

# Chiral EFT for nuclear forces with $\Delta(1232)$ degrees of freedom

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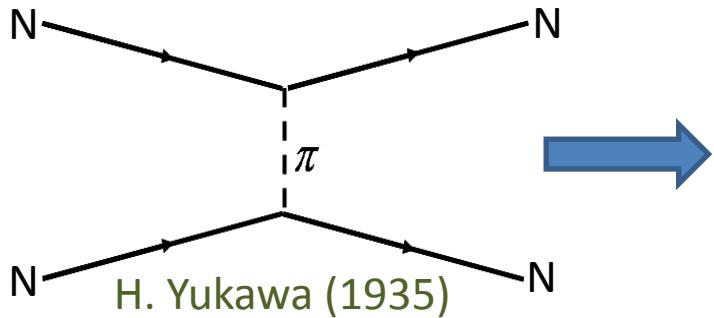
With V. Bernard, E. Epelbaum, U. Meißner

## Outline

- Nuclear forces in chiral EFT
- The role of  $\Delta(1232)$ -isobar
- Summary & Perspectives

# Nucleon-Nucleon forces

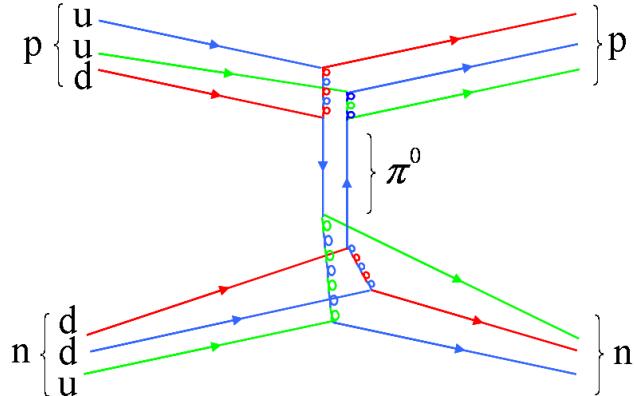
## Phenomenological description by Meson-exchange



- Boson-Exchange Models as basis for NN-force
- Highly sophisticated phen. NN potentials
- Excellent description of many experimental data
- Connection to QCD is unclear

## QCD Interpretation of NN forces

- NN force as residual strong interaction between hadrons



## Chiral EFT Interpretation of NN forces

- Model independent treatment
- At low energies NN force dominated by Goldstone Boson dynamics + short range int.
- Systematic perturbative description of few nucleon potentials
- Underlying QCD symmetries implemented by construction

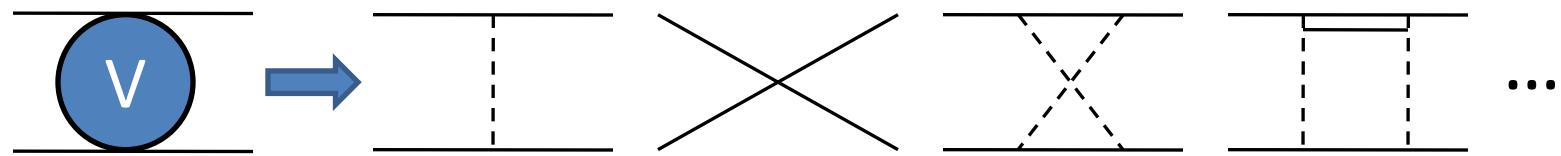
# Weinberg's scheme for NN

Weinberg, Nucl. Phys. B 363: 3 (1991)

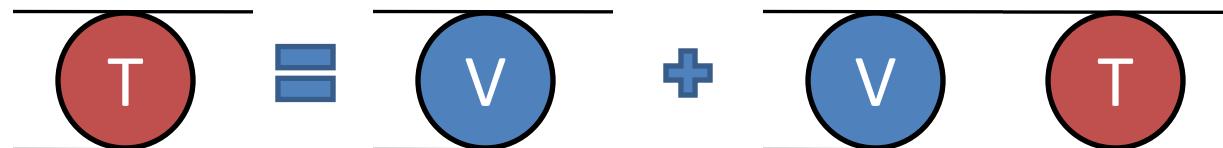
- No perturbative description for bound states



- Construct effective potential perturbatively



- Solve Lippmann-Schwinger equation nonperturbatively



# Nucleon-nucleon force up to N<sup>3</sup>LO

Ordonez et al. '94; Friar & Coon '94; Kaiser et al. '97; Epelbaum et al. '98, '03; Kaiser '99-'01; Higa et al. '03; ...

Chiral expansion for the 2N force:

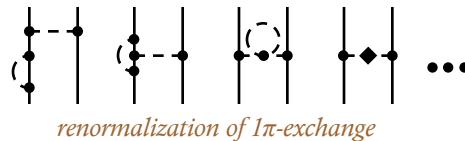
$$V_{2N} = V_{2N}^{(0)} + V_{2N}^{(2)} + V_{2N}^{(3)} + V_{2N}^{(4)} + \dots$$

• LO:



$\leftarrow 2 \text{ LECs}$

• NLO:



*renormalization of 1 $\pi$ -exchange*



*renormalization of contact terms*



*leading 2 $\pi$ -exchange*

• N<sup>2</sup>LO:

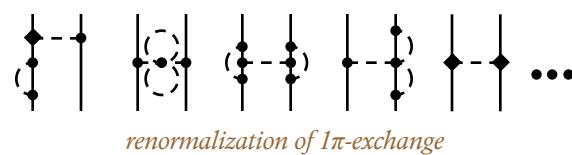


*renormalization of 1 $\pi$ -exchange*

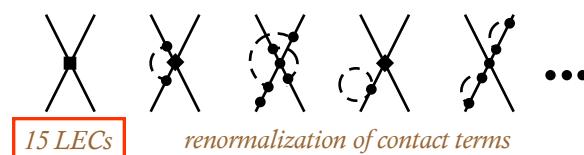


*subleading 2 $\pi$ -exchange*

• N<sup>3</sup>LO:

