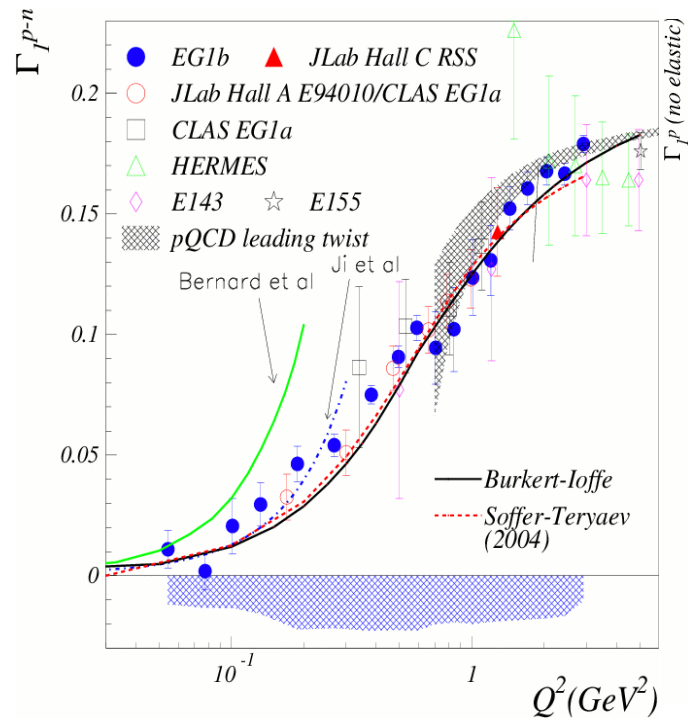
Transition from short to large scales: results on $\int g_1 dx$



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Transition from short to large scales: results on $\int g_1 dx$



Results on sum rules (higher moments)

Generalized forward spin polarizability:

$$\gamma_0 = \frac{4e^2 M^2}{\pi Q^6} \int x^2 (g_1 - \frac{4M^2}{Q^2} x^2 g_2) dx$$

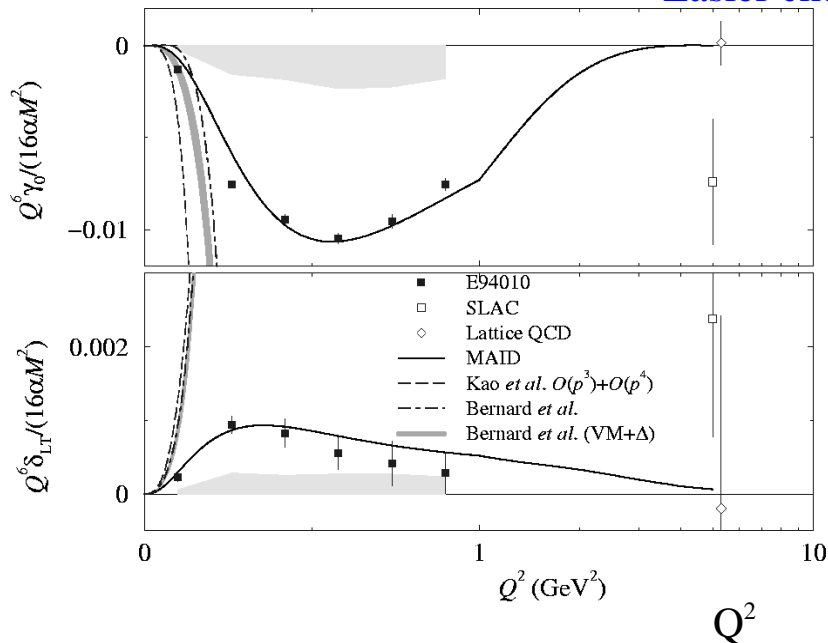
For Neutron

Longitudinal-Transverse polarizability:

$$\delta_{LT} = \frac{4e^2 M^2}{\pi Q^6} \int x^2 (g_1 + g_2) dx$$

Δ contribution suppressed

\Rightarrow Easier check of χ pT.



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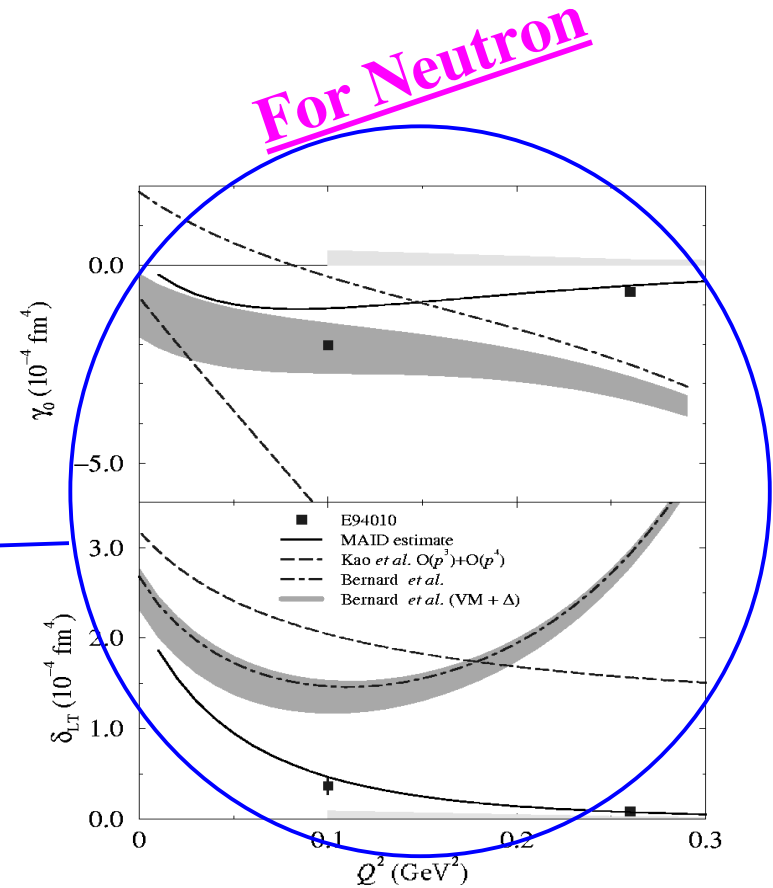
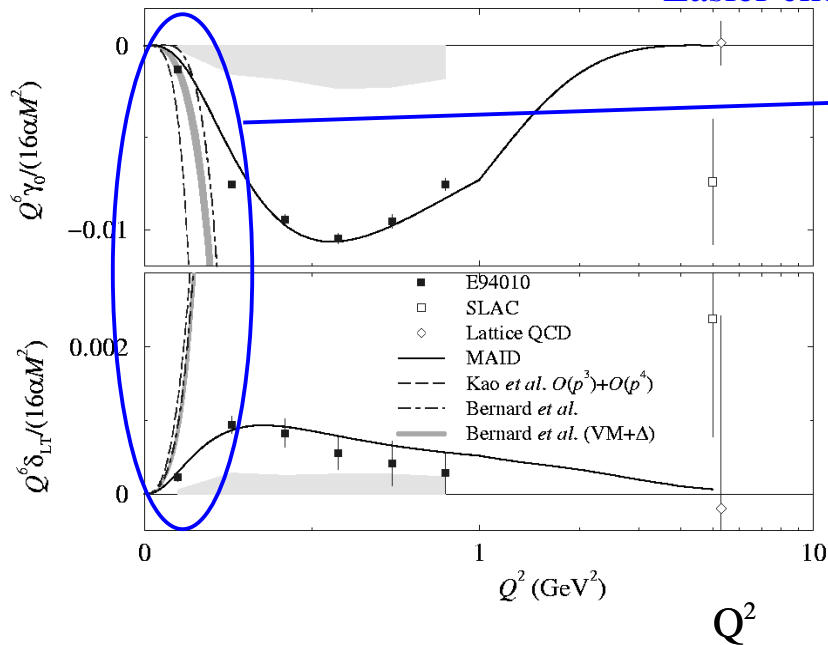
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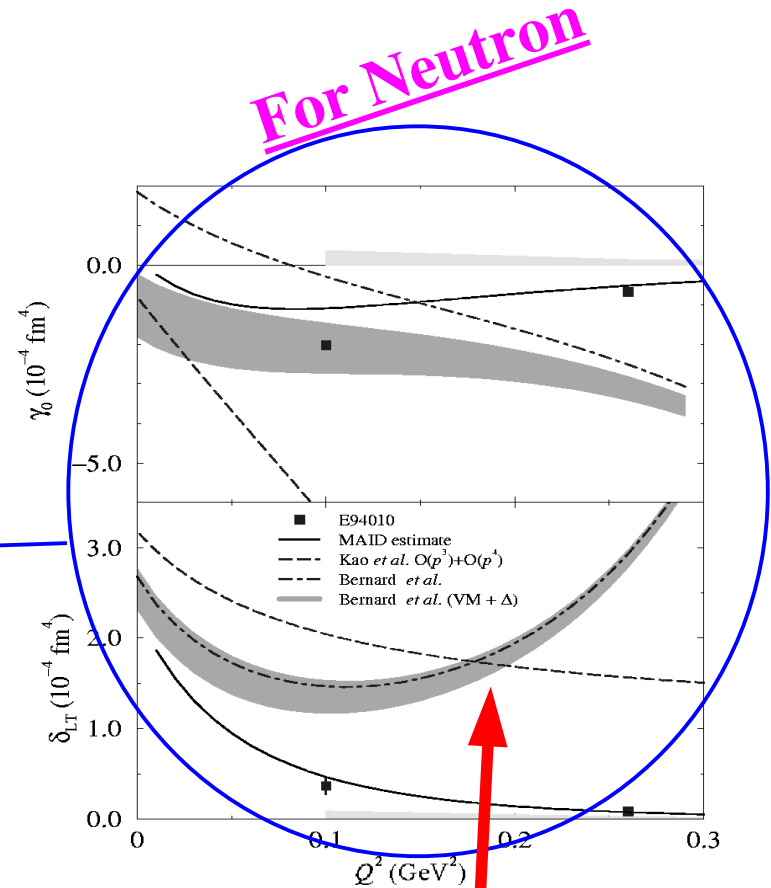
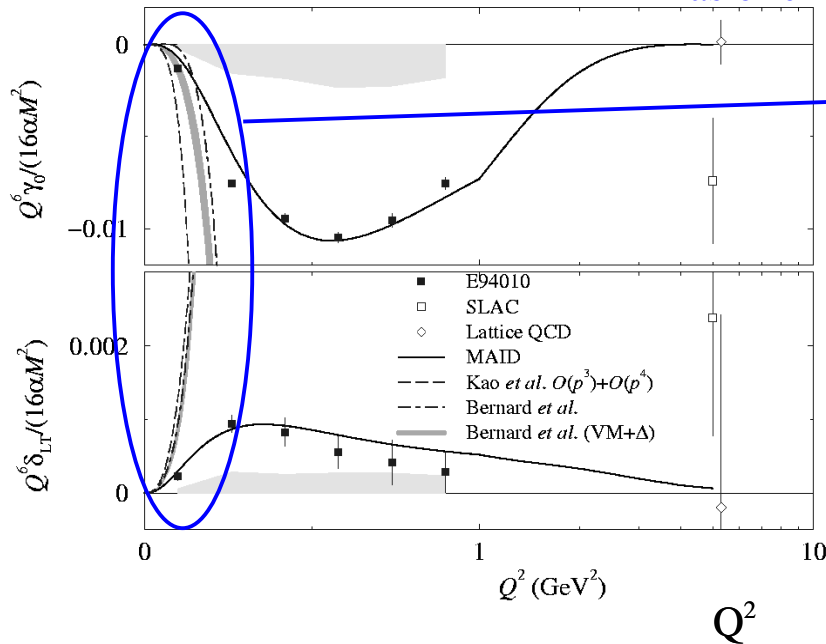
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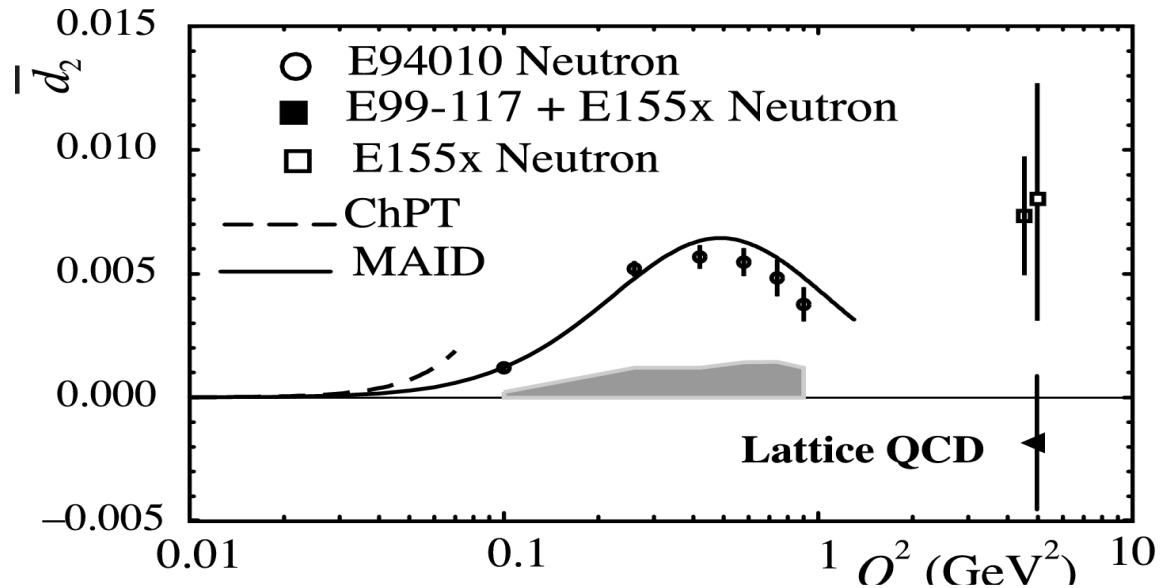