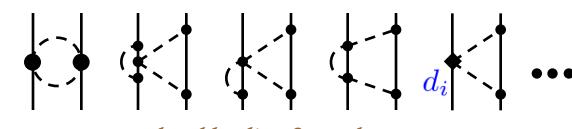


*renormalization of 1 $\pi$ -exchange*

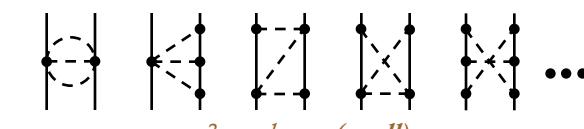


*15 LECs*

*renormalization of contact terms*



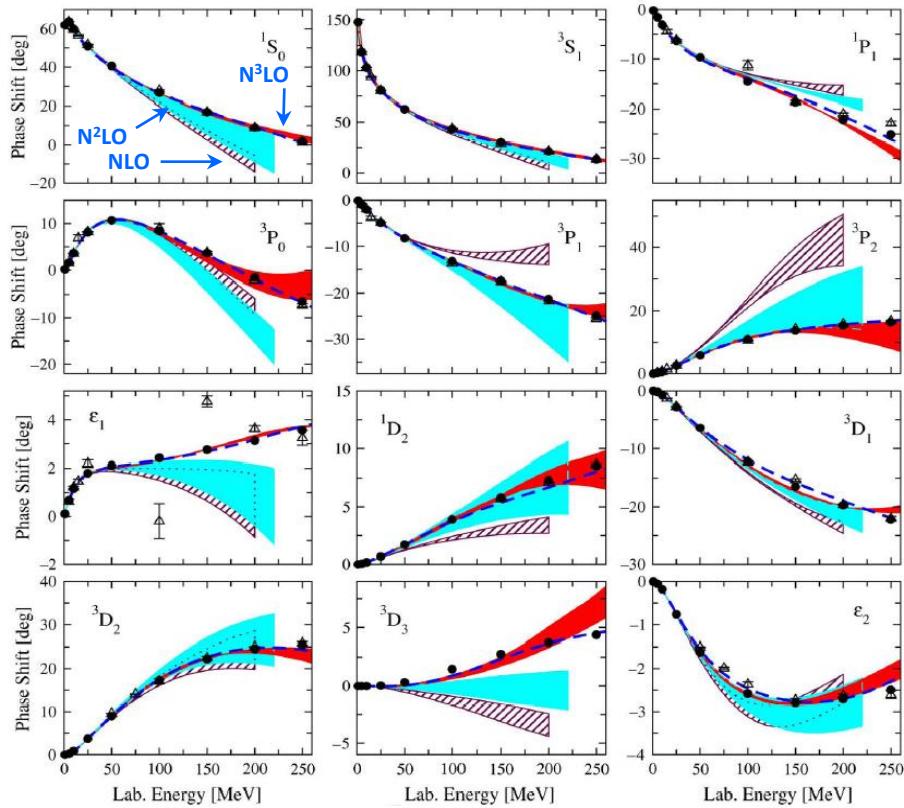
*sub-subleading 2 $\pi$ -exchange*



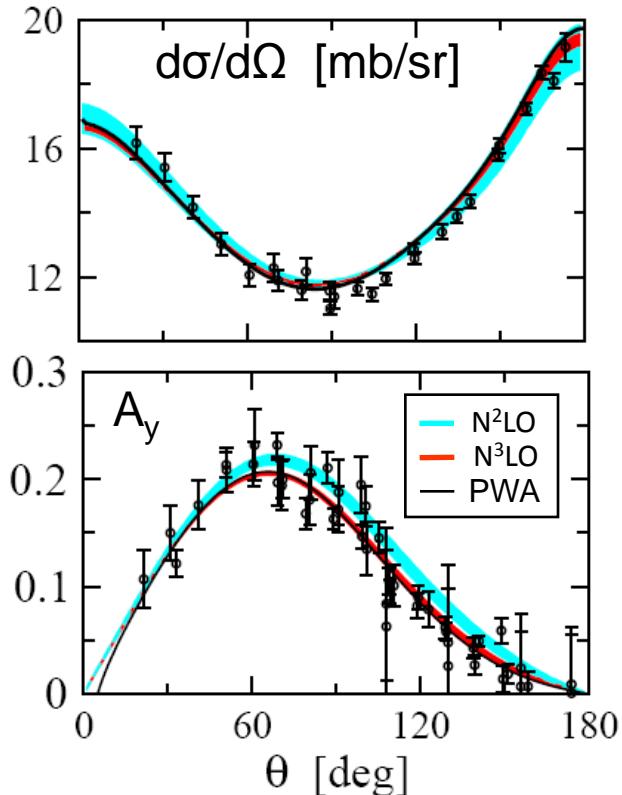
*3 $\pi$ -exchange (small)*

+ 1/m and isospin-breaking corrections...