**Failure of calculations
in spite of Δ suppression.**

Results on sum rules (higher moments)

$$d_2 = \int x^2 (2g_1 + 3g_2) dx$$

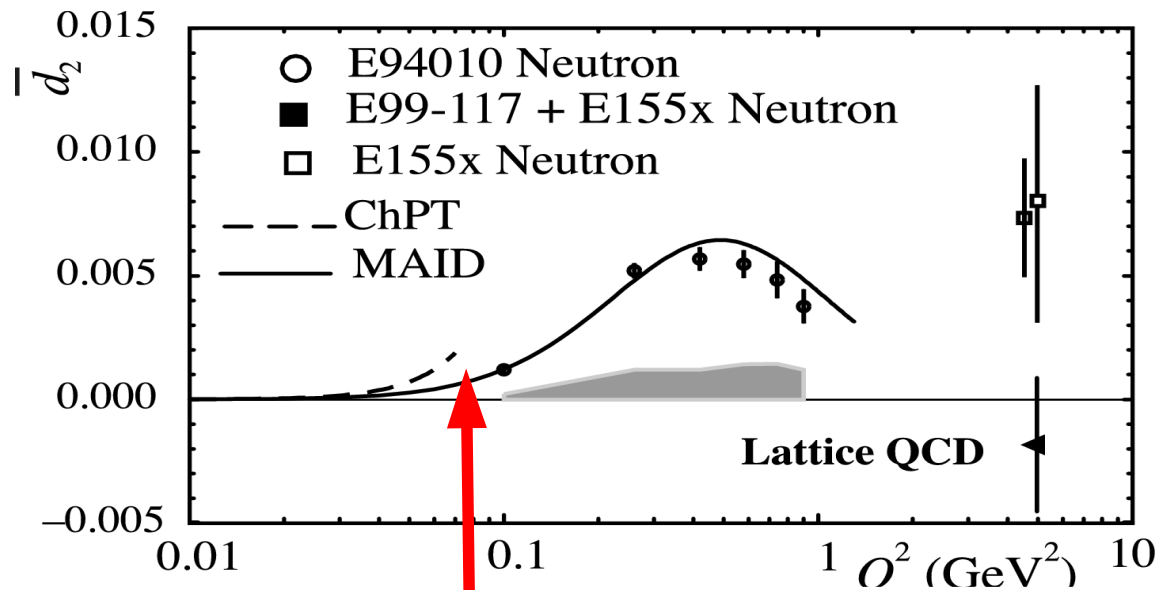
For Neutron



Results on sum rules (higher moments)

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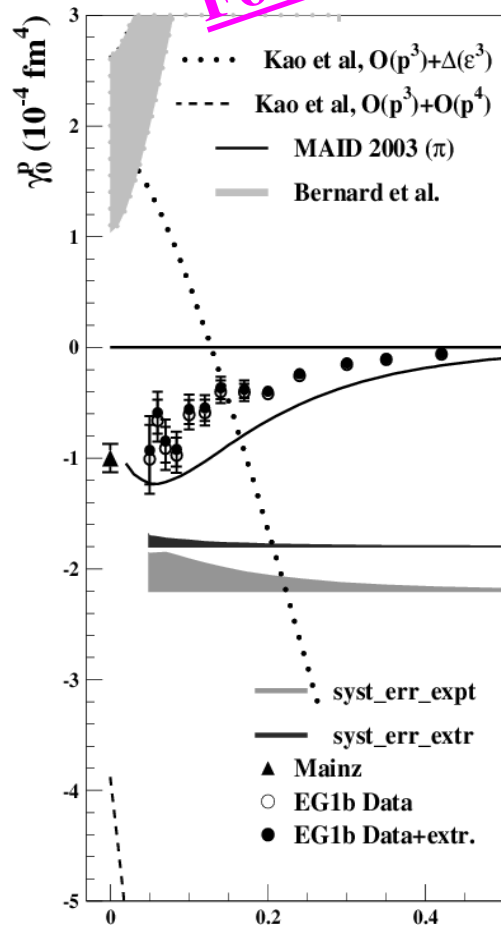
For Neutron



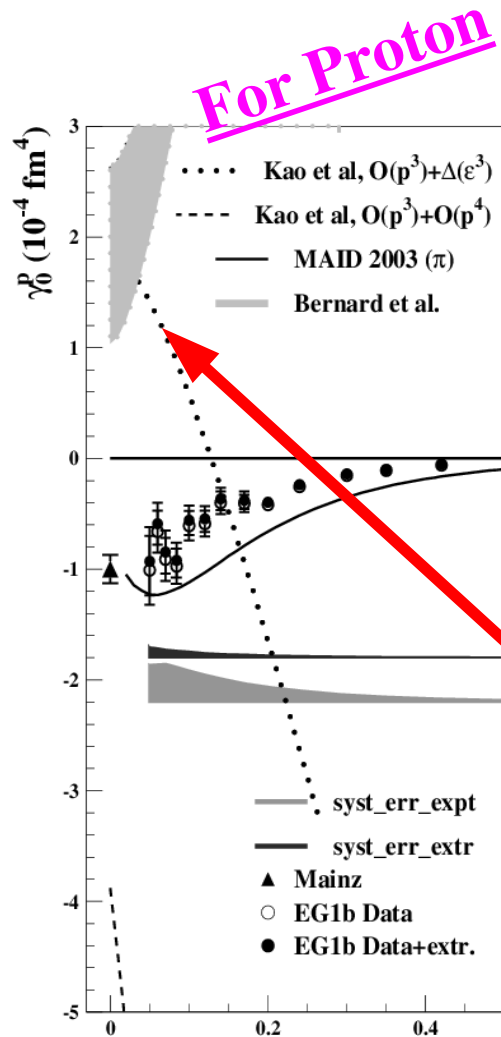
Failure of calculations
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Results on sum rules (higher moments)

For Proton



Results on sum rules (higher moments)

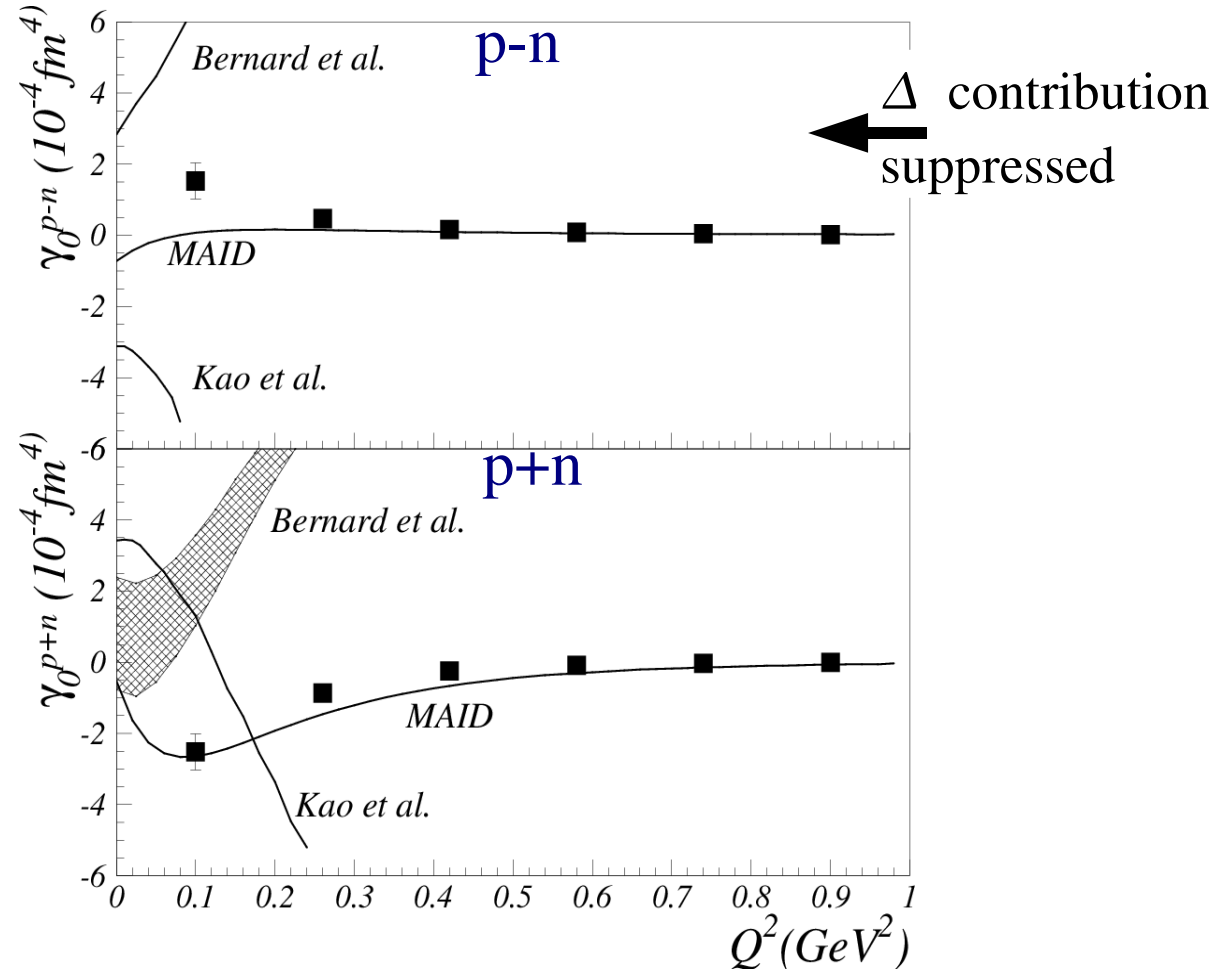
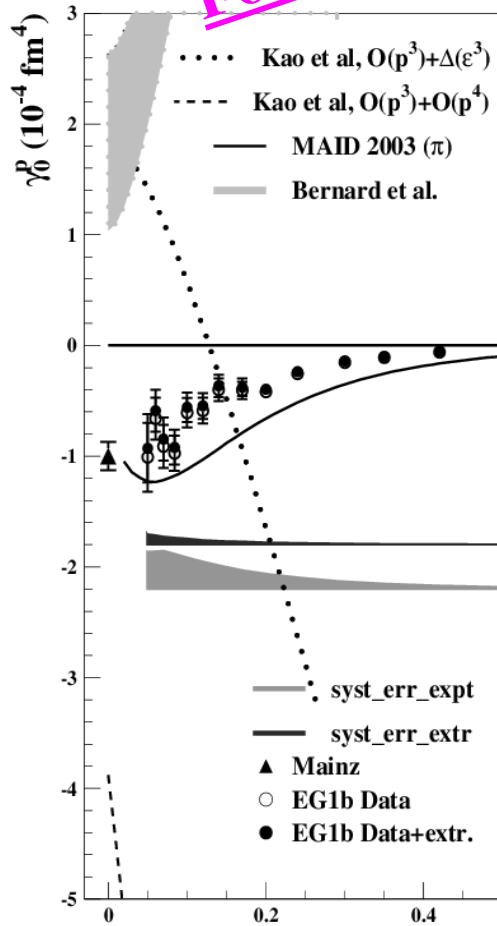


Failure of
calculations

Results on sum rules (higher moments)