## Neutron-proton phase shifts up to N<sup>3</sup>LO



## np scattering at 50 MeV



## Deuteron binding energy & asymptotic normalizations A<sub>s</sub> and η<sub>d</sub>

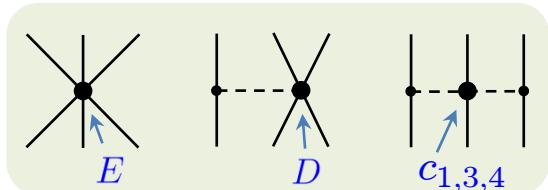
	NLO	N <sup>2</sup> LO	N <sup>3</sup> LO	Exp
E <sub>d</sub> [MeV]	-2.171 ... -2.186	-2.189 ... -2.202	-2.216 ... -2.223	-2.224575(9)
A <sub>s</sub> [fm <sup>-1/2</sup> ]	0.868 ... 0.873	0.874 ... 0.879	0.882 ... 0.883	0.8846(9)
η <sub>d</sub>	0.0256 ... 0.0257	0.0255 ... 0.0256	0.0254 ... 0.0255	0.0256(4)

Entem & Machleidt '03; Epelbaum, Glöckle & Meißner '05

# Three-nucleon forces

- Three-nucleon forces in chiral EFT start to contribute at NNLO

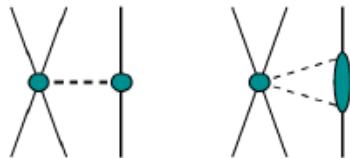
U. van Kolck '94; Epelbaum et al. '02; Nogga et al. '05; Navratil et al. '07



$c_{1,3,4}$  from the fit to  $\pi N$ -scattering data  
 $D, E$  from  $^3H, ^4He, ^{10}B$  binding energy +  
coherent  $nd$  scattering length

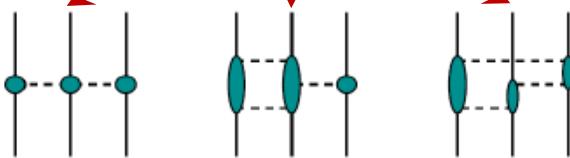
- Three-nucleon forces at  $N^3LO$

Short range contr.  
under construction



Bernard, Epelbaum, H.K. , Meißner '08; Ishikawa, Robilotta '07

Long range contributions



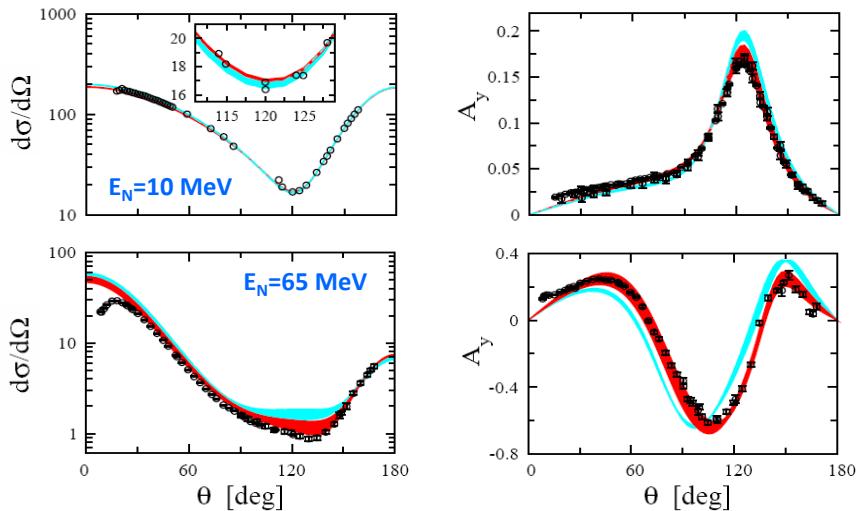
$$\begin{aligned} \text{---} &= \text{---} + \text{---} + \text{---} + \dots \\ \text{---} &= \text{---} + \text{---} + \text{---} + \dots \\ \text{---} &= \text{---} + \text{---} + \text{---} + \dots \end{aligned}$$

- No additional free parameters at  $N^3LO$
- Rich isospin-spin-orbit structure of  $N^3LO$  3NF

# Nd elastic scattering

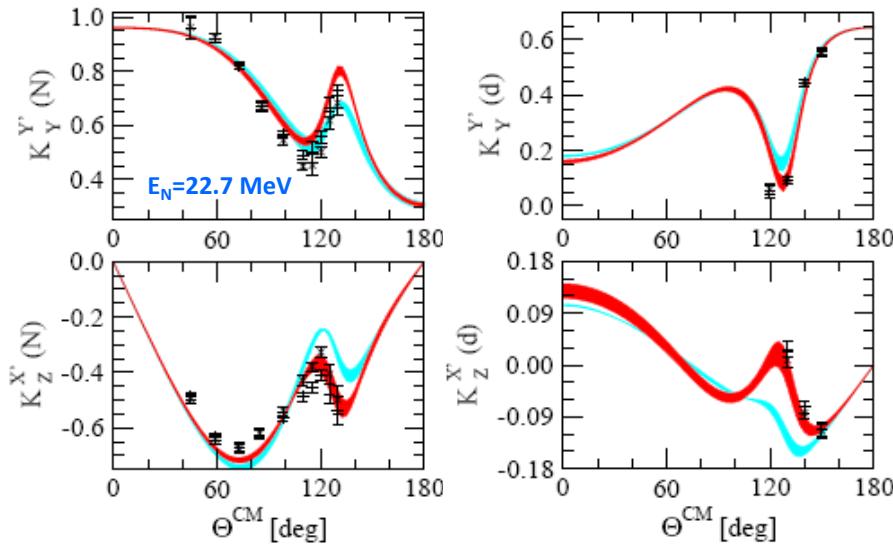
## Cross section & vector analyzing power

E.pelbaum, PPNP 57 (2006) 654



## Polarization transfer coefficients

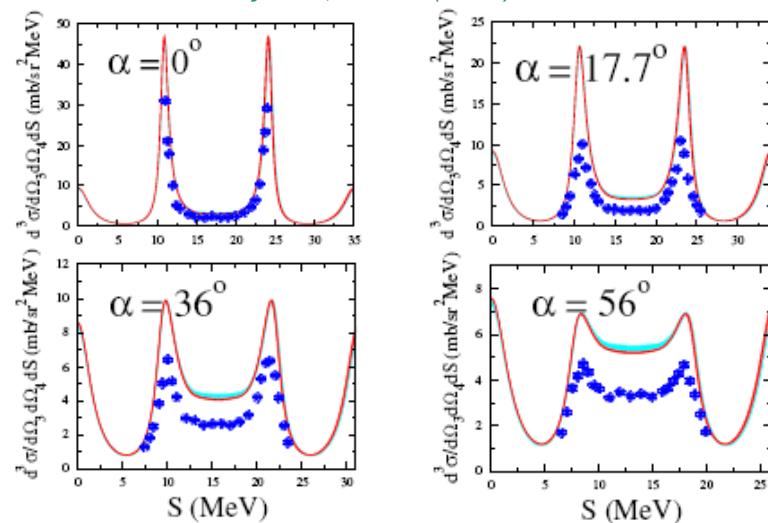
Witala et al., PRC 73 (2006) 044004



# Deuteron break-up

## SCRE configuration at $E_d=19$ MeV

Ley et al., PRC 73 (2006) 064001



- Promising NNLO results for Nd elastic scattering
- Satisfactory  $A_y$  description related to overprediction of triplet P-waves
- Systematic overestimation of deuteron break-up data
- Hope for improvement at N<sup>3</sup>LO

# Delta-less effective potential

- Standard chiral expansion:  $Q \sim M_\pi \ll \Delta \equiv m_\Delta - m_N = 293 \text{ MeV}$
- Small scale expansion:  $Q \sim M_\pi \sim \Delta \ll \Lambda_\chi$  (Hemmert, Holstein & Kambor '98)
- Delta contributions encoded in LECs (Bernard, Kaiser & Meißner '97)

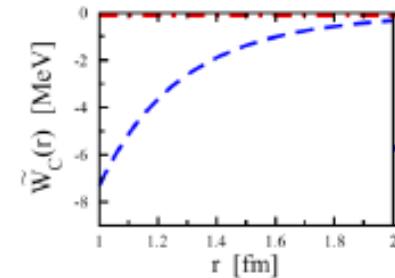
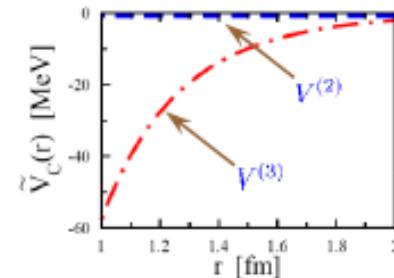
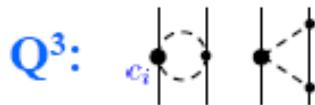
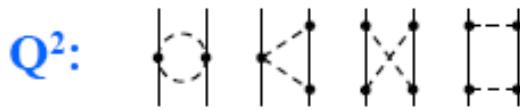


Delta-resonance saturation

$$c_3 = -2c_4 = c_3(\Delta) - \frac{4h_A^2}{9\Delta}$$

Enlargement due to  
Delta contribution

- Convergence of EFT potential



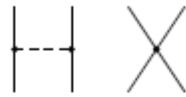
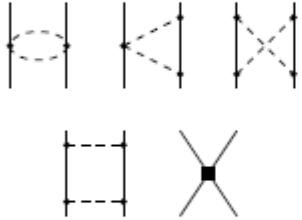
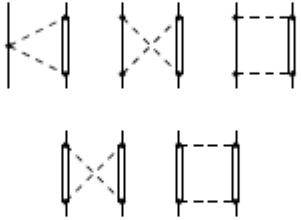
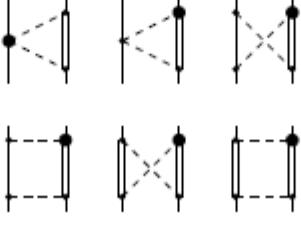
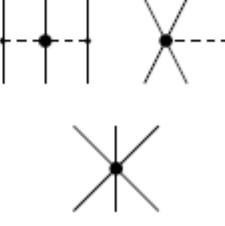
The subleading contribution is bigger than the leading one!

Expectation from inclusion of  $\Delta$  explicitly

- more natural size of LECs
- better convergence
- applicability at higher energies