For Proton

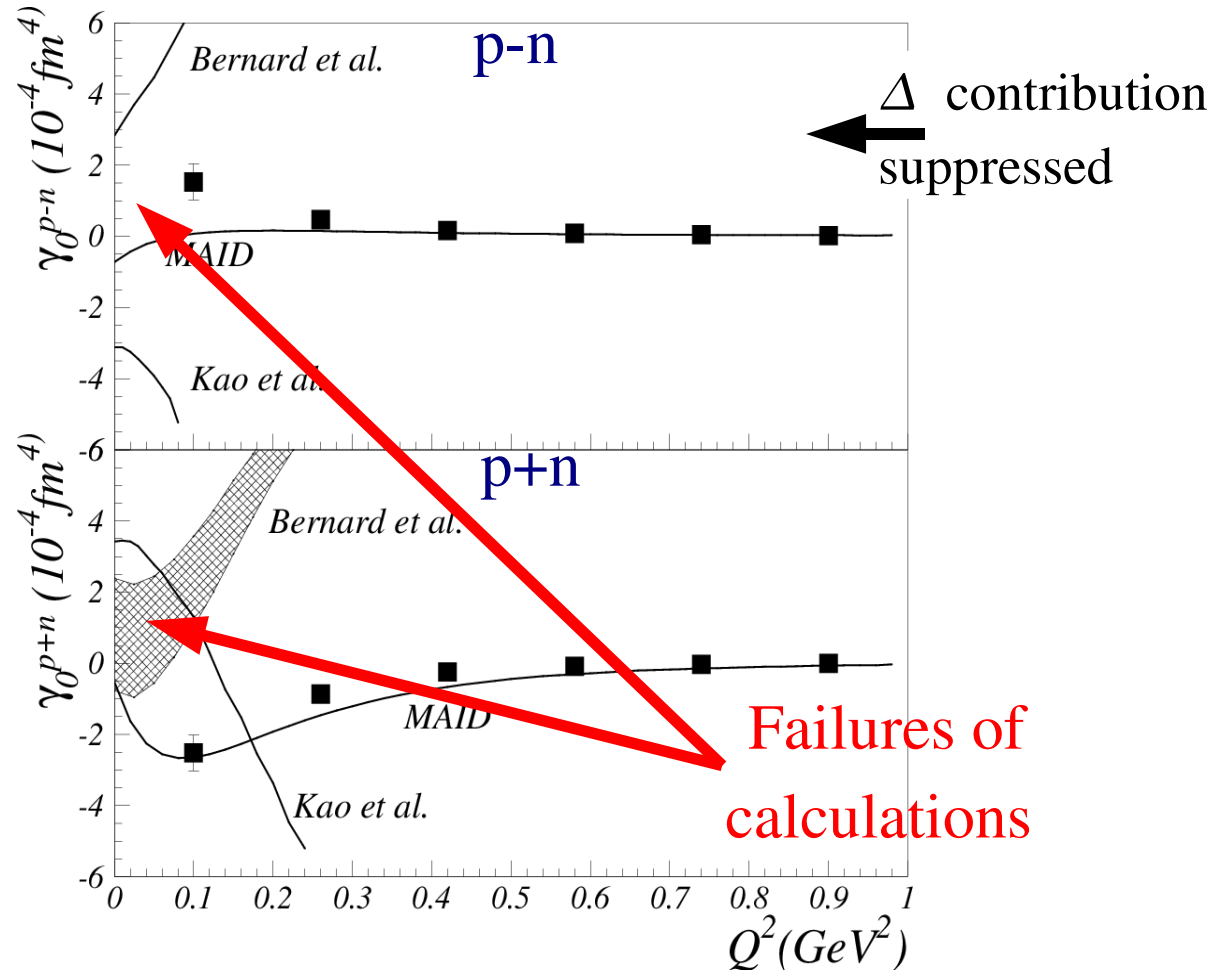
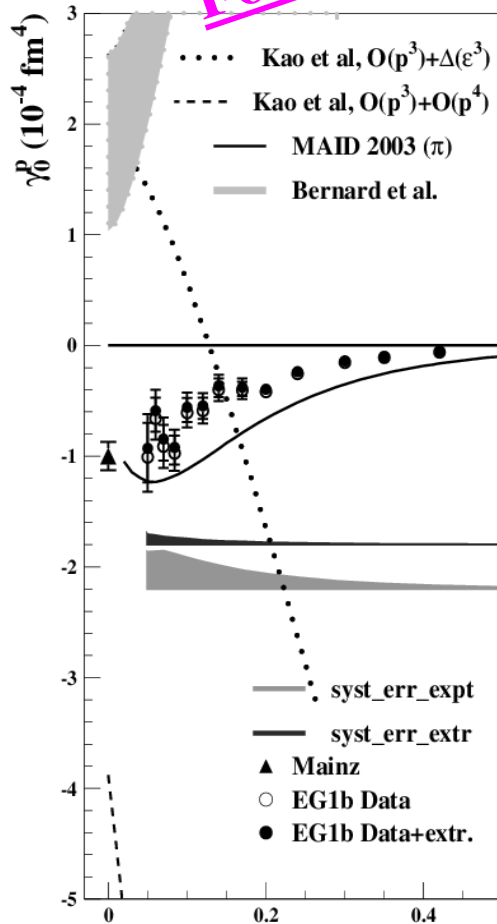
To study the influence of the Δ : Isospin decomposition of γ_0 using the Hall A and B data.



Results on sum rules (higher moments)

For Proton

To study the influence of the Δ : Isospin decomposition of γ_0 using the Hall A and B data.



Results on sum rules (higher moments)

For Proton

δ_{LT}^p is yet unmeasured. Desirable to measure it in order for isospin study of the “ δ_{LT} puzzle”

Summary:

χ pT:

No low-x



No low-x

No Δ



No low-x



No Δ \rightarrow

	Γ_1	γ_0	δ_{LT}	d_2
Proton	$a^{\text{exp}}=4.31\pm 0.31\pm 1.36$ $a^{\text{Ji}}=3.89$ Up to $Q^2\sim 0.08 \text{ GeV}^2$		No low Q^2 data	No low Q^2 data
Neutron		Up to $Q^2\sim 0.1 \text{ GeV}^2$ (Bernard <i>et al.</i> only)		
P-N	$a^{\text{exp}}=0.80\pm 0.07\pm 0.23$ $a^{\text{Ji}}=0.74, a^{\text{B}}=2.4$ Up to $Q^2\sim 0.3 \text{ GeV}^2$		No low Q^2 data	No low Q^2 data
P+N	$a^{\text{exp}}=6.97\pm 0.96\pm 1.48$ $a^{\text{Ji}}=7.11$ Up to $Q^2\sim 0.1 \text{ GeV}^2$		No low Q^2 data	No low Q^2 data

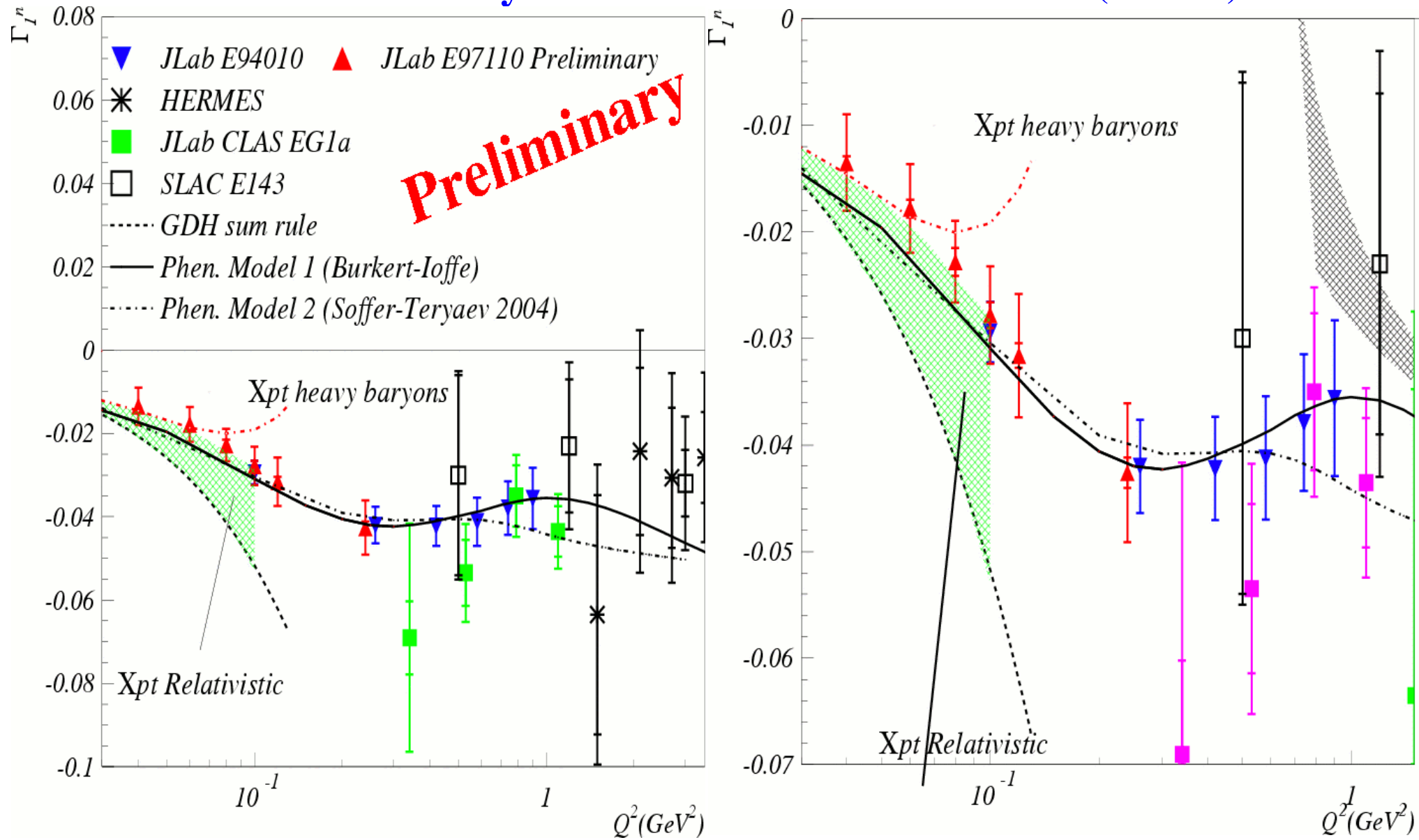
Preliminary Neutron Results from Hall A (E97110)