# Few-nucleon forces with the Delta

Isospin-symmetric contributions

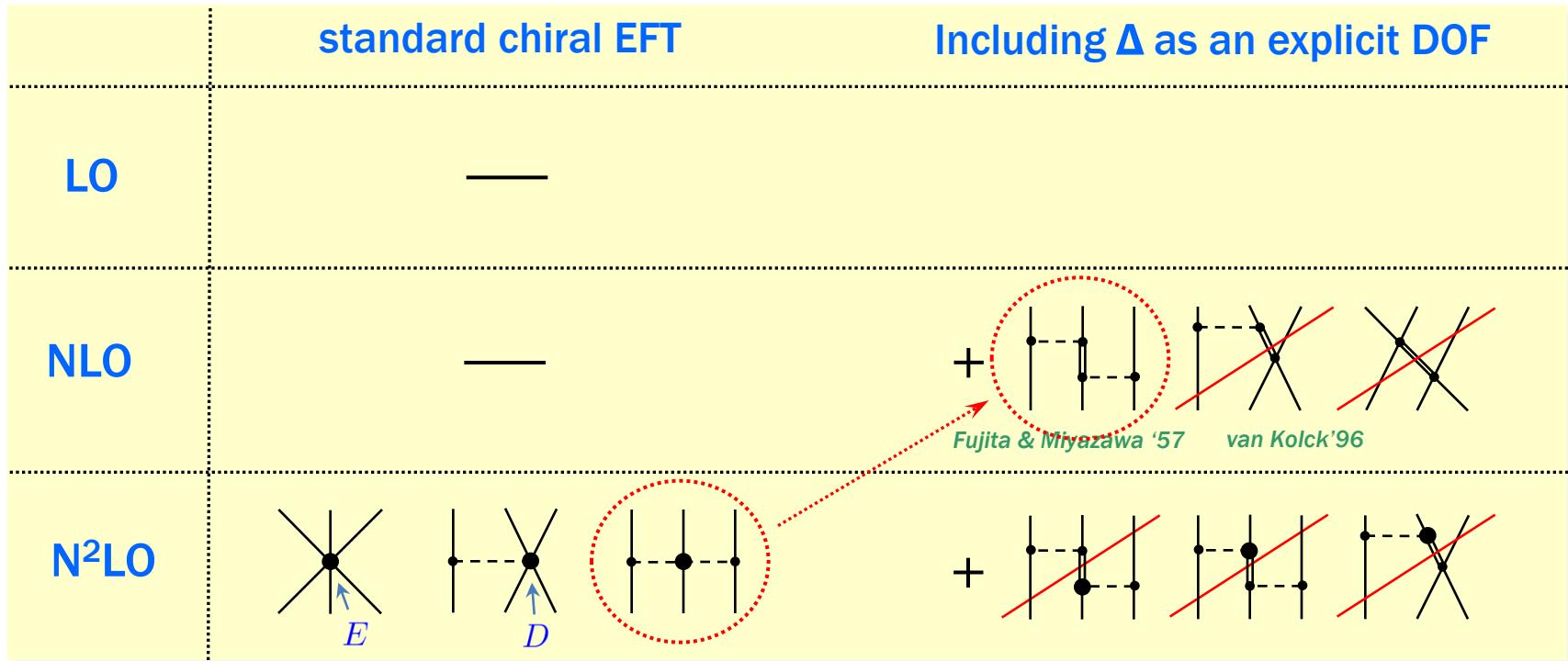
	<i>Two-nucleon force</i>		<i>Three-nucleon force</i>	
	$\Delta$ -less EFT	$\Delta$ -contributions	$\Delta$ -less EFT	$\Delta$ -contributions
<b>LO</b>		X	—	—
<b>NLO</b>			—	
<b>NNLO</b>				—

Ordonez et al.'96, Kaiser et al. '98

H.K., Epelbaum & Meißner '07

# Delta excitations and the three-nucleon force

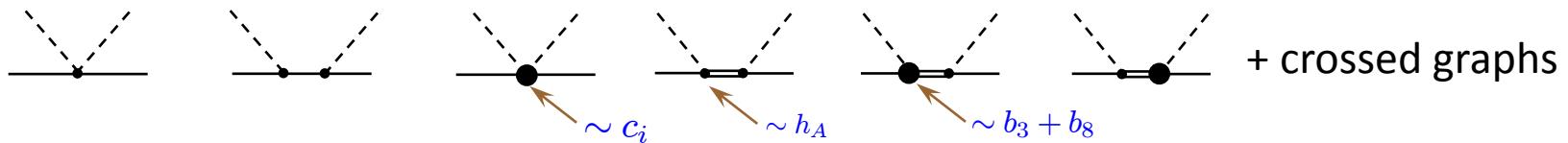
Epelbaum, H.K., Meißner, Nucl. Phys. A806 (2008) 65



- The LO NNN $\Delta$  contact interaction  $\bar{T}_i^\mu N \bar{N} S_\mu \tau^i N + \text{h.c.}$  vanishes due to the Pauli principle  
→ the LECs  $D$  and  $E$  are not saturated by the delta.
- No contributions from subleading  $2\pi$  –exchange due to  $\partial^0$  at the  $b_3 + b_8$  vertex.
- The entire effect of the  $\Delta$  is given by a partial shift of the  $N^2LO$  TPE 3NF to NLO...

# Fit of LECs to $\pi N$ scattering data

Epelbaum, H.K., Meißner, Eur. Phys. J. A32 (2007) 127



Two possible values for  $h_A$

- fit 1:  $h_A = \frac{3g_A}{2\sqrt{2}} \sim 1.34$  (**SU(4), large  $N_c$** )
- fit 2:  $h_A = 1.05$  (**Fettes & Meißner '01**)

LECs	$Q^2$ , no $\Delta$	$Q^2$ , fit 1	$Q^2$ , fit 2
$c_1$	-0.57	-0.57	-0.57
$c_2$	2.84	-0.25	0.83
$c_3$	-3.87	-0.79	-1.87
$c_4$	2.89	1.33	1.87
$h_A$	-	1.34*	1.05*
$b_3 + b_8$	-	1.40	2.95

Results of the fit

- Improved description of P-wave parameters when  $\Delta$  is included
- Strongly reduced values for  $c_i$
- Resulting  $c_i$  depend strongly on  $h_A$  while the thresh. param. do not

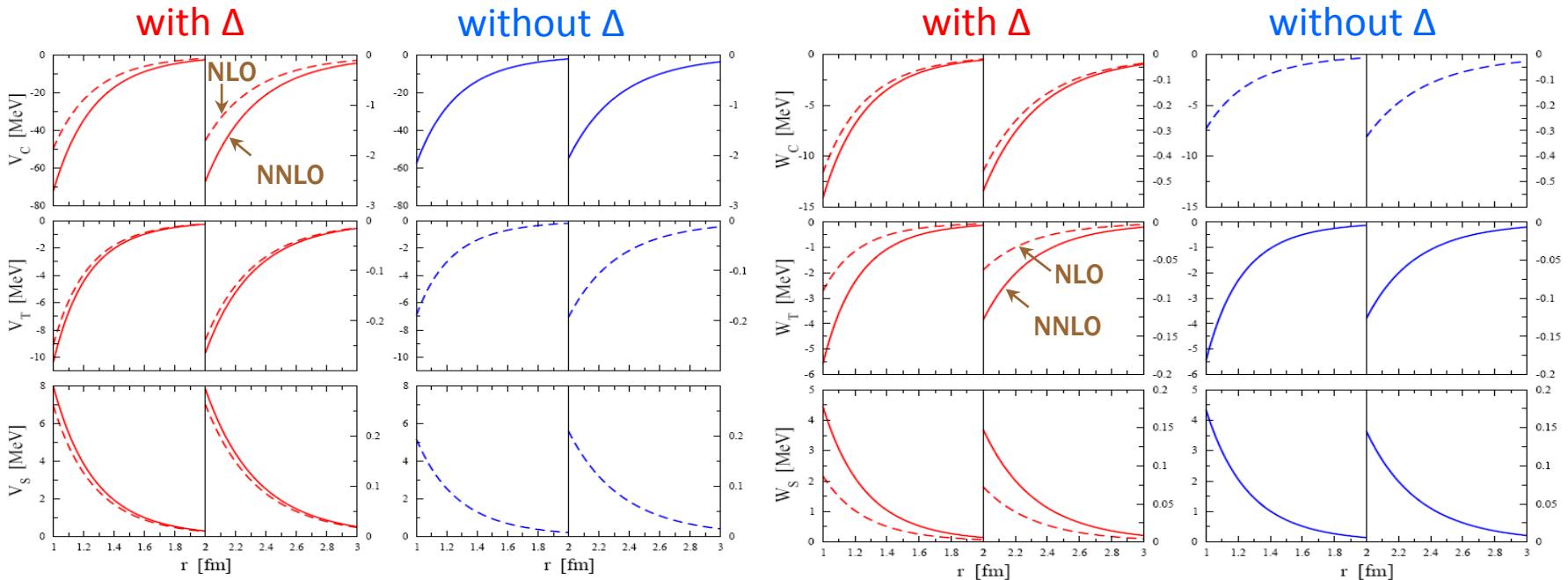
S- and P-wave threshold parameters

	$Q^2$ , no $\Delta$	$Q^2$ fits 1, 2	EM98
$a_{0+}^+$	0.41	0.41	$0.41 \pm 0.09$
$b_{0+}^+$	-4.46	-4.46	-4.46
$a_{0+}^-$	7.74	7.74	$7.73 \pm 0.06$
$b_{0+}^-$	3.34	3.34	1.56
$a_{1-}^-$	-0.05	-1.32	$-1.19 \pm 0.08$
$a_{1-}^+$	-2.81	-5.30	$-5.46 \pm 0.10$
$a_{1+}^-$	-6.22	-8.45	$-8.22 \pm 0.07$
$a_{1+}^+$	9.68	12.92	$13.13 \pm 0.13$

# NN potential with explicit $\Delta$

$$V_{\text{eff}} = V_C + W_C \vec{\tau}_1 \cdot \vec{\tau}_2 + [V_S + W_S \vec{\tau}_1 \cdot \vec{\tau}_2] \vec{\sigma}_1 \cdot \vec{\sigma}_2 + [V_T + W_T \vec{\tau}_1 \cdot \vec{\tau}_2] (3 \vec{\sigma}_1 \cdot \hat{r} \vec{\sigma}_2 \cdot \hat{r} - \vec{\sigma}_1 \cdot \vec{\sigma}_2)$$