Experiment specially designed to access low Q^2 :

- \parallel and \perp data on neutron (ran in 2003)
- New magnet added to high resolution spectrometer
- ^3He target redesigned for low angles.
- Target moved upstream.

(See V. Sulkosky's Talk)

Preliminary Neutron Results from Hall A (E97110)



Preliminary

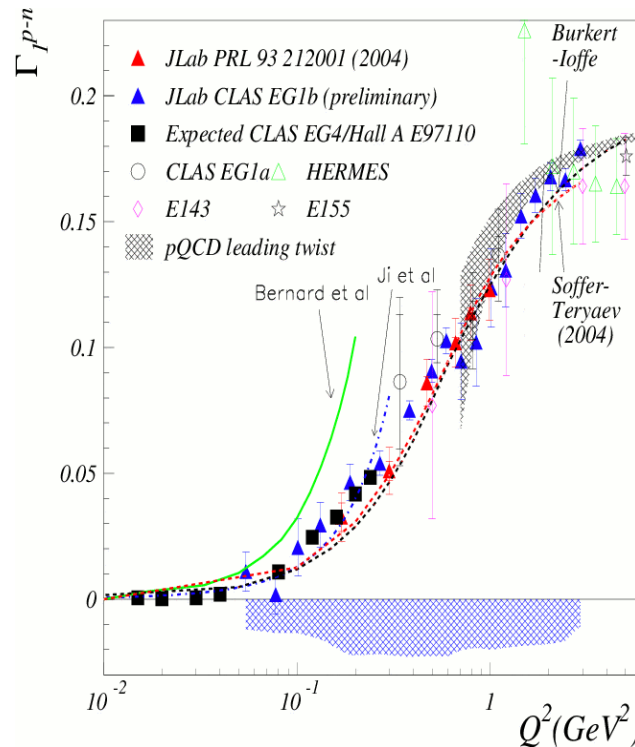
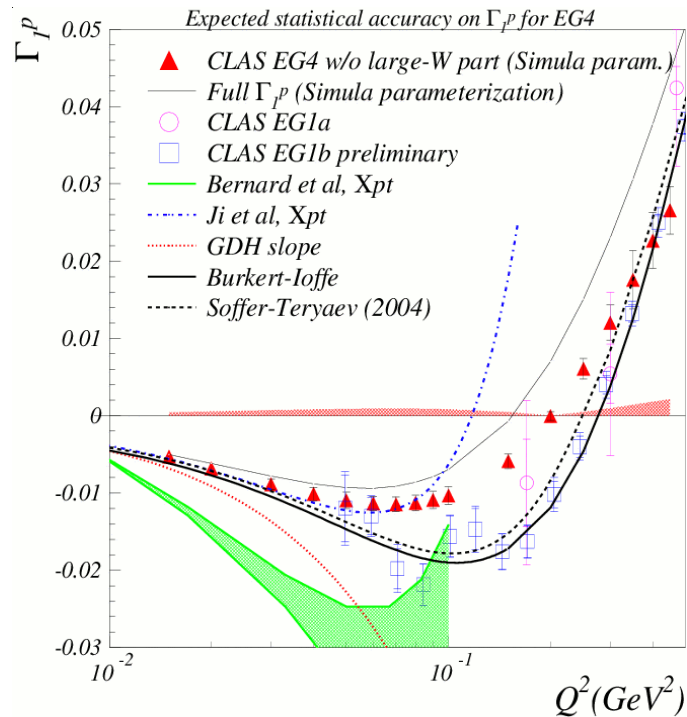
Preliminary Proton Results from Hall B (EG4)

Experiment specially designed to access low Q^2 :

- || data on proton and deuteron (ran in 2006)
- New detector added to large acceptance spectrometer
- NH_3 & ND_3 targets moved upstream.

Inclusive results not available yet.

(See S. Phillips' Talk)



Preliminary results on pion production asymmetries available (Xiaochao Zheng).

Observables in Pion Electroproduction

➤ NH3: $\vec{e} \vec{p} \rightarrow e' \pi^+ n$ and $\vec{e} \vec{p} \rightarrow e' \pi^0 p$

➤ ND3: $\vec{e} \vec{n} \rightarrow e' \pi^- p$ and $\vec{e} \vec{p} \rightarrow e' \pi^+ n$

● Cross section:

$$\frac{d\sigma}{d\Omega_\pi^*} \sim \frac{d\sigma_{unp}}{d\Omega_\pi^*} + P_e \frac{d\sigma_e}{d\Omega_\pi^*} + P_t \frac{d\sigma_t}{d\Omega_\pi^*} + P_e P_t \frac{d\sigma_{et}}{d\Omega_\pi^*}.$$

● Three independent asymmetries:

➤ Single-beam $A_e = \frac{d\sigma_e}{d\sigma_{unp}} = \frac{\sigma(+h_e) - \sigma(-h_e)}{\sigma(+h_e) + \sigma(-h_e)}$