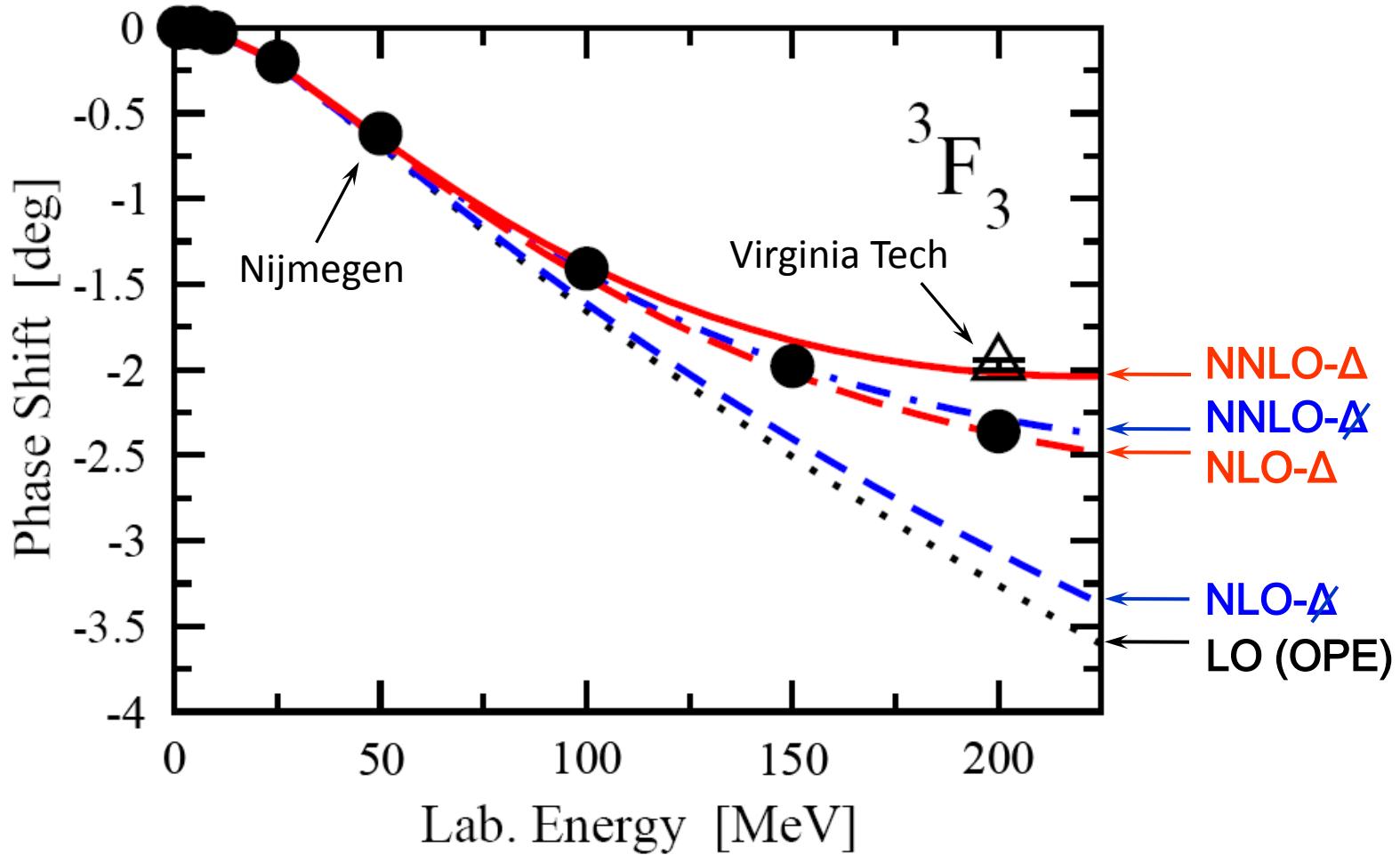
Chiral  $2\pi$ - exchange potential up to NNLO



Advantages when  $\Delta$  is included explicitly

- Dominant contributions already at NLO
- Much better convergence in all potentials

## $^3F_3$ partial waves up to NNLO with and without $\Delta$



(calculated in the first Born approximation)

# $\Delta$ -mass splitting in chiral EFT

Epelbaum, H.K., Meißner, Nucl. Phys. A806 (2008) 65

Tiburzi, Walker-Loud, Nucl. Phys. A 764 (2006) 274 (strong splitting)

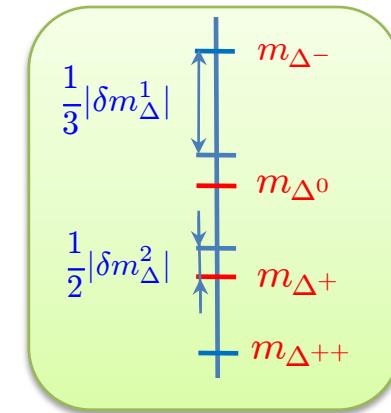
$$\mathcal{L}_{\Delta, \text{mass}}^{\text{LO}} = -\bar{T}_i^\mu \left[ -\delta m_\Delta^1 \frac{1}{2} \tau^3 \delta_{ij} - \delta m_\Delta^2 \frac{3}{4} \delta_{i3} \delta_{j3} \right] g_{\mu\nu} T_j^\nu.$$

Equidistant splitting: strong & em

$$\delta m_\Delta^1 = -4M_\pi^2 \epsilon c_5^\Delta - F_\pi^2 e^2 f_2^\Delta$$

Non-equidistant splitting: em

$$\delta m_\Delta^2 = -\frac{4}{3} F_\pi^2 e^2 f_2^\Delta$$



- Most recent data from the PDG:  $m_{\Delta^{++}} = 1230.80 \pm 0.30$  MeV,  $m_{\Delta^0} = 1233.45 \pm 0.35$  MeV

In addition, PDG's recommended value for the average mass:

$$m_\Delta = \frac{1}{4} (m_{\Delta^{++}} + m_{\Delta^+} + m_{\Delta^0} + m_{\Delta^-}) = \tilde{m}_\Delta + \frac{1}{4} \delta m_\Delta^2 = 1231 \dots 1233 \text{ MeV}$$

On the other hand:  $m_\Delta = 1233.4 \pm 0.4$  MeV (Arndt et al. '06)  $\rightarrow$  use:  $m_\Delta = 1233$  MeV

$\rightarrow \boxed{\tilde{m}_\Delta = 1233.4 \pm 0.7 \text{ MeV}, \quad \delta m_\Delta^1 = -5.3 \pm 2.0 \text{ MeV}, \quad \delta m_\Delta^2 = -1.7 \pm 2.7 \text{ MeV}}$

- Alternatively, use  $m_{\Delta^{++}}/m_{\Delta^0}$  & the QM relation:  $m_{\Delta^+} - m_{\Delta^0} = m_p - m_n$  (Rubinstein et al. '67)