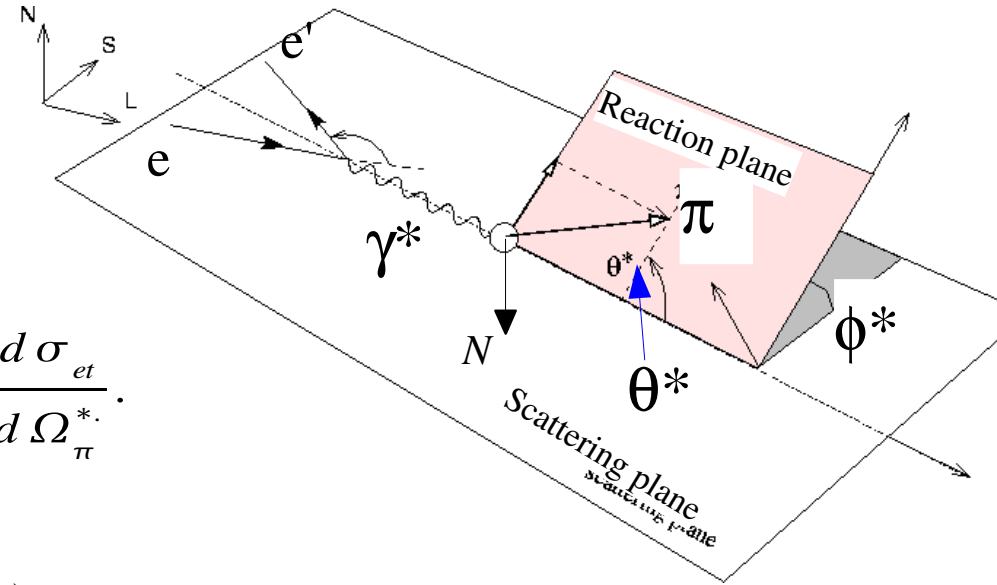
accessible from unpolarized target data

➤ Single-target $A_t = \frac{d\sigma_t}{d\sigma_{unp}} = \frac{\sigma(+h_N) - \sigma(-h_N)}{\sigma(+h_N) + \sigma(-h_N)}$

only accessible from polarized target data

➤ Double beam-target

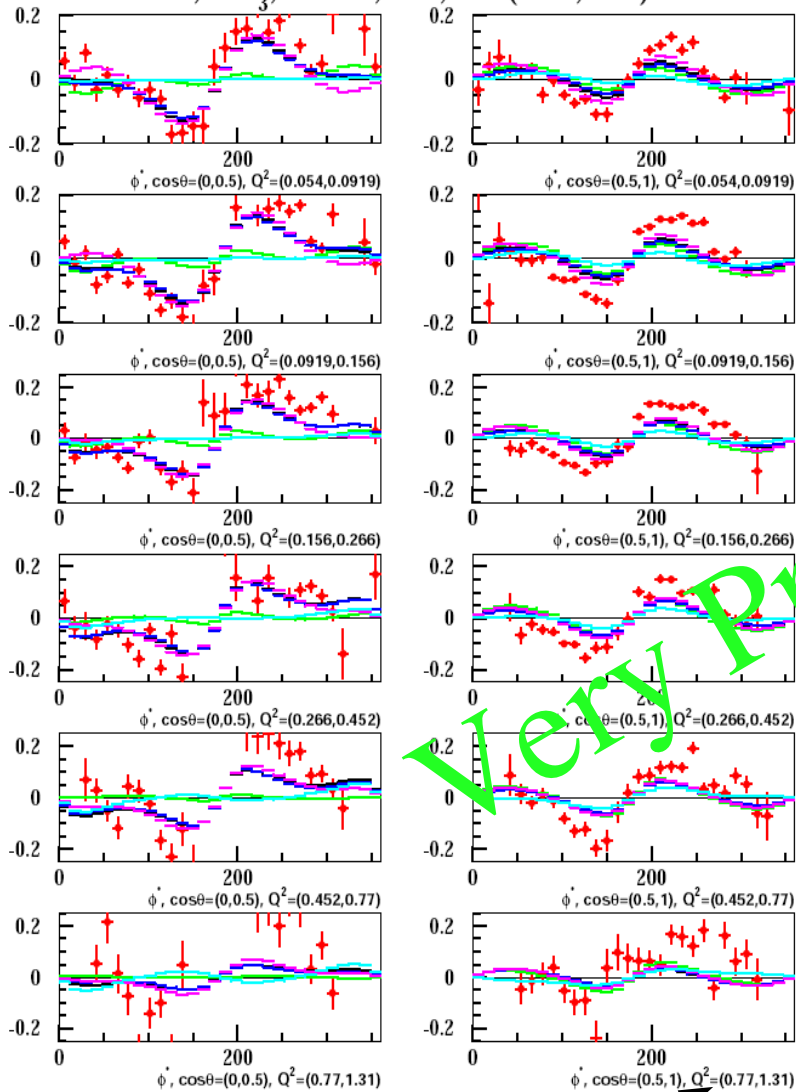
$$A_{et} = \frac{d\sigma_{et}}{d\sigma_{unp}} = \frac{\sigma(+h_e, +h_N) + \sigma(-h_e, -h_N) - \sigma(+h_e, -h_N) - \sigma(-h_e, +h_N)}{\sigma(+h_e, +h_N) + \sigma(-h_e, -h_N) + \sigma(+h_e, -h_N) + \sigma(-h_e, +h_N)}$$



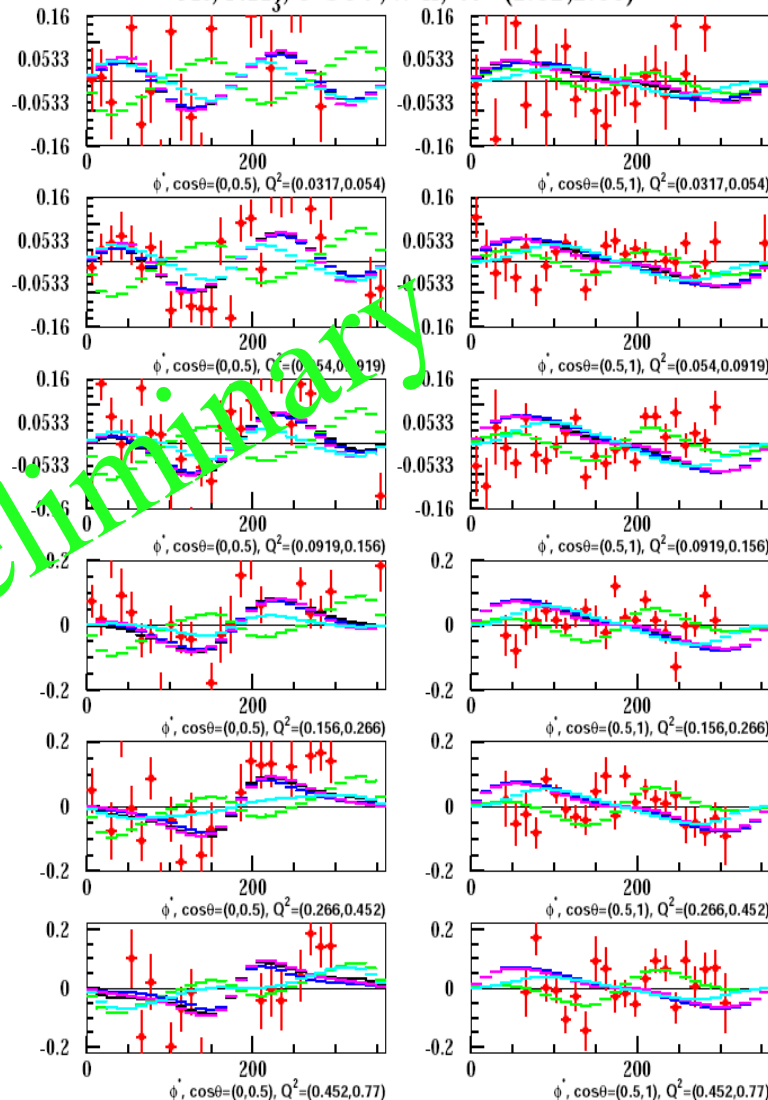
At vs. ϕ^*

Preliminary Results from Hall B (EG4)

At, NH₃, 3 GeV, π^+n , W=(1.58,1.82)



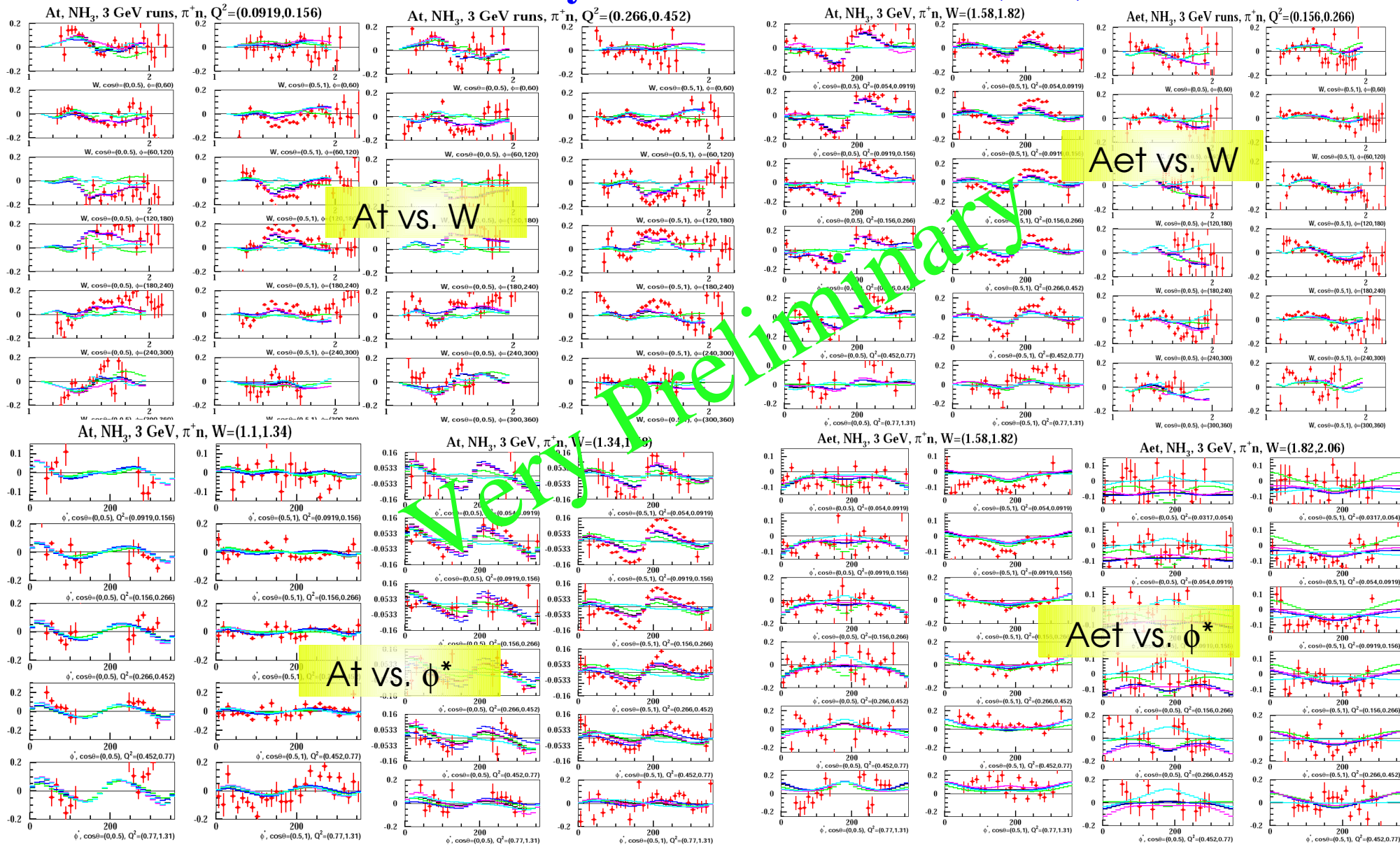
At, NH₃, 3 GeV, π^+n , W=(1.82,2.06)



3 GeV
(2 GeV,
1.3 GeV and
1.0 GeV
beam energies
available too)

Data; MAID2007; DMT; MAID2007(P11off); MAID2007(S11off); MAID2007(D13off)

Preliminary Results from Hall B (EG4)



At vs. W

Aet vs. W

At vs. ϕ^*

Aet vs. ϕ^*

Very Preliminary

Perspectives

E08027 (A. Camsonne, J.P. Chen and K. Slifer spokespersons)

- || **and** \perp data on proton (approved to run in Hall A).
- Use of new forward angle detection of Hall A.
- Import the NH_3 (& ND_3) target from Hall C.
- Dedicated to measure the missing $\delta_{\text{LT}}^{\text{p}}$.

Possibility for \perp data on proton and deuteron in Hall B is opening (HDice target).

- Target imported from BNL (LEGS, A. Sandorfi spokesperson and HD target group leader).
- Low dilution target.
- Used only with photons so far.
- Scheduled to be tested with electrons at end of 2010.
- If successful, open possibilities for transverse target polarization studies in Hall B.

Perspectives

		No low-x ↓	No low-x No Δ ↓	No low-x ↓
	Γ_1	γ_0	δ_{LT}	d_2
Proton	$a^{\text{exp}}=4.31\pm 0.31\pm 1.36$ $a^{\text{Ji}}=3.89$ Up to $Q^2\sim 0.08 \text{ GeV}^2$			
Neutron		Up to $Q^2\sim 0.1 \text{ GeV}^2$ (Bernard <i>et al.</i> only)		
No Δ → P-N	$a^{\text{exp}}=0.80\pm 0.07\pm 0.23$ $a^{\text{Ji}}=0.74, a^{\text{B}}=2.4$ Up to $Q^2\sim 0.3 \text{ GeV}^2$			
P+N	$a^{\text{exp}}=6.97\pm 0.96\pm 1.48$ $a^{\text{Ji}}=7.11$ Up to $Q^2\sim 0.1 \text{ GeV}^2$			

Conclusions

- Data on SSF moments at low Q^2 and χ pT do not consistently agree (or disagree).
 - (Implication of these discrepancies to χ pT extrapolation for lattice results?)
 - Δ cannot be the explanation for some disagreements.
 - Low- Q^2 fits provide quantitative comparisons. Importance of Q^6 terms.
 - Exclusive data on pion-electroproduction at Low- Q^2 . What can χ pT tells us?
 - Short term: high precision experiments at lower Q^2 :
 - E97110: \parallel and \perp on neutron (ran in 2003 in Hall A)
 - EG4: \parallel on proton and deuteron (ran in 2006 in Hall B)
 - Longer term:
 - E08027: \parallel and \perp on proton (approved for Hall A, to run in 2011)
- Possibility:
- \perp data on P and D in Hall B (HDice target)