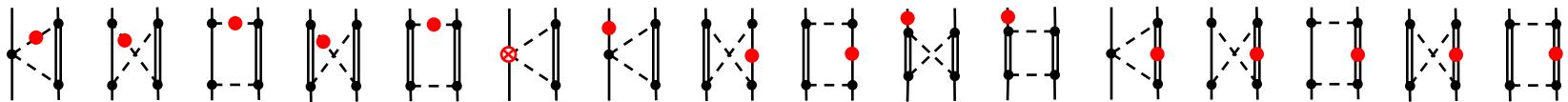
$\rightarrow \tilde{m}_\Delta = 1232.7 \pm 0.3 \text{ MeV}, \quad \delta m_\Delta^1 = -3.9 \text{ MeV}, \quad \delta m_\Delta^2 = 0.3 \pm 0.3 \text{ MeV}$

# Isospin-breaking NN potential

Epelbaum, H.K., Meißner, Phys. Rev. C77 (2008) 034006

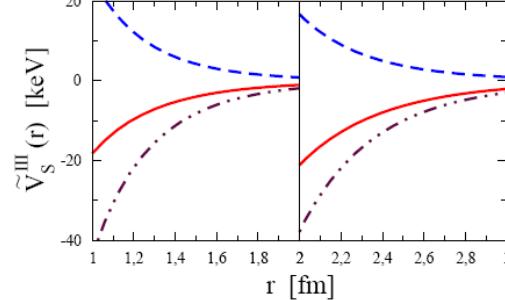
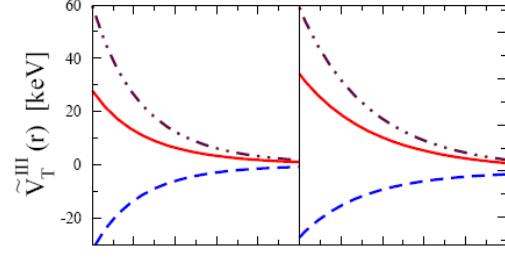
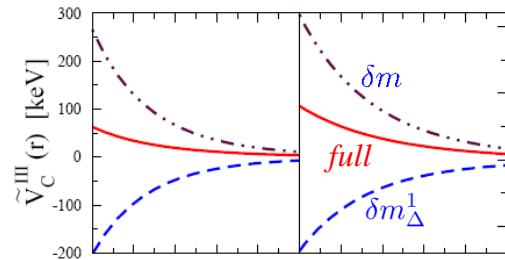
$2\pi$  – exchange contributions with explicit  $\Delta$

$$V = (\tau_1^3 + \tau_2^3) [V_C^{\text{III}} + V_S^{\text{III}} \vec{\sigma}_1 \cdot \vec{\sigma}_2 + V_T^{\text{III}} \vec{\sigma}_1 \cdot \vec{q} \vec{\sigma}_2 \cdot \vec{q}] + \dots$$

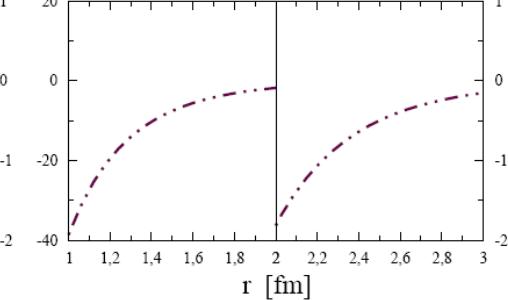
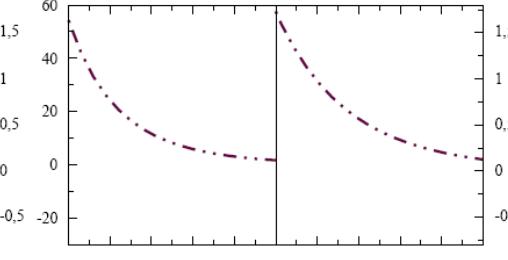
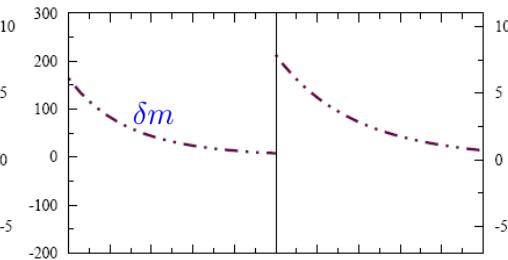


Charge-symmetry-breaking  $2\pi$ -exchange potential

$\Delta$ -full EFT, LO



$\Delta$ -less EFT, NLO



- Similar  $\sim \delta m$  contr. to  $\tilde{V}_{S,T}^{\text{III}}$  in the  $\Delta$ -less and  $\Delta$ -full EFT
- Sizeable deviation in  $\sim \delta m$  contr. for  $\tilde{V}_C^{\text{III}}$
- Strong cancellations between  $\sim \delta m$  and  $\sim \delta m_{\Delta}^1$  terms

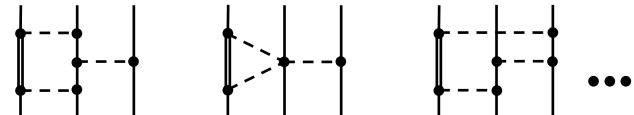


Big contributions beyond the subleading corrections in the  $\Delta$ -less EFT

# Summary

- Quantitative description of few-nucleon forces within chiral EFT
- Promising results for few-nucleon dynamics
- Better convergence of nuclear forces if  $\Delta$ -isobar is included explicitly

# Perspectives

- Complete numerical studies with 3NF and 4NF upto  $N^3LO$
- Few-nucleon forces with  $\Delta$ -DOF upto  $N^3LO$ :  
important effects expected for 3NF and 4NF 
- Electroweak reactions with few-nucleon systems